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 C_1 **-Symmetric)(C^N²)(O^O)]-***Tris***-Heteroleptic Iridium(III)- Complexes with the Preferentially Horizontal Orientation for High-Performance Near-Infrared (NIR) Organic Light-Emitting Diodes (OLEDs)**

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Abstract

The development of high-efficiency NIR-emitting Ir(III)-complex-based phosphors for reliable NIR-OLEDs (near-infrared organic light-emitting diodes) is still a formidable challenge. Herein, a molecule-engineered approach is developed to afford three *C*1-symmetric [Ir(C^N¹)(C^N²)(O^O)]-tris-heteroleptic Ir(III)-complexes ([Ir(iqbt)(dFppy)(acac)] (**1**), [Ir(iqbt)(ppy)(acac)] (**2**) and [Ir(iqbt)(dpqx)(acac)] (**3**)), whose good NIR-luminescent efficiency (*Φ*PL = 0.18 for **1** (*λ*em = 703 nm), 0.26 for **2** (*λ*em = 715 nm) or 0.28 for **3** (*λ*em = 707 nm)) originates from the strengthened 3 MLCT contribution (MLCT = metal-to-ligand charge transfer). Moreover, the quantitative molecular orientation determination of their doped emitting layers (EMLs) reveals that the preferential horizontal orientation of the emitting dipoles is beneficial to their highly efficient NIR-OLEDs with light out-coupling. Especially for the **NIR-OLED-2** with $\lambda_{\text{em}} = 715$ nm, its high performance ($\eta_{\text{EQE}}^{\text{Max}} = 5.30\%$ and negligible (< 2%) efficiency-roll-off) among the reported solution-processed NIR-OLEDs based on Ir(III)-complexes at the similar color gamut is realized. This result shows that *C*1 symmetric $[\text{Ir}(C^N)^1](C^N)^2](O^N)$]-*tris*-heteroleptic Ir(III)-complexes can provide a new platform to low-cost and large-area scalable NIR-OLEDs.

1. Introduction

Inspired by the emerging importance of near-infrared (NIR; $\lambda_{\rm em} > 700$ nm) organic lightemitting diodes (OLEDs) in night-vision display,^[1] information security,^[2] optical telecommunication^[3] and photodynamic therapy,^[4] concrete efforts^[5] have been devoted to the molecular design of NIR-emitters. In this regard, due to the concurrent harvesting of singlet and triplet excitons, NIR-emissive phosphors are more suitable candidates compared to fluorescent luminogens with an upper limit of the *η*_{IOE} (internal quantum efficiency) of 25%. Saliently for the octahedral Ir(III)-complexes with rather short triplet lifetimes, besides the high efficiency which is appreciably competitive to other triplet-based counterparts^[5] (Pt(II)complexes and TADF (thermal activated delayed fluorescence) molecules, etc.), the superiority of a significantly alleviated efficiency-roll-off renders them particular appealing^[6] to the studies of NIR-OLEDs. Nonetheless, as constrained by the so-called "energy-gap law",^[7] it remains a real challenge to develop new Ir(III)-complex dyes towards highperformance NIR-OLEDs.

To date, cyclometalated Ir(III)-complexes, possessing neutral ($[Ir(C^N)]$ 3]-homoleptic^[8] or [Ir(C^N)₂(L^X)]-heteroleptic; L^X = O^O^[9] or N^O^[10]) and cationic ([Ir(C^N)₂(N^N)]⁺)^[11] forms and acting as both the chromophore and the charge-transport medium, were demonstrated for reliable NIR-OLEDs. Apparently, in order to narrow the HOMO-LUMO (HOMO = highest occupied molecular orbital; LUMO = lowest unoccupied molecular orbital) gap of a specific Ir(III)-complex towards the restrictive NIR emission, the use of π conjugation expansion^[8-11] of the C ΔN -cyclometalated main ligand, especially with electronrich substituents, is generally preferred which can effectively destabilize the HOMO energy level. Considering the facile aggregation-caused quenching $(ACO)^{[12]}$ effect arising from the large π -conjugation of the C^{\wedge}N-cyclometalated main ligand, modification of the HL \wedge X ancillary ligand seems to be an accessible alternative way to stabilize the LUMO energy level of the $[Ir(C^N)/2(L^N)]$ -heteroleptic^[9-11] Ir(III)-complex. However, its exclusively reduced

HOMO-LUMO band-gap for NIR luminescence does not have an universal effect,^[13] and even the undesirable hypsochromic shift $[13]$ beyond the NIR regime results from the simultaneous stabilization of the HOMO energy level. Additionally, according to the "energygap law",[7] the inverse of *k*nr (non-radiative rate constant) depends exponentially on the excited-state $({}^{1}T)$ energy, and k_r (radiative rate constant) has the cubic dependence. Therefore, high photo-luminescence efficiency necessary to increase k_r and/or to reduce k_{nr} is essential for obtaining the desired NIR emission. Convincingly, these two factors, which are key to reliable NIR-OLEDs, are highly associated with the specific excited state situations of the Ir(III)-complexes (such as energy-level and electron/charge transfer (CT) among the ³MLCT/³LC (LC = ligand-centered; MLCT = metal-to-ligand charge transfer) of the ¹T state, etc.).^[14] In light of the transformation^[15] of the ¹T nature in the C_1 -symmetric *tris*-heteroleptic Ir(III)-complex composed of an Ir(III) ion and three different ligands, it is of notable interest on expanding that novel molecule-engineered strategy to Ir(III)-complex-based NIR-emitters, which should fill in the blank after previously reported NIR-emissive $[Ir(C^N)]^{8}]$ and $[\text{Ir}(C^N)/2(L^N)]$ -heteroleptic^[9-11] Ir(III)-complexes. Indeed, for *C*₁-symmetric *tris*heteroleptic Ir(III)-complexes, their evident advantage lies in the potentially enriched inventory ($[Ir(C^N)^1)(C^N)^2](L^N)$] (L^X = O^O or N^N),^[15-16] $[Ir(C^N)^1)(C^N)^2](C^N)^3]$,^[17] $[\text{Ir}(C^N\text{N}^1)(C^N\text{C}^2)(C^N\text{N}^3)]$,^[18] [Ir(C^NN¹)(N^NN¹)(N^NN²)],^[19] etc.) compared to conventional $[Ir(C^N)]3]^{-[8]}$ and $[Ir(C^N)]2(L^N)]$ -counterparts.^[9-11] Meanwhile, the breaking of the symmetry of molecular structures can help increasing the solubility of resultant Ir(III)complexes. Moreover, arising from the *C*1-symmetric *tris*-heteroleptic spatial configuration, the emitting dipole orientation (EDO) is rearranged. Especially through the specific multiligands' perturbation, the d-orbital participation in the ${}^{1}T$ state (3 MLCT contribution) can be redistributed. More importantly, for the emitting layers (EMLs) fabricated from *C*1-symmetric *tris*-heteroleptic Ir(III)-complexes, the molecular orientation is further specified, offering a promise to boost the light-coupling efficiency.^[20] Noticeably, although the effectiveness and

handiness of controlling the *fac*-[Ir(C^N)₃]-homoleptic^[21] or [Ir(C^N)₂(L^X)]-heteroleptic^[22] emitter's alignment parallel to the substrate, even without the need of micro-lens arrays, gratings or some other physical methods,^[23] were demonstrated to significantly improve their visible-light out-coupling efficiency (*η*out) by increasing the horizontal dipole ratio, this proofof-concept approach has not been carried forward to Ir(III)-complex-based NIR-emitting systems.

Herein, encouraged by the smooth color-tuning^[15-19] and good device performance^[15] within the vivid visible ($\lambda_{\text{em}} = 465{\text -}600$ nm) range of *C*₁-symmetric *tris*-heteroleptic Ir(III)complexes, it is anticipated that an attractive structural design strategy to our *C*1-symmetric [Ir(C^N¹)(C^N²)(O^O)]-*tris*-heteroleptic Ir(III)-complexes **1**-**3** (**Scheme 1**) could offer an opportunity to give NIR emission. Especially through the strengthened 3 MLCT effect from the *C*1-symmetric spatial configuration, large *k*r and/or reduced *k*nr can be motivated toward the efficient NIR-phosphorescence. Accordingly, we also expect that the photo-physical and electrochemical properties of these NIR-emitting phosphors can be smoothly governed by varying the $C^{\wedge}N^2$ ligand. Moreover, owing to the lower symmetry of the molecular configuration of the Ir(III)-complexes 1-3 than typical C_3 -symmetric-*fac*-[Ir(C^N)₃]^[21] and C_2 -symmetric-[Ir(C^N)₂(L^X)]^[22] counterparts, it is believed that more preferential horizontal orientation of the C_1 -symmetric $[\text{Ir}(C^N N^1)(C^N N^2)(O^N O)]$ -*tris*-heteroleptic Ir(III)-complexes should be beneficial for the NIR-light extraction^[20] and challenging the low efficiency limited by the so-called "energy-gap law".[7] However, related study is not yet available for *tris*heteroleptic Ir(III) complexes and was rarely reported for NIR systems. Thus, our present research renders the *C*1-symmetric *tris*-heteroleptic Ir(III)-complexes a new platform to the realization of efficient NIR-OLEDs by utilizing their preferential horizontal orientation.

Scheme 1. Synthetic scheme of the C_1 -symmetric $[\text{Ir}(C^N N^1)(C^N N^2)(O^N O)]$ -tris-heteroleptic Ir(III)-complexes **1**-**3** with different C^N² -ligands.

2. Results and Discussion

2.1. Synthesis and Characterization of C_1 -Symmetric $[\text{Ir}(C^{\wedge}N^1)(C^{\wedge}N^2)(O^{\wedge}O)]$ -Tris-**Heteroleptic Ir(III)-Complexes 1-3**

Suzuki coupling^[10c] of cost-effective 2-Cl-isoquinoline (instead of 2-Br-isoquinoline^[9e]) with benzo[*b*]thien-2-yl boronic acid gave the HC^N¹ main ligand **Hiqbt** in 73% yield. Also as shown in **Scheme 1**, straightforward metalation^[15] in the 2-ethoxyethanol/H₂O (v/v = 3:1) mixed-solvent of the **Hiqbt** and one of the HC^N² ligands (**HdFppy**, **Hppy** and **Hdpqx**) in an equimolar ratio with IrCl₃·3H₂O, gave rise to the corresponding *in situ* formed *μ*-chlorobridged dimer intermediate $[\text{Ir}(C^N N^1)(C^N N^2)(\mu - C)]_2$, which was not isolated from the inevitable interferents of $[\text{Ir}(C^N N^1)_2(\mu\text{-}Cl)]_2$ and $[\text{Ir}(C^N N^2)_2(\mu\text{-}Cl)]_2$. Further through the treatment of the *µ*-chloro-bridged dimer intermediates with **Hacac** in the presence of *t*-BuOK, the target C_1 -symmetric $(C^{\wedge}N^2)(O^{\wedge}O)]$ -*tris*-heteroleptic Ir(III)-complex [Ir(iqbt)(dFppy)(acac)] (**1**), [Ir(iqbt)(ppy)(acac)] (**2**) or [Ir(iqbt)(dpqx)(acac)] (**3**) as the corresponding desirable product, was obtained upon the subsequent chromatography

separation, respectively. Undoubtedly, in synthesizing each of the corresponding *C*1 symmetric [Ir(C^N¹)(C^N²)(O^O)]-*tris*-heteroleptic Ir(III)-complexes **1-3**, two kinds of *bis*heteroleptic Ir(III)-complexes $[Ir(C^N)^1]_2(O^N)$ (the same $[Ir(iqbt)₂(acac)]$ and [Ir(C^N²)₂(O^O)] ([Ir(dFppy)₂(acac)], [Ir(ppy)₂(acac)] or [Ir(dpqx)₂(acac)]) were also isolated as the by-products and structurally confirmed by spectroscopic characterization (shown in **ESI**), respectively. Worthy of note, probably contributing from the promoted solubility of the $[\text{Ir}(C^N\text{N}^1)(C^N\text{N}^2)(O^N\text{O})]$ -tris-heteroleptic Ir(III)-complexes **1-3** with the *C*₁-symmetricity, their resultant yields (21-23%) after isolation are desirably acceptable.

The *C*1-symmetric *tris*-heteroleptic Ir(III)-complexes, which are soluble in common organic polar or non-polar solvents, were well-characterized by elemental analysis, FT-IR, ¹H NMR and ESI-MS (see **ESI**). Especially in the ¹H NMR spectra (**Figure S1**) of the Ir(III) complexes **1-3**, the proton resonances (δ = 5.71-9.02 ppm for **1**; 6.18-8.92 ppm for **2** and 6.27-9.10 ppm for 3) of both the $(iqbt)$ and the corresponding $(dFppy)/(ppy)/(dpqx)$ ligands are presented. Moreover, different from the one set of proton signals (Figure S2) of the (acac) in the *C*₂-symmetric $[\text{Ir}(C^{\wedge}N^1)_2(O^{\wedge}O)]/[\text{Ir}(C^{\wedge}N^2)_2(O^{\wedge}O)]$, the H atoms on the two -CH₃ groups in every (acac) ancillary ligand were split into two sets of singlet signals, peaking at 1.86 and 1.76 ppm for **1**, 1.74 and 1.68 ppm for **2** or 1.73 and 1.61 ppm for **3**, respectively, indicating the desirable destruction of structural symmetry.[15-19] The *C*1-symmetric $[\text{Ir}(C^N)^1](C^N)^2$ (O[^]O)]-*tris*-heteroleptic character of the representative Ir(III)-complex 2 was further confirmed by its X-ray crystallographic analysis (**Tables S1-2**). As depicted in **Figure 1**, one (**iqbt**) ligand, one (**ppy**) ligand with the similar C[^]N chelation (C11[^]N1 or C18[^]N2) mode and one (acac) ancillary ligand with the O^O-chelate (O1^O2) mode coordinate to one Ir(III)-centre in a distorted octahedral geometry. Evidently, besides the O^O-chelation of the (acac)⁻, the *cis*-C,C and *trans*-N,N dispositions of the two different C[^]N ligands ((*iqbt*)⁻ and (ppy)⁻) render the mononuclear Ir(III)-complex 2 comparable to the typical asymmetric [Ir(C^N¹)(C^N²)(O^O)]-tris-heteroleptic Ir(III)-complexes.^[15] TGA (thermogravimetric

analysis) result (**Figure S3**) of the Ir(III)-complexes **1**-**3** showed their good thermal stability with 5% weight-reduction temperature $(\Delta T_{5\%})$ over 300 °C, which is sufficient for the following solution-processed device fabrication.

Figure 1. Perspective drawing of the C_1 -symmetric $[\text{Ir}(C^N N^1)(C^N N^2)(O^N O)]$ -tris-heteroleptic framework in the Ir(III)-complex **2** using thermal ellipsoids at a 30% probability level. All the hydrogen atoms and most of the labels are omitted for clarity.

2.2. Photophysical Properties of the *C***1-Symmetric Ir(III)-Complexes 1-3**

The photophysical properties of the Ir(III)-complexes **1**-**3** measured in solution at room temperature were examined, and the results were summarized in **Table S3**. In contrast to the limited absorptions (λ_{ab} < 450 nm; **Figure S4**) of the C^{\wedge}N ligands and **Hacac**, the Ir(III)complexes **1**-**3** exhibit the significantly broadened UV-visible-NIR absorption spectra (**Figure** 2), and besides the strong $\frac{1}{\pi} \pi^*$ absorption bands at the high-energy region (250-460 nm), the relatively weak absorptions in the lower-energy region $(\lambda_{ab} > 470 \text{ nm})$ are tentatively assigned to the admixed 1,3 LC/^{1,3}MLCT (LC including LLCT (ligand-to-ligand charge transfer) and ILCT (intraligand charge transfer)) transitions. Nonetheless, the differential ^{1,3}LC/^{1,3}MLCT absorption behaviors (λ_{ab} = 508 nm for **1**, 525 nm for **2** or 603 nm for **3**) of the Ir(III)-

complexes **1**-**3**, as confirmed by their DFT (density functional theory) and time-dependent DFT (TD-DFT) calculations (*vide infra*), should result from the different electronic effects with specific $(C^{\wedge}N^2)$ ligands involved.

Figure 2. The normalized UV-visible-NIR absorption and emission spectra of the Ir(III) complexes **1-3** in solution ($\lambda_{\text{ex}} = 477$ nm) and the PVK-OXD7 (65:30; weight ratio) in solidstate film $(\lambda_{ex} = 273 \text{ nm})$ at room temperature, respectively.

Upon photo-excitation, in contrast to the rigidochromic fluorescence (**Figure S5**) of the ligands, all the three C_1 -symmetric Ir(III)-complexes in CH_2Cl_2 solution emit the typical NIR emissions (**Figure 3**), peaking at 703 nm with a weaker shoulder at 779 nm for **1**, 715 nm and 768 nm (sh) for **2**, and 707 nm and 756 nm (sh) for **3**, respectively. The absence of the ligands' emissions should result from the efficient Dexter energy transfer^[24] for the Ir(III)complexes **1**-**3**. Meanwhile, their NIR-phosphorescent nature is further reflected from the corresponding T₁-decayed lifetime *τ* of 0.30 *μ*s (1), 0.42 *μ*s (2) or 0.45 *μ*s (3). On one hand, the well-defined vibronic structure of their spectral profiles implies that the 3 LC-state should be substantially involved in the T_1 radiative decay for these NIR-emissive Ir(III)-complexes in addition to the ³MLCT state. The almost no solvatochromic shift (**Figure S6**) of their NIR-

emissive bands in solvents of different polarities, together with the featureless FWHM (full width at half-maximum) up to *ca*. 70 nm, demonstrates their weak geometry distortion at the T_1 state. Moreover, in accordance with the calculated T_1 energy level order above, the NIRemission maximum (λ_{em} = 715 nm) of the (ppy)⁻based Ir(III)-complex 2 exhibits a significant bathochromic shift compared with those $(696-710 \text{ nm})$ of the previous C_2 -symmetric $[Irr(iqbt)]$ -counterparts^[9e,10c] or the other two *C*₁-symmetric Ir(III)-complexes 1 and 3 with the electron-withdrawing/enriching effects. As to the shortest lifetime unusually found for the (**dFppy**)⁻based Ir(III)-complex 1, it may be due to the weaker vibrational nonradiative coupling effect with the relatively smaller Huang-Rhys factor $(S_M)^{[25]}$ of 0.29 than those (0.32-0.38) for the other Ir(III)-complexes **2**-**3**. Noticeably, contributing from the weak vibronic coupling $(S_M = 0.29{\text -}0.38)$ between the T_1 and the S_0 states coupled with the distinctively shorter lifetimes ($\tau = 0.30$ -0.45 μ s) of the *C*₁-symmetric [Ir(C^N¹)(C^N²)(O^O)]*tris*-heteroleptic Ir(III)-complexes **1-3** than those ($\tau \sim 1.0 \mu s$) of conventionally *C*₂-symmetric $[Ir(C^N)/2(L^N)]$ -counterparts,^[9-11] their NIR-emissive efficiencies are significantly improved, exhibiting the $\Phi_{PL} = 0.18$ for **1** ($\lambda_{em} = 703$ nm), 0.26 for **2** ($\lambda_{em} = 715$ nm) and 0.28 for **3** ($\lambda_{em} =$ 707 nm), respectively. In good agreement with these results, their efficient NIRphosphorescence is further confirmed from their six-fold increase in k_r values (6.00-6.22×10⁵) s⁻¹). Excitingly, despite the fact that the situation is restricted to the so-called "energy gap law",[7] the augmented ³MLCT-characters (21.17% (**2**); 20.76% (**3**)) of the Ir(III)-complexes **2**-**3** relative to that of 18.30% (**1**) are reflected, as also confirmed by their DFT/DT-DFT calculations (*vide infra*). Therefore, both the large k_r value and the relatively smaller k_{nr} $(1.76\times10^6 \text{ s}^{-1})$; also **Table S3**), engender the Ir(III)-complex **2** an attractive efficiency (Φ_{PL} = 0.26) even with the NIR-emission wavelength extending to 715 nm.

Figure 3. The normalized NIR emission spectra of the *C*1-symmetric Ir(III)-complexes **1-3** in $CH₂Cl₂$ solution at room temperature.

2.3. Electronic Structure Calculations of the *C***1-Symmetric Ir(III)-Complexes 1-3**

To deeply understand the photophysical behaviors of the *C*1-symmetric Ir(III)-complexes **1-3**, DFT/TD-DFT calculations based on their optimized S_0 and T_1 geometries were performed, respectively. According to the theoretical calculation results (**Tables S4-5** and **Figure 4**), the Ir(III)-complexes **1**-**2** have the similar distribution patterns to all the HOMOs and the LUMOs, where the most contribution (*ca*. 94%) from the C^{\wedge}N² ligand ((**dFppy**)⁻ or (**ppy**)⁻) to the LUMO+1 and the almost entire contribution (*ca*. 93%) from the $C^N N^1$ ligand (**iqbt**) to the LUMO, are observed, respectively. Moreover, besides the negligible (2.07% or 2.31%) contribution from the (acac) ancillary ligand, the HOMO is mainly (63.88% or 56.42%) localized at the $C^N N^1$ ligand ($iqbt$) accompanied by some contributions from the Ir(III)-centre $(25.06\% \text{ or } 28.31\%)$ and the C^{\wedge}N² ligand (9.00% for (dFppy) or 12.95% for (ppy)). Due to the electron-withdrawing effect, the (**dFppy**)⁻-incorporation in the Ir(III)-complex 1 induces more contribution from the $C^{\wedge}N^1$ ligand ($iqbt$) but with significantly reduced contribution

from the $(dFppy)$ ⁻-C^{\wedge}N² ligand, the Ir(III)-centre and the (**acac**)⁻ ancillary to the HOMO as compared with the (ppy)-based Ir(III)-complex 2. However, the replacement of the C^N ² ligand does not change the HOMO-1 pattern between the complexes **1**-**2**, revealing the domination (*ca*. 50%) from the (**acac**) - ancillary ligand, the substantial proportions from the Ir(III)-center (*ca*. 34%) and the C^{\wedge N¹ ligand (**iqbt**)⁻ (*ca*. 10%) and the minor fraction (*ca*. 4%)} from the C^{\wedge}N² species. Therefore, considering the HOMO \rightarrow LUMO transition (*ca*. 94%) featuring the S_1 state (also **Table S4**) of the Ir(III)-complexes 1-2, although their low-energy UV-visible absorption is mainly caused by the 1 LC transition (1 ILCT (48.33%) and 1 LLCT (17.50%) for **1**; ¹ILCT (41.13%) and ¹LLCT (22.23%) for **2**) mixed with the ¹MLCT one (34.17% for **1**; 36.64% for **2**), the electron-withdrawing (**dFppy**) - -incorporation stabilizes both the HOMO (-4.89 eV) and the LUMO (-1.89 eV) of its Ir(III)-complex **1** in comparison to those (-4.73 and -1.79 eV) of the Ir(III)-complex **2**, leading to the slightly broader HOMO-LUMO gap ($E_g = 3.00$ eV (1) versus 2.94 eV (2)), and thereby the blue-shifted ($\lambda_{ab} = 508$ nm (1) versus 525 nm (2)) low-energy absorption. Using the (d pqx) with the larger π -conjugation as the C^{\wedge}N² ligand, its *C*₁-symmetric [Ir(C \wedge N¹)(C \wedge N²)(O \wedge O)]-*tris*-heteroleptic Ir(III)-complex **3** possesses a quite different HOMO/LUMO pattern from those of the Ir(III)-complexes **1**-**2**, in which besides the slightly increased contribution from the $(\mathbf{d}p\mathbf{q}x)$ ⁻-C^{\wedge}N² ligand to the HOMO (14.12%) or the HOMO-1 (6.58%), the LUMO is almost entirely (94.58%) localized on the (d pqx) - C^{\wedge}N² ligand with the most contribution (93.47%) from the C \wedge N¹ ligand (i qbt) to the LUMO+1. Hence, for the Ir(III)-complex **3**, the low-energy UV-visible absorption showing the narrowest HOMO-LUMO gap (2.67 eV) with the distinctive bathochromic shift $(\lambda_{ab} = 603 \text{ nm})$ is differentially originated from the ¹LLCT-dominated state ((C^N²-ligand \rightarrow $C^{\wedge}N^1$ -ligand; 56.37%) among the 1 LLCT/¹MLCT/¹ILCT-admixture. Interestingly, despite the large contribution (*ca*. 89%) from the HOMO \rightarrow LUMO transition to the corresponding T₁ state (1.775 eV (698 nm) of **1**; 1.748 eV (709 nm) of **2**) for the Ir(III)-complexes **1**-**2** while the HOMO \rightarrow LUMO+1 transition to that (1.754 eV (707 nm)) of the Ir(III)-complex 3, their

 $S_0 \rightarrow T_1$ absorption transitions (also **Table S5**) are consistent and reasonably assigned to the admixed ${}^{3}LC/{}^{3}MLCT$ -transitions with the ${}^{3}ILCT$ -domination.

Figure 4. The HOMO and LUMO patterns for the *C*1-symmetric Ir(III)-complexes **1**-**3** based on their corresponding optimized S_0 geometries.

NTO (natural transition orbital) calculations of the three *C*1-symmetric Ir(III)-complexes **1**- **3** on optimized T₁ geometries were carried out to further elucidate their NIR-emissive behaviors. The NTO analysis (**Table S6** and **Figure 5**) shows that both the hole (*ca*. 84%) and the particle (*ca*. 92%) orbitals are dominated by the $C^{\wedge}N^1$ (**iqbt**) ligand for each of the C_1 symmetric Ir(III)-complexes 1-3, and the introduction of the (ppy) -C^{\wedge}N² ligand gives the significantly increased percentage on the d_{π} orbital of the Ir(III) centre for the Ir(III)-complex 2. Nevertheless, based on the entire (100%) hole \rightarrow particle transition, the ³ILCT-dominated (78.10% for **1**; 73.82% for **2**; 72.79% for **3**; also **Table S5**) and the less prevalent ³MLCT (18.30% for **1**; 21.17% for **2**; 20.76% for **3**) transition should be responsible for their wellstructured NIR-emissive phosphorescence.

Figure 5. The NTO (natural transition orbital) hole and particle patterns for the *C*1-symmetric Ir(III)-complexes $1-3$ based on their optimized T_1 geometries.

14 **NTO** Further considering the transition dipole effect from the attractive *C*1-symmetry of the Ir(III)-complexes **1-3**, both the $\vec{\mu}$ -S₀ *and* $\vec{\mu}$ -T₁ dipole moments of the *C*₁-symmetric Ir(III)complexes 1-3 and the *C*₃-symmetric *fac*-[Ir(iqbt)₃]^[26] or the *C*₂-symmetric [Ir(iqbt)₂(dpm)]^[9e] with the record-high $\eta_{\text{EQE}}^{\text{Max}} = 3.07\%$ ($\lambda_{\text{em}} = 714$ nm) were explored for comparison upon the DFT calculations from their corresponding optimized T_1 and S_0 states, respectively. As summarized in **Table S7** and **Figures 6** and **S7**, significantly different from the directions of the $\vec{\mu}$ -S₀ and $\vec{\mu}$ -T₁ dipole moment vectors located at the *C*₂-axis of the *C*₂-symmetric $[\text{Ir}(i\text{qbt})_2(\text{dpm})]^{[9e]}$, the angle $(\theta; 32.44^{\circ} (1), 34.83^{\circ} (2)$ or 114.83° (3) between the two vectors is found for each of the C_1 -symmetric Ir(III)-complexes **1-3** and larger than that (23.17°) of the C_3 -symmetric fac -[Ir(iqbt)₃]^[26]. Accordingly, their transition dipole moment (TDM, $\Delta \vec{\mu}$ (T₁-S₀); 2.20 D (1), 2.36 D (2) and 3.60 D (3)) values are larger than that (2.01 D) of the C_3 -symmetric-*fac*-[Ir(iqbt)₃]^[26] and 4-5 times larger than that (0.56 D) of the C_2 -symmetric- $[Irr(iabt)_{2}(dpm)]^{[9e]}$. However, despite the transition dipole orientation available for the Ir(III)complexes **1**-**3** with the *C*1-symmetry, the hypothesis of their anisotropically oriented EMLs needs to be further verified (*vide infra*).

Figure 6 Chemical structure and TDMs (transition dipole moments) of the *C*₁-symmetric Ir(III)-complexes **1**-**3** by DFT calculations.

2.4. Electrochemical Properties of the *C***1-Symmetric Ir(III)-Complexes 1-3 with Different C^N² -Ligands**

Based on the electrochemical behaviours (**Figure 7**), all the *C*1-symmetric [Ir(C^N¹)(C^N²)(O^O)]-*tris*-heteroleptic Ir(III)-complexes **1**-**3** exhibit a reversible oxidation wave at 0.48 eV (**1**), 0.37 eV (**2**) or 0.46 eV (**3**), respectively. The redox process (HOMO level stabilization; also shown in **Table S4** and **Figure 4**) can be mainly attributed to the oxidation of the benzothiophene unit of the $C^N N^1$ ligand (**Hiqbt**) besides the Ir(III) centre involved. Clearly, in comparison to the Ir(III)-complex **2** or **3**, the stronger anodic shift (0.48 eV) in oxidation of the Ir(III)-complex **1** should result from the electron-withdrawing effect of two F atoms on the dFppy ligand. Meanwhile, owing to the (ppy) -C^{\wedge}N²-induced stronger ³MLCT effect, the higher electron density to the Ir(III) centre of the Ir(III)-complex 2 should be the reason to its lower oxidation potential (0.37 eV) than those (0.46-0.48 eV) of the Ir(III)-complexes **1** and **3**. Accordingly, both the HOMO (-4.77 eV) and LUMO (-2.63 eV) of the Ir(III)-complex **2** are destabilized in comparison with those (-4.88 and -2.65 eV for **1**; - 4.86 and -2.98 eV for **3**) of the other two Ir(III)-complexes, through which a much narrower

experimental HOMO-LUMO band-gap (1.88 eV) of the Ir(III)-complex **3** well verifies its redshifted low-energy absorption.

Figure 7. The cyclic voltammogram (CV) results of the Ir(III)-complexes **1-3** recorded versus Fc⁺/Fc in solution at room temperature under a N₂ atmosphere (scan rate = 100 mV/s).

2.5. Device Performance of NIR-OLEDs Using the *C***1-Symmetric Ir(III)-Complexes 1-3 with Different C^N² -Ligands as the Dopants**

Since the practical LUMO (-2.98 to -2.63 eV) and HOMO (-4.88 to -4.77 eV) energy levels of the C_1 -symmetric $[\text{Ir}(C^N N^1)(C^N N^2)(O^N O)]$ -tris-heteroleptic $\text{Ir}(III)$ -complexes **1-3** as the dopants fit perfectly within those $(-3.00 \text{ to } -2.00 \text{ eV})$ of LUMO and $(-6.30 \text{ to } -5.50 \text{ eV})$ of HOMO) of the PVK-OXD7 (PVK = poly(*N*-vinylcarbazole); OXD7 = 1,3-bis(5-(4-*tert*-butylphenyl)-1,3,4-oxadiazol-2-yl)benzene), the PVK-OXD7 with good carrier transport should be the suitable *co*-host for their solution-processed NIR-OLEDs. Moreover, considering the significant spectral overlap (also **Figure 2**) between the blue-light ($\lambda_{em} = 414$ nm) emission of the *co*-host and the LLCT/MLCT absorption of the dopants, effective Förster energy transfer^[27] should be motivated from the PVK-OXD7 to the Ir(III)-complex species upon the electrical excitation. Through the solution-processed procedure with the same configuration depicted in **Figure 8(a)**, a series of NIR-OLEDs using the Ir(III)-complexes **1**-**3** as the dopants at a 5 wt% doping level in PVK-OXD7 (65:30, wt%) were fabricated. PEDOT:PSS (poly(ethylenedioxythiophene):poly(styrenesulfonate)) acted as the hole-injecting material.

Figure 8. a) Device structures and energy level diagrams; b) normalized electroluminescent spectra; c) $J-V$ (current density versus voltage); d) $R-V$ (radiance versus voltage) and e) $\eta_{EOE} - J$ (external quantum efficiency versus current density) curves based on NIR-OLEDs from the *C*1-symmetric Ir(III)-complexes **1**-**3**.

In order to verify the effective charge transport of the *co*-host PVK-OXD7, hole-only and electron-only devices were fabricated with the configurations of $ITO/MoO₃$ (3 nm)/ PEDOT:PSS (40 nm)/PVK-OXD7 (65:30; weight ratio) (50 nm)/MoO₃ (3 nm)/Al (100 nm) and ITO/LiF (3 nm)/OXD7 (50 nm)/LiF (3 nm)/Al (100 nm), respectively. Meanwhile, based on the *J–V* curves (**Figure S8**) of the hole-only and electron-only devices, the hole/electron mobility was checked using the space-charge-limited current (SCLC) method. Accordingly, the hole mobility of PVK-OXD-7 was calculated to be *ca*. 1.62×10^{-6} cm² V⁻¹ s⁻¹, which is slightly larger than the electron mobility $(2.11 \times 10^{-7} \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1})$ of OXD-7. Clearly, towards a desirable charge transport balance, an extra layer of TmPyPB (1,3,5-tri[(3-pyridyl)-phen-3 yl]benzene) serving as the interface-modified function^[28] of both electron-transporting and hole-blocking, is necessary to ensure the good electroluminescent performance of the resultant NIR-OLEDs.

As shown in **Table S8** and **Figure 8(b)**, all the devices, upon illumination, exhibit the NIR electroluminescent spectra that well resembled the corresponding T_1 -based photoluminescence profiles (also **Figure 3**) of the Ir(III)-complexes **1**-**3**, indicating that the Dexter energy transfer with their corresponding T_1 state should be the major process^[24] despite the fact that the co -host to guest Förster energy transfer^[27] occurred. Moreover, under a proper forward bias, the respective electroluminescent spectra are independent of the applied voltages. By further checking the *J-V*, R -*V* and η_{EOE} -*J* curves shown in **Figures 8(c-e)** for the **NIR-OLED-1/2/3**, in contrast to the monotonous increase of the *J* or *R* with increasing the applied bias voltage, the η_{EOE} in each case increases instantly and then decreases steadily throughout the whole illumination. For the **NIR-OLED-1** (λ_{em} = 703 nm), after application of 6.5 V of the V_{on} (the voltage of $R = 5 \text{ W/sr} \cdot \text{cm}^2$), it displays an η_{EQE} ^{Max} of 3.59% with the $R =$ 16 W/sr·cm² at the $J = 1.18$ mA/cm². As to the other NIR-OLEDs based on the Ir(III)complexes 2 (λ_{em} = 715 nm) and 3 (λ_{em} = 705 nm), despite the distinctively lower R^{Max} (201

W/sr·cm² for the **NIR-OLED-2** upon $V_{on} = 7.2$ V and 84 W/sr·cm² upon $V_{on} = 8.1$ V for the **NIR-OLED-3**) than that $(R^{Max} = 1400 \text{ W/sr} \cdot \text{cm}^2)$ of the **NIR-OLED-1**, the η_{EQE}^{Max} values of 5.30% for the **NIR-OLED-2** and 4.24% for the **NIR-OLED-3** are significantly larger than that (3.59%) of the **NIR-OLED-1**, mirroring their trend in the photo-excited Φ_{PL} (0.18 (1), 0.26 (**2**) and 0.28 (**3**)). The slightly higher *V*on of 8.1 V for the **NIR-OLED-3** than those (6.5- 7.2 V) of the other **NIR-OLEDs-1-2**, should be caused by the larger LUMO energy barrier (0.36 eV) between the emitter **3** and TmPyPB. Nonetheless, profiting from the short T_1 decayed lifetimes ($\tau = 0.30$ -0.45 μ s), there is an almost negligible efficiency-roll-off (< 2%) for each of the three NIR-OLEDs.

Considering the record-high $\eta_{\text{EQE}}^{\text{Max}} = 3.07\%$ ($\lambda_{\text{em}} = 714$ nm) for the *C*₂-symmetric $[Ir(iqbt)₂(dpm)]^[9e]$ doping in the PVK-OXD7 while the cathode Ba (7 nm)/Al (100 nm) was used, it is of particular interest on its scientific comparison with the C_1 -symmetric $[\text{Ir}(C^N\text{N}^1)(C^N\text{N}^2)(O^N\text{O})]$ -tris-heteroleptic Ir(III)-complexes **1-3**. Through the same device structure as the **NIR-OLED-1/2/3**, the $[Ir(iqbt)_{2}(dpm)]$ -doped reference device was fabricated, and its electroluminescent behaviour (also **Table S8**) was shown in **Figure S9**. Despite the replacement of TmPyPB $(45 \text{ nm})/\text{LiF}$ $(1 \text{ nm})/\text{Al}$ (100 nm) , the $[\text{Ir}(i\text{qbt})_2(\text{dpm})]$ -doped reference device gave the almost identical η_{EQE} ^{Max} = 3.46% (λ_{em} = 711 nm) to that^[9e] with the cathode Ba (7 nm)/Al (100 nm). However, their inferior electroluminescent performance $(\eta_{\text{EQE}}^{\text{Max}} = 3.07$ -3.46% and ~10% efficiency-roll-off) relative to those $(\eta_{\text{EQE}}^{\text{Max}} = 3.59$ -5.30% and negligible $(< 2\%)$ efficiency-roll-off) of the **NIR-OLED-1/2/3** fabricated from the C_1 symmetric $[\text{Ir}(C^N\text{N}^1)(C^N\text{N}^2)(O^N\text{O})]$ -tris-heteroleptic Ir(III)-complexes **1-3** is evident.

2.6. Effect of the Preferential Horizontal Orientation on the Enhanced Out-Coupling Efficiency of the NIR-OLEDs-1-3

Intriguingly, as compared with those solution-processed devices from the typical *C*3 symmetric-*fac*-[Ir(C^N)₃]^[8] and *C*₂-symmetric [Ir(C^N)₂(L^X)]-counterparts,^[9-11] the significantly improved electroluminescent performance of the **NIR-OLEDs-1-3** fabricated from the C_1 -symmetric $[\text{Ir}(C^N N^1)(C^N N^2)(O^N O)]$ -tris-heteroleptic Ir(III)-complexes **1-3** are highlighted. According to the theoretical equation 1,

$$
\eta_{\text{EQE}} = \eta_{eh} \times \eta_{PL} \times \eta_{\text{exciton}} \times \eta_{\text{out}}
$$
\n(1)

 η_{eh} is the recombination efficiency of injected holes and electrons, η_{PL} is the radiative emitter's quantum yield*, ηexciton* is the radiative exciton ratio and *ηout* is the light out-coupling efficiency, the high efficiency (η_{EQE} ^{Max} = 3.59% of the **NIR-OLED-1**; 5.30% of the **NIR-OLED-2**; 4.24% of the **NIR-OLED-3)** of our cases should inherently benefit from their inherent high quantum yield ($\Phi_{PL} = 0.18$ for **1** ($\lambda_{em} = 703$ nm); 0.26 for **2** ($\lambda_{em} = 715$ nm); 0.28 for **3** (λ_{em} = 707 nm)) associated with the weak vibrational nonradiative coupling effect.^[25] Meanwhile, relying on the suitable $TmPyPB$ -interfaced^[28] HOMO/LUMO-level alignment (also **Figure 8(b)**) towards the enhanced carrier injection/transport, more ${}^{1}T$ -excitons' confinement and recombination should be in a subordinate position. Especially contributing to the TmPyPB-induced negligible electron injection/transport barrier for the **NIR-OLED-2**, the desirable carrier balance might take into effect to give a better performance than that of the **NIR-OLED-3**. More importantly, even assuming the η_{eh} of 100% and the $\eta_{exciton}$ of 100%,^[29] the calculated *ηout* of 20.0% for the **NIR-OLED-1**, 20.4% for the **NIR-OLED-2** or 15.1% for the **NIR-OLED-3** means that molecular orientation of the *C*1-symmetric $[\text{Ir}(C^N\text{N}^1)(C^N\text{N}^2)(O^N\text{O})]$ -tris-heteroleptic Ir(III)-complexes **1-3** in the EMLs should be beneficial for the NIR-light extraction.[20-22]

To further confirm this hypothesis, the quantitative orientation distribution of the EMLs doped with the *C*1-symmetric Ir(III)-complexes **1**-**3** was experimentally determined by variable-angle spectroscopic ellipsometry (VASE) method.^[30] Based on the following equations of the horizontal dipole ratio $h/(h + v)$ and the order parameter *S* relative to the

TDMs' anisotropy, respectively,

$$
\frac{h}{h+v} = \frac{2(1-S)}{3}
$$
 (2)

$$
S = \frac{3(\cos^2 \theta) - 1}{2} = \frac{k_e^{\max} - k_o^{\max}}{k_e^{\max} + 2k_o^{\max}}
$$
 (3)

where *S* (defined by the *e* angle between the the transition dipole moment vector and the direction vertical to the substrate, or the ordinary and extraordinary coefficients k_0 and k_e) of 0, -0.5 or 1 represents the isotropic, completely horizontal or vertical alignment to the substrate, respectively. As shown in **Figure 9**, both the ordinary refractive index (*n*o) and the extraordinary refractive index (n_e) of the EMLs-1-3 are highly sensitive to the wavelength of light, while their large values even at the corresponding emissive wavelength (1.637 and 1.661 at 703 nm (EML-**1**); 1.644 and 1.634 at 715 nm (EML-**2**); 1.683 and 1.648 at 707 nm (EML-**3**)) indicate that each of the EMLs-**1**-**3** should be horizontally orientated.[31] Moreover, their θ angles are up to 56.21-62.54°. Accordingly, their *S* (-0.036 of the EML-1; -0.073 of the EML-2 or -0.181 of the EML-3) and the corresponding $h/(h + v)$ (exceeding 69%; 69.0%, 71.5% or 78.7%) results are summarized in **Table S9**, where the preferential orientation parallel to the substrate^[20-22,30] for all the EMLs-1-3 contributes to the high efficiency of their **NIR-OLEDs-1-3**. Worthy of note, the **NIR-OLED-2**, probably with the best overall optoelectronic properties, represents the most efficient one ($\eta_{EQE}^{Max} = 5.30\%$; $\lambda_{em} = 715$ nm) among other solution-processed NIR-OLEDs^[9a,9c-9(f),9(i),10-11] based on Ir(III)-complexes (**Table S10** and **Figure 10**). Besides the superiorities on cost-effectiveness and large-area, its satisfactory performance (high-efficiency and low efficiency-roll-off; $\lambda_{\text{em}} = 715$ nm), to the best of our knowledge, can even compete with those from vacuum-deposited NIR-OLEDs^[8,9(b),9(g)-9(h)] based on Ir(III)-complexes.

Figure 9. Wavelength-relative refractive indices (*n*; a-c) and coefficients (*k*; d-f) in ordinary (n_0/k_0) and extraordinary modes (n_0/k_0) of the EMLs doped with the corresponding C_1 symmetric Ir(III)-complexes **1**-**3**.

Figure 10. The λ_{EL} -relative η_{EOE} comparison between the **NIR-OLEDs-1/2/3** in this work with previously reported Ir(III)-complex-based vacuum-deposited NIR-OLEDs (NIR-OLEDs-V) and solution-processed small-molecule-hosted NIR-OLEDs (NIR-OLEDs-S) or polymerhosted NIR-OLEDs (NIR-OLEDs-P).

3. Conclusions

In conclusion, through the C_1 -symmetric $[\text{Ir}(C^N N^1)(C^N N^2)(O^N O)]$ -tris-heteroleptic molecular design strategy, the Ir(III)-complexes $1-3$ showing the ${}^{3}LC/{}^{3}MLCT$ -admixed NIRphosphorescence have been successfully obtained. The electronic effects endowed by different (C^{\wedge}N²)-ligands can govern the color-tuning within the NIR (λ _{em} = 703-715 nm) region to some extent. Moreover, as a result of their doped polymer films with the preferential horizontal orientation for reliable solution-processed NIR-OLEDs, the attractive device performance especially with the $\eta_{\text{EQE}}^{\text{Max}}$ ($\lambda_{\text{em}} = 715$ nm) up to 5.30% is the best among solution-processed NIR-OLEDs and even comparable to those of vacuum-deposited NIR-OLEDs fabricated from Ir(III)-complexes with the similar NIR-emissive gamut. This result offers Ir(III)-complex featuring a *C*1-symmetric [Ir(C^N¹)(C^N²)(O^O)]-*tris*-heteroleptic structure a new platform to efficient NIR-OLEDs.

Supporting Information

Electronic Supporting Information including experimental details, synthesis and characterization, and supplemental Tables and Figures is available from the Wiley Online Library or from the author.

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Conflict of Interest

The authors declare no conflict of interest.

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Keywords

Cyclometalated Ir(III)-complex, *tris*-heteroleptic configuration, transition dipole moment; near-infrared phosphorescence, solution-processed organic light-emitting diode

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