

# Gate-controlled Transport Properties in Dilute Magnetic Semiconductor (Zn, Mn)O Thin Films

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**Ionic liquid (IL) gating of functional oxides has drawn significant attentions, since it can provide reversible changes in carrier concentration ( $\sim 10^{14} \text{ cm}^{-3}$ ) at the interface, permitting the manipulation of electrical and magnetic properties of oxide films with low voltages. In this work, we demonstrated the electric-field manipulation of transport properties in dilute magnetic oxide (DMS) of  $\text{Zn}_{0.98}\text{Mn}_{0.02}\text{O}$  (MZO) using an electric-double-layer transistor (EDLT) geometry through the ionic liquid electrolyte. The MZO EDLTs exhibited a reversible control of resistance up to 33 % at 230 K. Moreover, magnetoresistance (MR) measurement revealed the influence of gate voltage ( $V_g$ ) on the magnetotransport behavior, which exhibited a positive MR in the low field region and a negative-MR in high magnetic field (up to 9 T). An increase in low-field positive MR ( $< 1$  T) upon switching the  $V_g$  from -2 to 2 V implied an enhanced ferromagnetic state of MZO due to increased electron carrier concentration. The present results therefore demonstrated controllable carrier concentration by electric-field effect plays an important role in the manipulation of magnetism in MZO.**

*Index Terms*— Electric effect, Mn-doped ZnO, DMS and magnetism

## I. INTRODUCTION

**E**LECTRIC-FIELD manipulation of magnetic properties provides a promising way for preparing fast and compact data storage devices [1-4]. Recently, intensive studies are made on the control of magnetism in artificial multiferroic heterostructures through the manipulation of charge carriers or inducing strains [4-11]. However, these heterostructure are often limited by the compatibility between the ferromagnetic and ferroelectric components, such as lattice mismatching issue and interfacial diffusion due to high deposition temperatures [7, 11]. Moreover, for observable tunability of magnetism one often requires thin ferroelectric layers or large-switching voltages for complete reversal of the polarization directions [12], and often results in the breakdown of the ferroelectric layer [6].

For low-switching voltages and highly tunability of carrier concentrations in the ferromagnetic interface, ionic-liquid (IL)-gating modulation has demonstrated interfacial carrier concentration changes up to  $10^{14} \text{ cm}^{-2}$  in ferromagnets, which is even larger than the tuning by electrostatic doping via magnetoelectric coupling [13-17]. The ionic-liquid-gating technique has been widely deployed for achieving high tunability in different systems such as metal-to-insulator transition in the complex oxides [13, 14], controlling superconductivity behavior in two-dimensional materials [18] and exchange-bias effect in antiferromagnets [19].

Generally, the electric-field control of carrier density is more effective in manipulating semiconductor as compared metallic systems, since the carrier density in semiconductors is 2-3 order of magnitude smaller [1]. A high tunability in carrier

density is expected as the screening depth is increased. Among the candidates of ferromagnetic semiconductors, manganese-substituted zinc oxide (MZO) has the advantage that the carrier concentration is independent of the total moment of MZO film since the zinc atom is isovalent with manganese [20]. Previous studies have shown that the magnetic moment of MZO can be controlled via manganese doping or the oxygen content inside the film [20, 21].

In this work, we report the fabrication and characterization of electric-double-layer transistors (EDLTs) for electric-field manipulation of magnetotransport in *n*-type MZO. The devices were characterized by several techniques, including X-ray diffractometry (XRD), atomic force microscopy (AFM), Hall measurements, resistance-switching tests and magnetotransport measurements. Our results showed the gate voltage ( $V_g$ ) manipulation in the accumulations and depletion of charge carriers of the MZO films through a low magnitude ( $\leq 2\text{V}$ )  $V_g$ . Moreover, our results also demonstrate that the low-temperature magnetotransport in MZO can be manipulated by  $V_g$ .

## II. EXPERIMENTAL PROCEDURE

10 nm-thick  $\text{Zn}_{0.98}\text{Mn}_{0.02}\text{O}$  (hereafter MZO) thin films were deposited on (0001)  $\text{Al}_2\text{O}_3$  single crystal substrates by pulsed laser deposition. The film deposition was performed at 300°C with an oxygen pressure of  $5 \times 10^{-4}$  Pa. The as-grown films were cooled down to room temperature in the same oxygen ambient at a rate of  $\sim 10^\circ\text{C}/\text{min}$  for promoting the oxygen stoichiometry of the deposited films. Thickness of as-grown films was estimated by the number of laser pulses used and was further confirmed with AFM.

For electric-field manipulation of MZO devices, the films were patterned into Hall-bars (channel width: 50  $\mu\text{m}$ , channel length: 110  $\mu\text{m}$ ) by photolithography and wet etching, using dilute HCl (1% HCl, etching rate: 60 nm/min) as etchant [22]. Fig. 1(a) depicts the schematic of the device employed in this

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study. The MZO thin film Hall bar devices were capped by 2 nm of  $\text{SiO}_2$  to prevent the direct contact between the MZO film and the IL electrolyte (N, N-diethyl-N-(2-methoxyethyl)-N-methylammonium-bis-(trifluoromethylsulfonyl)-imide, DEME-TFSI). Coplanar electrodes for IL and contact electrodes for MZO films were prepared by depositing Ti (5 nm)/Au (50 nm) using electron-beam-evaporation. Prior to the gating experiment, IL was baked at  $80^\circ\text{C}$  in vacuum to get rid of water contamination: it is known that absorbed water in IL not only affects the performance of EDL layer, but may also modify the surface of gated materials [23, 24]. To avoid the water contamination, the devices were immediately cool down to 230 K in helium or vacuum environment before applying any  $V_g$ .

Magnetoresistance and resistance-switching behavior of the devices were measured in a physical property measurement system (PPMS) with a constant current of  $1\ \mu\text{A}$  using the four-point-probe geometry.  $V_g$  on IL was applied using a Keithley 2400 sourcemeter. Transport measurements in the temperature range of 10 K to 180 K were performed after maintaining  $V_g$  for 60 mins at 230 K and removing the  $V_g$  at 180 K; since the ionic mobility of IL at 230 K is low, a prolonged duration is required to stabilize the EDL layer. The freezing point of IL used in this work is 210 K, below which the electric field induced in the EDL cannot be dynamically changed. Reheating devices above the freezing point (230 K) is essential to alter the carrier concentration of the MZO surface via the application of  $V_g$ .

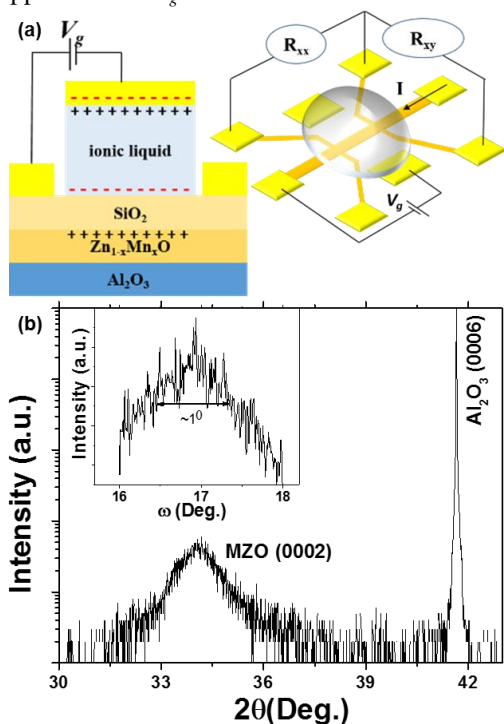


Fig. 1. (a) Schematic diagram of MZO EDLT device (left), and the Hall-bar device configuration with ionic liquid electrolyte as top electrode (right). (b) XRD  $2\theta$  scan of 10 nm-thick MZO thin film. Inset:  $\omega$ -scan of (0002) MZO peak.

### III. RESULTS AND DISCUSSIONS

Fig. 1(b) presents the XRD scan of MZO (10 nm) grown on  $\text{Al}_2\text{O}_3$  (0001) substrate. Only the (0002) diffraction peak of MZO was observed, indicating that highly  $c$ -axis-oriented MZO films were grown. No impurity phase such as  $\text{Mn}_2\text{O}_3$ ,  $\text{MnO}_2$  and  $\text{Mn}_3\text{O}_4$  was observed. The  $c$ -axis lattice constant ( $5.26\ \text{\AA}$ ) of the MZO film matches well with the corresponding lattice constant of reported PLD-grown Mn-substituted ZnO films [25, 26]. Fig. 1 (c) shows the XRD  $\omega$ -scan of (0002) MZO diffraction peak, with a full-width at half-maximum (FWHM) of  $1^\circ$ .

Fig. 2(a) illustrates the  $V_g$ -dependent longitudinal resistance ( $R_{xx}$ ) of MZO EDLT, showing the reversible switching between two distinct resistance levels. The results were collected at 230 K, at which the electrochemical reaction between IL and MZO can be suppressed while the ionic mobility of IL at 230 K is enough to induce accumulation or depletion of electron charge densities at MZO surfaces [27].  $\Delta R_{xx}$  increases 33% upon switching  $V_g$  from 2 V to -2 V, which is consistent with the scenario of accumulated (depleted) negative charge carriers at the interfaces between IL and MZO [28]. Moreover, the resistance switching behavior in MZO EDLTs exhibits good dynamic response and reversible manipulation via the  $V_g$  reversal. Such modulations of  $R_{xx}$  are due to electron charge movement in MZO rather than the contribution of gate current: the drain-source current is higher than the gate-source current by at least two orders of magnitude, which is shown in Fig. 2(b).

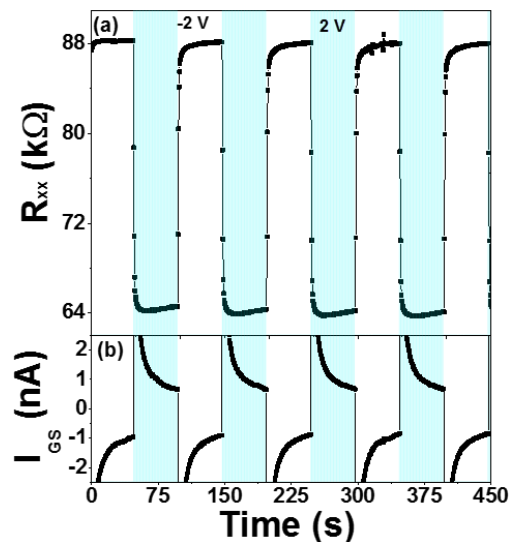


Fig. 2. Time profile of (a) longitudinal resistance ( $R_{xx}$ ) and (b) gate-source current ( $I_{GS}$ ) of MZO EDLT with  $V_g = 2\ \text{V}$  (highlighted region) and  $-2\ \text{V}$ .

High- to low-resistance states of MZO EDLT are accompanied by switching  $V_g$  from  $-2\ \text{V}$  to  $2\ \text{V}$ , resulting in the modulation of transport behavior in MZO. To identify that, we carried out measurements of temperature dependent resistivity ( $\rho$ ) and carrier concentration ( $\eta$ ) of MZO EDLTs using the Hall bar configuration. Although a 2-nm  $\text{SiO}_2$  protective layer is present in between IL and MZO, the changes in carrier concentration so induced is still much larger

than simply using thick dielectric layers [19, 29]. Fig. 3(a) and (b) show clear correlations of  $\rho$  and  $\eta$  of MZO with different polarities of  $V_g$ . As the temperature decreases from 180 K to 20 K, the resistivity for +2 V-poled device ( $\rho(2V)$ ) is smaller than that poled by  $V_g = -2$  V ( $\rho(-2V)$ ). Regarding the polarity direction of IL switching from +2 V and -2 V, the change in resistivity ( $\Delta\rho$ , defines by  $(\Delta\rho)/\rho=(\rho(2V)-\rho(-2V))/(\rho(-2V))$ ) reaches 3.7 % at 180 K and decreases to 3% at 20 K. Moreover, the Hall coefficient of MZO is negative for all temperatures, implying an electron-doping characteristic of the MZO films. The depletion or accumulation of electron carriers in MZO EDLTs by electric-field effect can be reflected from the change in carrier concentration with different gating voltages.

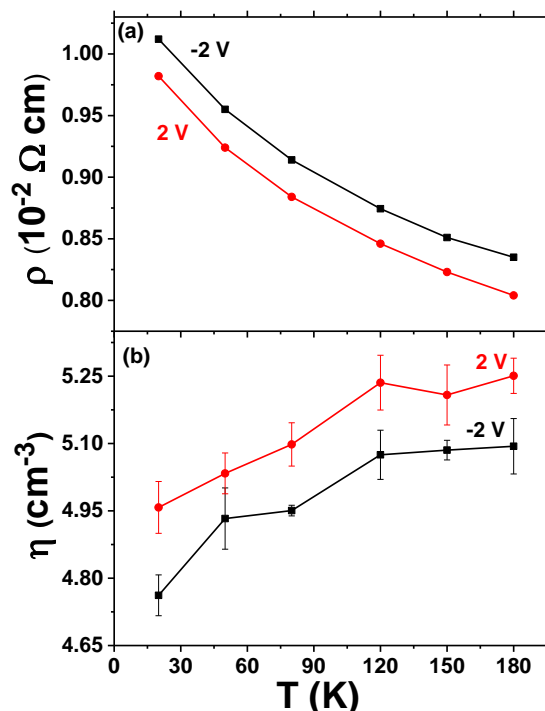


Fig. 3. Temperature dependence of resistivity (a) and carrier concentration (b) with  $V_g = -2$  V (black squares) and 2V (red dots).

To understand the nature of electric-field-induced magnetotransport behavior, we focus on the application of  $V_g$  in the cooling process. Fig. 4 displays the magnetoresistance (MR) curves with out-of-plane applied field for MZO EDLTs when  $V_g$  are set at -2, 0 and 2 V in the cooling process. Here MR is defined as  $\text{MR}=(R_{xx}(H)-R_{xx}(0))/R_{xx}(0)$ , where  $R_{xx}(H)$  and  $R_{xx}(0)$  are the  $R_{xx}$  values with external magnetic fields of  $H$  and zero, respectively. From the comparison in Fig. 4(a), we note that the application of  $V_g$  in the cooling process triggers the manipulation of magnetotransport at 10 K. When the  $V_g$  set to 0 V, the MR exhibits an increasing trend from 0 to 2.8 T, and reaches 0.4% at 2.8 T. The trend of MR becomes a decrease trend from 2.8 to 9 T, and reaches -3% at 9 T. As the negative polarity of  $V_g$  is applied, the MR almost traces similar behavior as compared with MR with 0V while the positive MR disappears. After the  $V_g$  set 2 V, the change in the peak value of positive and negative MR is observed. The peak

positive MR increases from 0 to 1.8% and the negative-MR (measured at 9 T) decreases from -4.5% to -0.6% when  $V_g$  increases from -2 to 2 V. As shown in Fig. 4b, negative MR greatly decreases with increasing temperature and become insensitive towards  $V_g$  modulations.

The positive MR behaviors at low temperature has been attributed to the spin splitting of  $s$ - $d$  exchange coupling between carrier and localized electrons conduction band in the DMS, leading to the increase in MR at the low-field regime[30, 31]. On the other hand, the negative MR originates from the interaction between localized electrons and the orbitals of  $3d$  transition metal ions ( $\text{Mn}^{2+}$ ), which forms the bound magnetic polarons[32]. As the temperature decreases, the localized electron contributing from  $V_o$  tends to freeze out, and a hopping channel is established. With certain magnetic field, the random orientation of magnetic polarons tends to align and results in a decrease in resistance of MR measurement. Moreover, negative MR at high-field regime is due to magnetic scattering in  $3d$  transition ions, which contributes to the paramagnetic or antiferromagnetic signals. It is well known that not all magnetic ions form magnetic polarons, while some of magnetic ions possess in paramagnetic and antiferromagnetic signal through the mediation of  $V_o$ [32].

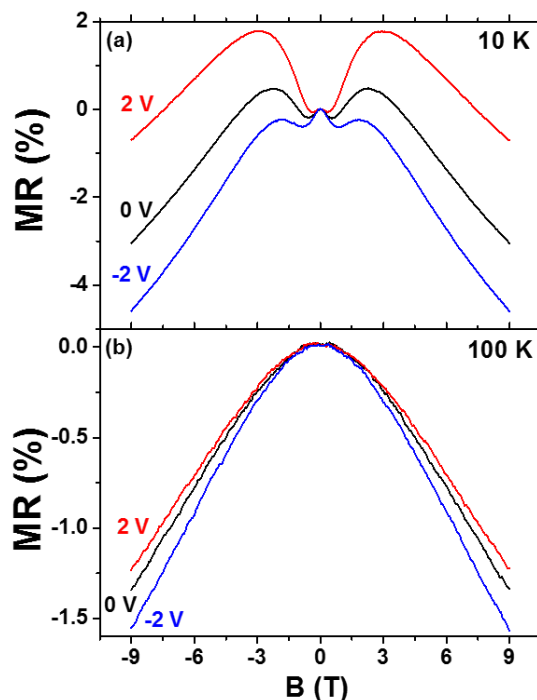


Fig. 4. Magnetoresistance of MZO EDLTs for  $V_g = -2, 0$  and 2 V measured at (a) 10 K and (b) 100 K.

Comparing the MR values and change in carrier concentration for different polarity of  $V_g$ , an apparent increase in negative MR was observed as the  $V_g$  switched from -2 to 2 V. At 10 K, increase in carrier concentration of MZO EDLTs lead to increase the positive MR at low-field regime and decreases the negative MR. The negative MR at  $T = 100$  K, however, only slightly changes upon the  $V_g$  reversal from -2 to 2 V. These strong dependences between carrier concentration

and magnetotransport is consistent with the previous report in Mn:ZnO film grown on  $\text{Pb}(\text{Mg}_{1/3}\text{Nb}_{2/3})\text{O}_3$ - $\text{PbTiO}_3$  [11]. The impurity band model in DMS proposed by Coey *et al.* that an increase in electron concentration lead to increase the interaction between each polaron, and finally enhance the ferromagnetism in the MZO EDLTs [33]. This increase in ferromagnetism is evidenced by electric-field manipulation in magnetoresistance ( $T < 100$  K). Moreover, the required voltage for manipulating the magnetism in MZO is much less than the previous reports [11, 30, 34]. The present results therefore demonstrate electric-field controllable accumulation/depletion of electron charge carrier in MZO EDLTs through low-voltage gating process, which plays an important role in the manipulation of magnetism in the MZO.

#### IV. CONCLUSION

In summary, we fabricated and studied electric-field manipulation of magnetotransport properties of MZNO EDLTs via IL gating. Resistance-switching behavior was observed upon the application of different gate voltages, which lead to a reversible control in the transport phenomena at the interface. The MZO EDLTs showed strong dependence between magnetoresistance and  $V_g$ . This result also indicated that the ferromagnetism of MZO can be affected by the application of external electric field. Our results suggest that electric-field manipulation through IL gating is a promising method to reversibly control the electrical and magnetic properties of oxide-based materials.

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