This version of the article has been accepted for publication, after peer review (when applicable) and is subject to Springer Nature's AM terms of use(https:// www.springernature.com/gp/open-research/policies/accepted-manuscript-terms), but is not the Version of Record and does not reflect post-acceptance improvements, or any corrections. The Version of Record is available online at: http://dx.doi.org/10.1007/s11661-021-06204-w.

Atomistic Simulation on the Mechanical Properties of Diffusion Bonded Zr-Cu Metallic Glasses with Oxidized Interfaces 3

Tian Li^a, Guangping Zheng^{a,*}

^a Department of Mechanical Engineering, Hong Kong Polytechnic University, Hung Hom,

Kowloon, Hong Kong SAR, 10000, China

* To whom correspondence should be addressed. E-mail: mmzheng@polyu.edu.hk

ABSTRACT

10 A novel welding technique of diffusion bonding for Zr-Cu metallic glass with pre-oxidized surfaces is proposed in this study, which is systematically investigated by molecular dynamics 11 12 (MD) simulation. Compared with the conventional welding technique, the diffusion bonding 13 process can be well implemented below the crystallization temperature of metallic glass. The obtained structure possesses glass-glass interfaces (GGIs) similar with those in nano-glasses. 14 15 As revealed by MD simulation, the diffusion bonded metallic glasses possess enhanced mechanical strength and ductility that generally do not exist in nano-glasses and their bulk 16 17 metallic glass counterparts. The GGIs are found to hinder the propagation of shear bands, where 18 there is strong bonding between Zr and O and the segregated Cu and ZrO₂ clusters could induce 19 extra plasticity. The results demonstrate that the diffusion bonding of metallic glass with pre-20 oxidized surfaces could provide an alternative approach in solving the longstanding issue of size limitation on metallic glasses. 21

22

4

5

6 7

8

9

23 I. INTRODUCTION

24 Amorphous alloys, also known as metallic glasses (MGs), possess outstanding mechanical 25 properties including high mechanical strength and excellent corrosion resistance, and are promising for structural applications [1]-[2]. Even though different MGs with a variety of 26 compositions have been fabricated in the past decades, the MGs are still not large enough in 27 physical dimensions (< 10 cm) due to the fast quenching rate required for suppressing 28 29 crystallization. Meanwhile, conventional welding technique cannot be applied to MGs due to 30 high temperature profiles at heat affected zone (HAZ) that might induce crystallization. The mechanical properties of MG might be greatly deteriorated if low-temperature (< 500 °C) 31 solders are used in the welding process. Therefore, to overcome the size limitation on MGs, 32 33 other welding techniques for MGs should be further investigated and developed. By far, only a 34 few studies have been reported, demonstrating the feasibility of friction, electron-beam, spark, pulse-current and explosive jointing techniques in wedding MGs [3]-[7]. However, those 35 36 alternative approaches are not suitable for massive manufacturing and the processing parameters must be precisely controlled. Moreover, the mechanical performances of MGs 37 welded by those approaches are at best compatible to those of as-prepared MGs. 38

Recently, a novel type of nanomaterial, named as nano-glass, has been successfully prepared, which consists of numerous glass-glass interfaces (GGI) connecting different nanosized MG clusters or particles [8]–[11]. It is mainly synthesized by inert gas condensation, magnetron sputtering, serve plastic deformations and electrodeposition. Among these techniques, inert gas condensation is the most commonly used one, which produces nanoglasses by compacting nano-sized MG clusters under high external pressure (>100 MPa) and

at ambient or elevated temperatures [12]-[13]. Such formation of nano-glasses can be 45 46 considered as diffusion bonding, and could be promising in solving the longstanding issue of size limitation on MGs. Moreover, the ductility of as-prepared nano-glasses is found to be 47 enhanced, compared to those of their bulk MG counterparts [14]–[25]. On the other hand, the 48 49 viability of diffusion bonding for nano-glass fabrication under ambient conditions, e.g., without the protection of inert gases, becomes critical and the formation of GGIs among different MG 50 clusters must be facile to eliminate the manufacturing barrier, for example, through the powder 51 52 metallurgy manufacturing route. Nonetheless, one of the challenging issues of handling nanosized MG clusters or powders in those manufacturing routes is the oxidation inevitably occurred 53 54 on the surfaces of MGs. Thus, it is necessary to understand the effects of oxidized layers of the 55 nano-sized MG clusters on the diffusion bonding of nano-glasses.

In this study, we investigate the diffusion bonding of MG blocks with oxidized surfaces since surface oxidation of nano-sized MG blocks cannot be avoided at ambient processing conditions. From molecular dynamics (MD) simulation, the pre-oxidized surfaces are found to facilitate the formation of GGIs at different diffusion bonding temperatures under atmospheric pressure, firmly connecting the nano-sized MG blocks and resulting in multilayer structures. Furthermore, the mechanical properties of the diffusion bonded MGs are predicted and the size effects of oxidized MG blocks on their mechanical properties are investigated.

63

64 II. SIMULATION METHODOLOGIES

65 The MD simulations were performed using Large-Scale Atomic/Molecular Massively 66 Parallel Simulator (LAMMPS) [26]. A representative crystalline Zr₅₀Cu₅₀ alloy with 67 dimensions of $5 \times 5 \times 5$ nm³, roughly 7000 atoms, was constructed using ATOMSK software [27] and initially relaxed at 300 K for 100 ps under zero external pressure. Periodic boundary 68 conditions (PBCs) were applied and the embedded atom method (EAM) [28] was used to 69 describe the inter-atomic interactions of Zr-Zr, Zr-Cu and Cu-Cu atoms. Then the system was 70 heated gradually and equilibrated at 2000 K (above the melting temperature $T_m = 1400$ K) for 71 2 ns to allow for the complete melting of Zr₅₀Cu₅₀, and the melt was quenched to 300 K at a 72 cooling rate of 10¹² K s⁻¹ and relaxed for 500 ps. The obtained Zr₅₀Cu₅₀ was further analyzed 73 by radial distribution function that confirmed its amorphous structure. Subsequently, a 74 75 replicated MG block was placed at the other side and an additional O₂ layer was inserted 76 between two MG blocks. As shown in Fig. 1(a), the surface of each MG block was exposed to 77 O₂. Oxidation at the exposed MG surfaces under zero external pressure was simulated at 300-78 600 K for 2 ns by using canonical-ensemble (NVT) MD. The Charge-Optimized Many-Body 79 (COMB3) potential [29] was used to describe the inter-atomic interactions of Zr-O, Cu-O and 80 O-O atoms. After the oxidation, the extra O2 molecules were removed and the remained MG blocks with oxidized surfaces were consolidated for 2 ns under a hydrostatic pressure of 1 bar 81 at 400, 500, 600 and 700 K (the glass transition temperature $T_g = 670$ K), separately. After the 82 83 bonding process, the system was equilibrated at 300 K for 1 ns, resulting in the diffusion bonded 84 MG structure with one GGI, as shown in Fig. 1 (b). In order to investigate the size effects on 85 the mechanical properties of the diffusion bonded MGs, MG blocks with oxidized surfaces (Fig. 86 1c) which had roughly 10000, 20000, 30000, 40000 and 60000 atoms and were 2.5, 5, 10, 15, 87 20 nm in thickness, were consolidated at 600 K under a hydrostatic pressure of 1 bar, as 88 simulated by isobaric-isothermal (NPT) ensemble MD. Different from the previously obtained

systems with one GGI only, the systems containing MG blocks with different thicknesses
contained two GGIs, as shown in Fig. 1 (d). In the aforementioned MD simulations of oxidation
and diffusion bonding processes, PBC was applied in each direction and the finally obtained
diffusion bonded MGs formed the multilayer structures, as shown in Fig. 1.



Fig. 1. Schematics of diffusion bonded Zr₅₀Cu₅₀ MGs with (a) one GGI, and (c) two GGIs, whose atomic
 structures are shown in (b) and (d), respectively.

96

93

All diffusion bonded systems were equilibrated at 300 K for 2 ns before the uniaxial tensile loading was applied in *x*-direction at a strain rate of $3.33 \times 10^7 \text{ s}^{-1}$. There were lateral tractions at *y*- and *z*-directions during the tensile loading process, and the temperature was kept at 300 K by using a Nose-Hoover thermostat in NVT ensemble MD simulation. The atomistic mechanisms of diffusion bonding and mechanical deformation were analyzed by self-scripted MATLAB code and visualized by OVITO software packages [30].

103 104

III. RESULTS AND DISCUSSION

105 The diffusion bonded MGs with one and two GGIs are shown in Fig. 1 (b) and (d), respectively, clearly indicating that the GGIs have different atomic structure with the MG 106 blocks. The mechanical properties of diffusion bonded Zr-Cu MGs with one oxidized GGI, 107 which contain 10 nm Zr-Cu MG blocks welded at different temperatures, are investigated by 108 the tensile tests. The engineering stress vs. strain curves, presented in Fig. 2 (a), reveal that the 109 110 Zr-Cu MGs are well bonded together, without showing a sudden stress drop with a strain lower than the elastic limit of Zr-Cu MGs (~2%). After reaching its ultimate tensile strength (UTS) at 111 1.52 -1.75 GPa, the engineering stress starts decreasing until the fracture occurs. It is found that 112 the welded Zr-Cu MGs with oxidized GGI could have enhanced UTS and ductility as measured 113 by the plastic strain at the UTS, in comparison with those of pristine MGs (Fig. 2a). Among the 114 systems investigated, Zr-Cu MG bonded at 500 K exhibits the best combination of mechanical 115 strength and ductility, i.e., an extra plasticity of 6.2% and a UTS of 1.75 GPa. 116

117





119 Fig. 2. The mechanical properties of diffusion bonded MGs with MG blocks 10 nm in thickness. (a)

Stress vs strain curves for the MGs bonded at different temperatures. (b) The influences of weldingtemperature on fracture strain and UTS.



122

Fig. 3. The atomic shear strain mapping for MGs bonded at 400, 500, 600 and 700 K with an engineering
strain of 7.5 %. The dash lines mark the GGI regions.

125

To have a better understanding on the welding temperature dependence mechanical properties, the mapping of atomic shear strain for each diffusion bonded MG at an engineering strain of 7.5% is shown in Fig. 3, and the shear bands can be identified as those localized regions with high shear strains. Here, the atomic shear strain is defined as [31-34]:

130
$$\eta_i = \frac{1}{2} (J_i J_i^T - I)$$
, (1)

131
$$\eta_{i}^{\text{Mises}} = \sqrt{\eta_{xy}^{x} + \eta_{xz}^{2} + \eta_{yz}^{2} + \frac{(\eta_{xx} - \eta_{yy})^{2} + (\eta_{xx} - \eta_{zz})^{2} + (\eta_{yy} - \eta_{zz})^{2}}{6}}, \quad (2)$$

132 where η_i is the local Lagrangian strain matrix for each atom *i*, J_i is the local transformation 133 matrix and η_i^{Mises} is the atomic shear strain. It can be observed that in the MGs bonded at 500 134 K, the GGI actively involves in the plastic deformation and the roles of GGI are two folded. 135 The shear bands can be generated from GGI and propagate into the MG blocks, or they are 136 generated inside the MG block and are hindered by the GGI, resulting in the enhanced plasticity and mechanical strength, respectively. The mapping of shear strain in the MG blocks at both sides of GGI, as shown in Fig. 3, indicates that in the MGs bonded at 500 K there are 5 localized shear regions that can be connected with GGI. For other diffusion bonded MGs at different temperatures (400, 600 and 700 K), on the contrary, the GGI is less involved during the plastic deformation and the number of those regions are only 2 to 3. It could be the reason that their mechanical properties are not compatible to those of MGs bonded at 500 K.

To investigate the size effects on the mechanical properties of diffusion bonded Zr-Cu 143 MGs, MG blocks with different sizes in thickness are oxidized and welded at 600 K. The stress 144 145 vs. strain curves for diffusion bonded MGs are analyzed, as shown in Fig. 4 (a). The Young's 146 modulus and tensile strength of diffusion bonded MGs are enhanced as the thickness of MG 147 block decreases, while the elongation to failure of diffusion bonded MGs is deteriorated. When the thickness of MG block is lower than 5 nm, its tensile strength is dramatically increased up 148 149 to 3.7 GPa, as shown in Fig. 4 (b). The tensile strength of diffusion bonded MGs is close to that of pristine MGs when the thickness of MG block is increased to 20 nm. Moreover, the ductility 150 151 is dramatically improved from 4.8% to nearly 10% when the thickness of MG blocks is reduced to 2.5 nm. 152



153

Fig. 4. (a) Stress vs strain curves of the diffusion bonded MGs with MG blocks in different sizes
(thicknesses); L=5 nm. (b) The size effects on UTS. Atomic shear strain mapping of diffusion bonded
MGs with MG blocks 5-nm (left column) and 20-nm in thickness (right column), which are deformed at

157 strains of 2.5% (c), 5% (d), 7.5% (e), 10% (f). The dash lines mark the GGI regions.

158

The mapping of atomic shear strain for the diffusion bonded MGs with 5 and 20 nm MG 159 blocks in thickness is shown in Fig. 4. When the diffusion bonded MGs are under plastic 160 deformation (strain>2%), both systems are deformed heterogeneously with the initiation of 161 shear bands either inside the MG blocks or from the GGIs. At an engineering strain of 5%, the 162 shear bands are found to propagate along the direction with an incline angle of about 45° with 163 the loading direction. In the diffusion bonded MGs with 20 nm MG blocks, most shear bands 164 do not interact with GGIs, and grow across the system in the lateral direction, directly leading 165 166 to the catastrophic failure of the bonded MGs. However, in the diffusion bonded MGs with 5 167 nm MG blocks, the shear bands are hindered by GGIs, and as a result, the mechanical strength is enhanced. Fig. 4 (c)-(f) thus demonstrates that the roles of GGI may be different in the 168 169 diffusion bonded MGs containing MG blocks with different sizes. Therefore, as shown in Fig. 4 (b), the mechanical strength of diffusion bonded MGs is found to significantly increase when 170 the ratio of thickness to lateral dimension of the MG blocks is less than 1.0. 171





Fig. 5. The chemical composition analysis on oxidized interfaces in the bonded Zr-Cu MGs with L=5 nm: (a) atomic count variation vs position, (b) visualization of atomic volume, (c) atomic volume at the

interface bonded at different welding temperatures, and (d) density at the interface bonded at differenttemperatures. The dash lines mark the GGI regions.

177

Since the oxidized interface plays an important role in low-temperature welding and in 178 enhancing mechanical properties of the diffusion bonded MGs, it is necessary to conduct a 179 comprehensive investigation about the structure and composition of oxidized interface for a 180 better understanding about the mechanisms of welding and the enhanced mechanical properties 181 of bonded MGs. In Fig. 5 (a), the composition variation is plotted with respect to the position 182 in the system. It can be seen from the plots that the atoms are evenly distributed in the MG 183 184 blocks and the composition variation can be observed at the positions close to the oxidized 185 interface. In general, the oxidized interface could be divided into two different regions that consist of different chemical compositions. The Cu atoms are highly concentrated in the center 186 187 of oxidized interface, where Zr and O atoms are more preferentially located on the MG blocks at both sides of the interface. Compared with the MG blocks, the oxidized interface exhibits a 188 reduced atomic volume, as shown in Fig. 5 (b). However, as the welding temperature is 189 increased, the atomic volume at oxidized interface is further decreased while the density at the 190 191 interface is enhanced, as shown in the Fig. 5 (c) and (d). Moreover, the density in the center of 192 oxidized interface is determined to be the minimum, less than that of MG blocks. In contrast, the density at both sides of the oxidized interface exhibits the maximum, where Zr and O atoms 193 194 are enriched. By analyzing the atomic structure, quite a few ZrO₂ clusters, as shown in Fig. 6 (a), have been found at the oxidized interface. Moreover, as shown in Fig. 6(b), a small amount 195 196 of CuO clusters can be observed at the oxidized interface where Cu atoms seem to be less 197 favorable to interact with O atoms as compared with Zr atoms. The formation of ZrO₂ clusters is critical to the enhanced mechanical properties of the diffusion bonded MGs since the 198 mechanical strength of ZrO_2 crystal is as high as 2000 MPa [35]. On the one hand, the ultimate 199 tensile strength of diffusion bonded MGs could be much enhanced since the ZrO₂ clusters could 200 be hard obstacles that impede the propagation of shear bands generated in the interiors of MG 201 blocks. On the other hand, the formation of ZrO_2 clusters leads to the structural heterogeneity 202 203 around the GGI, where shear bands could be generated close to the resultant ZrO_2 clusters. 204 Because of the sparse distribution of ZrO_2 clusters, multiplication of shear bands at the ZrO_2 205 clusters could lead to the enhanced ductility of the diffusion bonded MGs.



206

207 Fig. 6. Analyses on the chemical composition of (a) Zr-O, and (b) Cu-O clusters formed at glass-glass

208 interfaces in the MGs bonded at different temperatures.

209



Fig. 7. (a) The coordination analysis on Zr-, and Cu-centered clusters in the MGs bonded at different temperatures. (b) Internal energy of MG and diffusion bonded MGs. (c) Absolute energy of the glassglass interface.

214

210

215 The atomic coordination of Zr and Cu atoms at the MG surface oxidized under different 216 temperatures are analyzed, as shown in Fig. 7 (a). For all oxidization temperatures, the O atoms, 217 with an approximated average amount of 4 for Zr-centered atom and 2-3 for Cu-centered atom, are more preferential to interact with the Zr atoms. Meanwhile, Cu-centered atom tends to 218 closely interact with other Cu atoms, leading to the formation of Cu clusters during the 219 220 oxidation process. Therefore, the majority of Cu clusters, concentrated in the center of oxidized 221 interface, could be also responsible for the enhanced ductility of the diffusion bonded MGs. In 222 other words, during the welding process, those Cu clusters play an important role, as low-223 temperature solder, firmly bonding the MG blocks together while providing high ductility of 224 the diffusion bonded MGs under plastic deformation. The oxidized interface is found to be 225 stable, as shown in Fig. 7 (b). As the welding temperature is increased, the internal energy of 226 MGs increases and becomes less stable. In contrast, the internal energy of the diffusion bonded MGs decreases at higher welding temperature. Such a phenomenon is obviously related with the oxidized interface or GGI and indicates a much stable structure of oxidized interface. In addition, the absolute average GGI energy is positively correlated with the welding temperature, which implies an excellent stability of oxidized interface at high welding temperatures, as shown in Fig. 7 (c).



Fig. 8. Diffusion analysis on O atoms during the oxidization at the MG interface at 500 K by applying
(a) Fick's second law of diffusion (at t=2 ns), and (b) Einstein relation. The inset in (a) is the fits of
diffusion coefficient D by using Arrhenius equation.

236

232

In order to explain the heterogeneous atomic structure and position-dependent chemical composition around the GGIs, the analysis on diffusions of Zr, Cu and O atoms during oxidation processes is conducted. At the surface of MG under oxidization, Fick's second law or nonsteady state diffusion could be used to analyze the diffusion coefficients of Zr, Cu and O atoms and the relation is expressed in the following equation

242
$$\frac{c_x - c_0}{c_s - c_0} = 1 - \operatorname{erf}\left(\frac{x}{2\sqrt{Dt}}\right) \qquad , \tag{3}$$

243 where c_x is the concentration at position x, c_s is the surface concentration of diffusing atoms, c_0 244 is the uniform concentration of diffusing atoms away from the surface, D is the diffusion 245 coefficient and t is time. For the diffusion coefficient of each atom type inside the MG block, 246 Einstein relation is applied, and the corresponding equation could be expressed as

247 MSD = 6 D t + C, (4) 248 where MSD is the mean square displacement for each atom type and C is the constant. Finally, 249 the activation energy of diffusion for each element is determined by Arrhenius relation, as 250 follows:

251
$$D = D_0 e^{\frac{E_a}{RT}}$$
, (5)

where E_a is the activation energy, R is the universal gas constant, T is the temperature in Kelvin and D_0 is the maximal diffusion coefficient. Diffusion analysis at the interface under oxidization is performed in every 100 ps at different oxidization temperatures. Meanwhile, the diffusion coefficient inside the MG block is obtained when mean square displacement and its corresponding time follow a linear relationship. Taking the diffusions of O atoms as examples, Fig. 8(a) and (b) shows the determination on the diffusion coefficients D at the surface and in the interior of MG blocks during oxidation processes, using Eqs. (3) and (4), respectively. The

activation energy of diffusion of O atoms is determined by Eq. (5), as shown in the inset of Fig.
8(a).

261

262 Table I. The diffusion coefficients and activation energies of Zr, Cu and O atoms at the surface and in 263 the interior of MG blocks during oxidization

Element Type		Diffusion Coefficient (m ² s ⁻¹)			Activation
		300 K	500 K	700 K	Energy E _a (kJ mol ⁻¹)
Surface	Zr	5.71 x 10 ⁻¹³	1.13 x 10 ⁻¹²	2.20 x 10 ⁻¹²	5.88
of MG	Cu	7.29 x 10 ⁻¹²	8.32 x 10 ⁻¹²	7.56 x 10 ⁻¹¹	11.8
block	Ο	5.32 x 10 ⁻¹²	6.93 x 10 ⁻¹²	2.44 x 10 ⁻¹¹	8.24
Interior of MG block	Zr	4.03 x 10 ⁻¹³	2.02 x 10 ⁻¹²	5.33 x 10 ⁻¹²	11.1
	Cu	4.27 x 10 ⁻¹³	2.73 x 10 ⁻¹²	9.32 x 10 ⁻¹²	13.3
	Ο	7.60 x 10 ⁻¹²	4.38 x 10 ⁻¹¹	7.09 x 10 ⁻¹¹	10.0

264

265 As listed in Table I, the diffusion coefficient of Zr atoms at the surface under oxidization is much lower than those of Cu atoms and O atoms. Consequently, the Cu atoms could diffuse 266 to and aggregate at the outer layer of the oxidized surface of the MG block, while O atoms 267 could diffuse into the interiors of MG block and interact with Zr atoms to form the ZrO₂ clusters, 268 as shown in Fig. 7(a). Thus, the fact that after the diffusion bonding processes there is broken 269 270 concentration homogeneity of Cu and Zr atoms in the GGIs, as shown in Fig. 5, can be well 271 explained. In the diffusion bonded MGs, Cu atoms in the center of GGI could have the highest concentration, as shown in Fig. 5(a) and Fig. 7(a), and could serve as low-temperature solders 272 during the welding process. In contrast, Zr atoms tend to form stable ZrO₂ clusters away from 273 the center of GGI because of the fast diffusion of O atoms from the oxidized surface into the 274 275 interior of MG blocks.

276

277 IV. CONCLUSION

278 In summary, a novel low-temperature welding technique for MG is investigated by MD simulation. The diffusion bonding of Zr-Cu MGs with pre-oxidized surfaces is found to result 279 in the enhanced mechanical properties of MGs when the bonding temperature and aspect ratio 280 of the MG blocks are properly chosen. The GGIs in the diffusion bonded MGs are characterized 281 282 by the presence of segregated Cu and ZrO₂ clusters, mainly resulting from the fast diffusions 283 of Cu atom to the center of interface and the diffusion of oxygen atoms into the interior of MG blocks. The effects of GGIs on the plastic deformation of diffusion bonded MGs are 284 investigated. This study paves a new avenue to overcome the longstanding issue of size 285 limitation on MGs. 286

287

288 Acknowledgements

This work was supported by a grant from the Research Grants Council of the Hong Kong
Special Administrative Region, China (PolyU152607/16E).

291

292	2 References				
293	1	Y. Shen, X.C. Zheng, and G.P. Zheng, Metall. Mater. Trans. A, 2011, vol. 42, pp. 211-			
294		8.			
295	2	T. Burgess and M. Ferry, Mater. Today, 2009, vol. 12, pp. 24–32.			
296	3	Y. Kawamura, T. Shoji, and Y. Ohno, J. Non. Cryst. Solids, 2003, vol. 317, pp. 152-7.			
297	4	J. Kim and Y. Kawamura, Scr. Mater., 2007, vol. 56, pp. 709–12.			
298	5	C.H. Wong and C.H. Shek, Scr. Mater., 2003, vol. 49, pp. 393-7.			
299	6	Y. Kawamura and Y. Ohno, Scr. Mater., 2001, vol. 45, pp. 279-85.			
300	7	Y. Kawamura and Y. Ohno, Scr. Mater., 2001, vol. 45, pp. 127–32.			
301	8	H. Gleiter, Beilstein J. Nanotechnol., 2013, vol. 4, pp. 517-33.			
302	9	H. Gleiter, T. Schimmel, and H. Hahn, Nano Today, 2014, 9, vol. 9.			
303	10	P. Zhou, Q. Li, P. Gong, X. Wang, and M. Zhang, Microelectron. Eng, 2020, vol. 229,			
304		рр. 1113-63.			
305	11	C. Guo, Y. Fang, B. Wu, S. Lan, G. Peng, X.L. Wang, H. Hahn, H. Gleiter, and T. Feng,			
306		Mater. Res. Lett., 2017, vol. 5, pp. 293–9.			
307	12	O. Adjaoud and K. Albe, Acta Mater., 2018, vol. 145, pp. 322-30.			
308	13	O. Adjaoud and K. Albe, Acta Mater., 2016, vol. 113, pp. 284–92.			
309	14	D. Şopu, Y. Ritter, H. Gleiter, and K. Albe, Phys. Rev. B, 2011, vol. 83, pp. 1-4.			
310	15	B. Cheng and J.R. Trelewicz, J. Mater. Res., 2019, vol. 34, pp. 2325-36.			
311	16	O. Adjaoud and K. Albe, Acta Mater., 2019, vol. 168, pp. 393-400.			
312	17	D. Şopu and K. Albe, Beilstein J. Nanotechnol., 2015, vol. 6, pp. 537-45.			
313	18	Z.D. Sha, P.S. Branicio, Q.X. Pei, Z.S. Liu, H.P. Lee, T.E. Tay, and T.J. Wang,			
314		Nanoscale, 2015, vol. 7, pp. 17404–9.			
315	19	K. Albe, Y. Ritter, and D. Şopu, Mech. Mater., 2013, vol. 67, pp. 94–103.			
316	20	S. Adibi, P.S. Branicio, Y.W. Zhang, and S.P. Joshi, J. Appl. Phys, 2014. vol. 116, pp.			
317		043522-1.			
318	21	Y. Ritter, D. Opu, H. Gleiter, and K. Albe, Acta Mater., 2011, vol. 59, pp. 6588–93.			
319	22	S. Adibi, P.S. Branicio, and S.P. Joshi, Sci. Rep., 2015, vol. 5, pp. 1–9.			
320	23	Y. Zhao, X. Peng, C. Huang, B. Yang, N. Hu, and M. Wang, 2019, vol. 9, pp. 1–17.			
321	24	S.H. Nandam, O. Adjaoud, R. Schwaiger, Y. Ivanisenko, M.R. Chellali, D. Wang, K.			
322		Albe, and H. Hahn, Acta Mater., 2020, vol. 193, pp. 252-60.			
323	25	M. Zhang, Q.M. Li, J.C. Zhang, G.P. Zheng, and X.Y. Wang, J. Alloys Compd., 2019,			
324		vol. 801, pp. 318–26.			
325	26	S. Plimpton, J. Comput. Phys., 1997, vol. 117, pp. 1-42.			
326	27	P. Hirel, Comput. Phys. Commun., 2015, vol. 197, pp. 212-9.			
327	28	Y.Q. Cheng and E. Ma, Prog. Mater. Sci., 2011, vol. 56, pp. 379-473.			
328	29	T. Liang, T.R. Shan, Y.T. Cheng, B.D. Devine, M. Noordhoek, Y. Li, Z. Lu, S.R.			
329		Phillpot, and S.B. Sinnott, Mater. Sci. Eng. R Reports, 2013, vol. 74, pp. 255–79.			
330	30	A. Stukowski, Model. Simul. Mater. Sci. Eng., 2010, vol. 18, pp. 015012-7.			
331	31	M.L. Falk and J.S. Langer, Phys. Rev. E, 1998, vol. 57, p. 14.			
332	32	F. Shimizu, S. Ogata, and J. Li, Mater. Trans., 2007, vol. 48, pp. 2923-7.			
333	33	J. Li and F. Shimizu, <i>Report</i> , 2005, vol. 0, pp. 4–6.			
334	34	Y.Q. Cheng, A.J. Cao, and E. Ma, Acta Mater., 2009, vol. 57, pp. 3253-67.			
335	35	C. Piconi, G. Maccauro, Biomaterials, 1999, vol. 20, pp. 1-25.			