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# Bifurcation Analysis and Experimental Study of a Multi-Operating-Mode Photovoltaic-Battery Hybrid Power System

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Abstract—The stand-alone hybrid renewable power generation systems for local applications have gained popularity in recent years. However, due to the intermittent nature of the renewable resources, the hybrid renewable power generation systems are often designed to operate with multiple structures and multiple operating modes. The design for stable operation of such systems requires consideration of the stability conditions for all possible structures and operating modes. As a system of practical importance, the stand-alone photovoltaic-battery hybrid power system is studied for illustrating the possible complex behavior in this paper. We reveal smooth bifurcation in this system, including slow-scale Neimark-Sacker bifurcation, fast-scale period-doubling bifurcation as well as coexisting bifurcation. Under certain conditions, when the system switches its operating mode, a non-smooth bifurcation, manifested as a jump between stable and unstable behavior, can also be observed. Moreover, a detailed analysis based on a discrete-time mapping model is performed to identify these bifurcation phenomena and evaluate the stability boundaries of the system. Extensive experiments verify the theoretical analysis and simulated results.

Index Terms—Bifurcation, fast-scale instability, low-scale instability, photovoltaic-battery hybrid system.

#### I. INTRODUCTION

▼ LOBAL environmental concerns and increased demand **J** for energy consumption, coupled with a steady development of renewable energy technologies, are opening up new opportunities for developing renewable power generation systems [1]. For localized applications in some remote nonelectrified regions, constructing stand-alone power generation systems is often more cost-effective than the conventional deployment of the grid-connected power generation system [2]-[9]. However, as climate, solar radiation, wind speed, etc. vary from time to time, the availability of energy from renewable sources is not always predictable. Thus, a power system that depends entirely upon multiple renewable energy sources can be unreliable. In order to provide continuous electrical power to the load, the use of energy storages in stand-alone renewable power generation systems is usually adopted. The renewable energy sources and energy storages can be connected to a

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common dc bus via a number of power converters [4]–[10]. In some applications, the power sources are close to the load, and multiple-input converters [11] or multiple-port bidirectional converters [2], [12] can be used as interface conversion systems. Regardless of the exact topologies used, the switching circuits and control systems are required to regulate the dc bus voltage and maximize the use of renewable energy, as well as balance the power from all sources. Then, the renewable power conversion unit may operate in a maximum power point tracking (MPPT) mode or off-MPPT mode, and at the same time, the energy storage unit may provide power to the load or store energy from the renewable power sources. The whole system is thus designed to operate with multiple structures and multiple operating modes [3], [5]-[12]. As a result, the dynamics of such a system is rather complex. The closed-loop design has to take into consideration the different modes of operation. It can thus be appreciated that the stability issue is non-trivial as the system assumes different structures in different operating modes under different ranges of parameters.

For most power electronics systems, normal stable operation refers to stable period-1 operation, which is often the expected operating regime for practical applications. Operation states that are not the normal operation are practically regarded as unstable states, and the bifurcation refers to the change from one type of operation to another. As of now, a rich set of nonlinear phenomena have been investigated in power electronic circuits, such as period-doubling bifurcation [13]-[18], Neimark-Sacker bifurcation [19], saddle-node bifurcation [20], border collision [20]-[22] and chaos [14]-[21]. Due to the change of switching structure, border collision as a type of non-smooth bifurcation has been found in many switching systems, which is characterized by a discontinuous 'jump' in the eigenvalues of the Jacobian as some parameters are varied continuously [21]. In such a case, the duty cycle becomes saturated as it is bounded between 0 or 1. In recent years, the nonlinear stability analysis has been extended to more complex conversion systems, such as the high order Cuk converter [23], [24], single-inductor dual-switching dcdc converters [25], multichannel converters [26], [27], and grid-connected power converters [28]–[30]. However, the dynamics of a multi-structure multi-operating-mode system has rarely been formally discussed in the literature [31], [32]. To illustrate the possible complex behavior in this kind of system, a commonly used multi-operating-mode stand-alone photovoltaic-battery hybrid power system (PBHPS) is studied in this paper, which is based on a solar power system and a

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battery connected to an output dc bus via a boost converter and a bidirectional buck-boost converter, respectively.

## II. PHOTOVOLTAIC-BATTERY HYBRID POWER SYSTEM

# A. PBHPS Description

In PBHPS, a solar power system and a battery are connected to an output dc bus via a boost converter and a bidirectional buck-boost converter, respectively. Shown in Fig. 1 is a commonly used practical PBHPS, where  $v_{of} = v_o H_{of}$  and  $v_{batf} =$  $v_{\text{bat}}H_{bf}$ . Here,  $H_{of}$  and  $H_{bf}$  are the sampling coefficients for the output voltage and the battery voltage, respectively. They are proportionality constants implemented by resistor dividers. Also,  $K_s$  is the sampling coefficient for inductor current  $i_{Lb1}$ and is also a proportionality constant. The basic control law in the input power distribution aims to provide power to the load primarily by photovoltaic (PV) panels, and when solar power becomes excessive or insufficient, the battery serves as a power storage or backup power source. In practice, MPPT for the PV power system is implemented with the Perturb-and-Observe algorithm [33], which regulates the reference  $I_{mppt}$ while the light intensity varies, and forces  $i_{Lb1}$  to follow the reference current under the action of the peak current control loop. When the system enters the steady state,  $I_{mppt}$  is tracked to the current reference at the maximum power point. The compensation ramp is applied to increase the stability range of peak current controlled converters.

Depending on the amount of available solar power  $p_{\rm pv}$  and the state of the battery, the PBHPS has three normal operating modes, as shown in Table I, where  $P_o$  is the output power,  $V_{b\rm min}$  and  $V_{b\rm max}$  are the permitted minimum and maximum battery voltages, and  $I_{b\rm max}$  is the charging current limit.

Operating Mode  $M_1$  — If  $p_{pv} < P_o$ , the solar power is controlled by the peak current control loop in Fig. 1(b), with the MPPT algorithm enabled, i.e., the peak current reference  $I_{ref} = I_{mppt}$ . The battery provides complementary power through the bidirectional buck-boost converter operating in boost mode to regulate  $v_o$ . Thus, in this operating mode, diodes  $D_c$ ,  $D_2$  and  $D_3$  are off, and  $D_1$  is on.

If  $p_{pv}$  is larger than  $P_o$  by a small margin, the excessive solar power is used to charge the battery. In this case, the battery charge current is controlled by the energy flow balance, i.e.,  $p_{pv} = P_o + p_{bat}$ . Assuming converters are lossless,  $p_{bat} = v_{bat}i_{bat}$  is the power absorbed by the battery.

Operating Mode  $M_2$  — If more excessive power is available from the PV panels, causing the charge current  $i_{\text{bat}}$  to reach the limit  $I_{b\text{max}}$ , then the bidirectional buck-boost converter operates in buck mode to control charge current, i.e.,  $D_2$  is on and  $D_1$  is off . Meanwhile, as excessive power is delivered to the load,  $v_o$  increases, causing the error signal  $v_e$  to decrease. When it falls below zero, diode  $D_c$  is turned on, and  $v_o$  is regulated through the voltage loop with an inner peak current control loop, i.e.  $I_{\text{ref}} = I_{\text{mppt}} + v_e$ . Under this condition, the MPPT algorithm is disabled as  $v_e$  is smaller than zero and  $I_{\text{ref}}$  is regulated by  $v_e$ .

Operating Mode  $M_3$  — If  $p_{pv} \ge P_o$  and the battery is fully charged, the constant voltage charging (float charging) loop is enabled to prevent self-discharge of the battery, causing  $D_3$  to turn on and  $D_2$  to turn off.



Fig. 1. Photovoltaic-battery hybrid power system. (a) Main topology; (b) control system; (c) typical waveforms for peak current control loop; (d) PI compensator.

TABLE I OPERATING MODES AND CONDITIONS

	$p_{\rm pv} < P_o$	$p_{\rm pv} \ge P_o$	$p_{\rm pv} \ge P_o$
		$i_{\rm bat} < I_{b\rm max}$	$i_{\text{bat}} \geq I_{b\max}$
$v_{\rm bat} \leq V_{bmin}$	shut down	$M_1$	$M_2$
$V_{b\min} < v_{bat} < V_{b\max}$	$M_1$	$M_1$	$M_2$
$v_{\rm bat} \ge V_{b\max}$	$M_1$	$M_3$	$M_3$

$T_s$	$V_{ocpv}$	$r_{\rm pv}$	$I_o$	$L_{b1}$	$L_{b2}$	$C_{\rm pv}$	$C_f$	$C_{\rm bat}$	$V_{ocbat}$	$r_{\rm bat}$	$V_o$	$V_{bmin}$	$V_{bmax}$
$10 \mu s$	20–34V	$50 \mathrm{m}\Omega$	0–5A	$48 \mu H$	$48 \mu H$	$200 \mu F$	$1880 \mu F$	$400 \mu F$	34–40V	$24 \mathrm{m}\Omega$	48V	34V	40V
$k_{p1}$	$ au_{f1}$	$k_{p2}$	$\tau_{f2}$	$k_{p3}$	$ au_{f3}$	$V_{oref}$	$V_{bref}$	$I_{bmax}$	$I_{\rm mppt}$	$H_{of}$	$H_{bf}$	$K_s$	$V_p$
1	0.001–0.6 <i>m</i> s	1	0.01 <i>m</i> s	10	0.012ms	2.5V	2.5V	3A	0–3A	1/19.2	1/16	0.1	0-0.2V

 TABLE II

 PARAMETERS FOR COMPONENT AND CONTROL SYSTEM USED IN SIMULATION.



Fig. 2. Typical simulated waveforms in operating mode  $M_1$  for different values of  $V_{ocpv}$ , with  $V_{ocbat} = 38$  V,  $I_{mppt} = 0.3$  A,  $\tau_{f1} = 0.475$  ms,  $V_p = 0$  V and  $I_o = 3.3$  A are kept constant. Upper trace: actual simulated waveforms, lower trace: sampled-data waveform at the beginning of each switching period. (a) Stable period-1 operation with  $V_{ocpv} = 32$  V; (b) NS bifurcation with  $V_{ocpv} = 26$  V; (c) coexisting fast-scale and slow-scale bifurcation with  $V_{ocpv} = 24$  V.

### B. Bifurcation Behavior from Simulations

Circuit simulations are performed to examine the stability of PBHPS. In our study, the converter parameters are chosen so that CCM is the default operating mode. The switching frequency of the two converters are the same and trailing edge modulation is used for "PWM Comparator2". Firstorder proportional-integral (PI) regulators are employed, as shown in Fig. 1(d), and the transfer function is  $G_{PI}(s) =$  $k_p(1+1/(s\tau_f))$ , where  $k_p = R_2/R_1$  and  $\tau_f = R_2C_1$  are the proportional gain and the integral time constant. The parameters used in the simulations are shown in Table II, where  $V_o$ is the direct component of  $v_o$ ;  $k_{pj}$  and  $\tau_{fj}$  (j = 1, 2, 3) are parameters for the *j*th regulator in Fig. 1(b).

For practical purposes, we examine the dynamic response with  $I_{mppt}$ ,  $V_{ocpv}$ ,  $R_L$ ,  $V_{ocbat}$  and closed-loop parameters serving as the bifurcation parameters, as shown in Figs. 2 and 3. From Fig. 2, the system becomes unstable with slow-scale oscillation when  $V_{ocpv}$  is decreased to 26 V. Here the system operates in  $M_1$ , and the instability can be recognized as a Neimark-Saker (NS) bifurcation. As  $V_{ocpv}$  is further reduced to 24 V, a coexisting fast-scale and slow-scale oscillation will occur [13]. From Fig. 3, we see that the system is stable in the initial operating mode, but it becomes unstable as it switches to another operating mode when the load steps up. Thus, this bifurcation is a non-smooth one. For brevity, we will report the representative smooth and non-smooth bifurcations and will omit the various repetitive details of such bifurcations as they bear no practical consequence to the present study.

### III. ANALYSIS BASED ON DISCRETE-TIME MAP

From the foregoing simulation results, we have observed that the system may become unstable through slow-scale bifurcation or fast-scale bifurcation in some particular regions of the parameter space. In this section, we will use a discretetime model to analyze these bifurcation phenomena. For the purpose of illustration, we consider operating mode  $M_2$ , and the same analysis procedure can be applied to other operating modes.

From Fig. 1, the state variable vector of the system can be got as  $\boldsymbol{x} = [v_o \ i_{Lb1} \ i_{Lb2} \ v_{pv} \ v_{bat} \ v_e \ v_c]^T$ . As CCM operation is assumed, diode  $D_{pv}$  and switch  $Q_3$  are always in complementary states to switches  $Q_1$  and  $Q_2$ . Then, the state equations for the main circuit are

$$\frac{\mathrm{d}x_1}{\mathrm{d}t} = \frac{-1}{R_L C_f} x_1 + \frac{1-q_1}{C_f} x_2 + \frac{1-q_2}{C_f} x_3$$

$$\frac{\mathrm{d}x_2}{\mathrm{d}t} = \frac{q_1 - 1}{L_{b1}} x_1 + \frac{1}{L_{b1}} x_4$$

$$\frac{\mathrm{d}x_3}{\mathrm{d}t} = \frac{q_2 - 1}{L_{b2}} x_1 + \frac{1}{L_{b2}} x_5$$
(1)
$$\frac{\mathrm{d}x_4}{\mathrm{d}t} = \frac{-1}{C_{\mathrm{pv}}} x_2 - \frac{1}{r_{\mathrm{pv}} C_{\mathrm{pv}}} x_4 + \frac{1}{r_{\mathrm{pv}} C_{\mathrm{pv}}} v_{ocpv}$$

$$\frac{\mathrm{d}x_5}{\mathrm{d}t} = \frac{-1}{C_{\mathrm{bat}}} x_3 - \frac{1}{r_{\mathrm{bat}} C_{\mathrm{bat}}} x_5 + \frac{1}{r_{\mathrm{bat}} C_{\mathrm{bat}}} v_{ocbat}$$

where  $q_1$  and  $q_2$  are the switching functions, given as  $q_i = 1$  for  $Q_i$  is on and  $q_i = 0$  for  $Q_i$  is off.



Fig. 3. Simulated transient waveforms with  $V_{ocpv} = 26$  V and  $I_{mppt} = 0.8$  A. (a) Load current stepped up from 20% load to full load with  $V_{ocbat} = 38$  V and  $\tau_{f1} = 0.39$  ms. The system switches from operating mode  $M_2$  to  $M_1$ , and slow-scale bifurcation occurs in  $M_1$ . (b) Load current stepped down from full load to 20% load with  $V_{ocbat} = 39.99$  V and  $\tau_{f1} = 0.6$  ms. The system switches from operating mode  $M_2$  to  $M_3$ , and it becomes unstable in  $M_3$ .

From the transfer function of PI control and Fig. 1(d), we can write (1, 1)

$$v_2 = k_p \left( 1 + \frac{1}{s\tau_f} \right) \left( V_{\text{ref}} - v_1 \right) \tag{2}$$

Upon differentiating both sides of equation in (2), we get

$$\frac{\mathrm{d}v_2}{\mathrm{d}t} = -k_p \frac{\mathrm{d}v_1}{\mathrm{d}t} - \frac{k_p}{\tau_f} v_1 + \frac{k_p}{\tau_f} V_{\mathrm{ref}} \tag{3}$$

In  $M_2$ , the boost converter operates with a dual loop to regulate the bus voltage, whereas the buck-boost converter operates in buck mode to control the charging current, as shown in Fig. 1(b). For the "Output Voltage Regulator1", we substitute  $v_1 = v_{of} = H_{of}v_o$  and  $v_2 = v_e = x_6$  into (3), and get

$$\frac{\mathrm{d}x_6}{\mathrm{d}t} = -k_{p1}H_{of}\left(\frac{\mathrm{d}x_1}{\mathrm{d}t} + \frac{x_1}{\tau_{f1}}\right) + \frac{k_{p1}}{\tau_{f1}}V_{oref} \tag{4}$$

For the "Charge Current Regulator2", due to the negative logic, we have  $x_7 = v_c = -v_2$ ,  $v_1 = -i_{Lb2}$ , and  $V_{ref} = I_{bmax}$ . Then, using (3), we have

$$\frac{\mathrm{d}x_7}{\mathrm{d}t} = -k_{p2}\frac{\mathrm{d}x_3}{\mathrm{d}t} - \frac{k_{p2}}{\tau_{f2}}x_3 - \frac{k_{p2}}{\tau_{f2}}I_{b\mathrm{max}}$$
(5)

In one switching period, the system toggles among three switch states and the sequence takes the following order: (i):  $Q_1$  and  $Q_2$  are on; (ii):  $Q_1$  is on and  $Q_2$  is off; (iii):  $Q_1$  and  $Q_2$  are both off. Here, as switch  $Q_2$  of the bidirectional buckboost converter is controlled so that its output voltage is equal to the boost converter's. From Table II and  $V_{ocpv} \leq V_{ocbat}$ , we see that the duty cycle of  $Q_1$  is larger than that of  $Q_2$ . Accordingly, the switch state " $Q_1$  is off and  $Q_2$  is on" will not occur. The corresponding state equations for the three switch states can be expressed as

$$\dot{\boldsymbol{x}} = A_{2k}\boldsymbol{x} + B_{2k}\boldsymbol{E}_2 \tag{6}$$

where k = 1, 2, 3 is corresponding to the switch state (i), (ii) and (iii), respectively;  $E_2$  is the vector of input voltage and reference signals in operating mode  $M_2$ , i.e.,  $E_2$  =  $[V_{ocpv} \ V_{ocbat} \ V_{oref} \ I_{bmax}]^T$ . System matrices  $A_{2k}$  and  $B_{2k}$  are given in (7) which is shown on the top of next page, and (8) below.

Denote  $x_n$  and  $x_{n+1}$  as the state variable at the start and the end of the *n*th switching period, respectively. Using (6), we can derive the solution directly as

where  $d_1$  and  $d_2$  are the duty cycles of  $Q_1$  and  $Q_2$  in the *n*th switching period, respectively, which is usually a function of the system's state variables, as will be discussed later;  $\boldsymbol{x}_{21}$  and  $\boldsymbol{x}_{22}$  are the solution at the end of switch state (i) and (ii);  $\boldsymbol{\Phi}_{2k}(\xi)$  and  $\boldsymbol{\Psi}_{2k}(\xi)$  (k = 1, 2, 3) are the corresponding transition matrix and system matrix for  $M_2$ , given as

$$\Phi_{2k}(\xi) = e^{A_{2k}\xi} = I + \sum_{n=1}^{\infty} \frac{1}{n!} A_{2k}^n \xi^n$$
$$\Psi_{2k}(\xi) = \int_0^{\xi} \Phi_{2k}(\xi - \tau) B_{2k} d\tau$$
(10)

By stacking the state variables at the end of each switch state in (9), the discrete-time model can be obtained as

$$\begin{aligned} \boldsymbol{x}_{n+1} &= \boldsymbol{\Phi}_{23} \big( (1-d_1)T_s \big) \boldsymbol{\Phi}_{22} \big( (d_1-d_2)T_s \big) \boldsymbol{\Phi}_{21} (d_2T_s) \boldsymbol{x}_n \\ &+ \boldsymbol{\Phi}_{23} \big( (1-d_1)T_s \big) \boldsymbol{\Phi}_{22} \big( (d_1-d_2)T_s \big) \boldsymbol{\Psi}_{21} (d_2T_s) E_2 \\ &+ \boldsymbol{\Phi}_{23} \big( (1-d_1)T_s \big) \boldsymbol{\Psi}_{22} \big( (d_1-d_2)T_s \big) E_2 \\ &+ \boldsymbol{\Psi}_{23} \big( (1-d_1)T_s \big) E_2 \\ &\stackrel{\text{def}}{=} f(\boldsymbol{x}_n, d_1, d_2, E_2) \end{aligned}$$
(11)

$$A_{2k} = \begin{bmatrix} \frac{-1}{R_L C_f} & \frac{1-q_1}{C_f} & \frac{1-q_2}{C_f} & 0 & 0 & 0 \\ \frac{q_1-1}{L_{b1_1}} & 0 & 0 & \frac{1}{L_{b1}} & 0 & 0 \\ \frac{q_2-1}{L_{b2}} & 0 & 0 & 0 & \frac{1}{L_{b2}} & 0 \\ 0 & \frac{-1}{C_{pv}} & 0 & \frac{-1}{r_{pv}C_{pv}} & 0 & 0 \\ \frac{k_{p1}H_{of}}{R_L C_f} - \frac{k_{p1}H_{of}}{\tau_{f1}} & \frac{(q_1-1)k_{p1}H_{of}}{C_f} & \frac{(q_2-1)k_{p1}H_{of}}{C_f} & 0 & 0 \\ \frac{(1-q_2)k_{p2}}{L_{b2}} & 0 & \frac{-k_{p2}}{T_{f2}} & 0 & \frac{-k_{p2}}{L_{b2}} & 0 & 0 \end{bmatrix}$$

To complete the discrete-time map, the relationship between the duty cycles and the state variables should be derived. During switch states (i) and (ii),  $Q_1$  is always on. Thus, inductor current  $i_{Lb1}$  rises, and when the sampled current  $K_s i_{Lb1}$  reaches the reference level  $I_{ref}$ , switch  $Q_1$  is turned off. Thus, the switching function for  $Q_1$  can be derived as

$$s_{21} = I_{\text{mppt}} + v_e(d_1T_s) - m_c d_1T_s - K_s i_{Lb1}(d_1T_s)$$
  
=  $I_{\text{mppt}} - m_c d_1T_s + C_1 x_{22}$  (12)

where  $C_1 = [0 - K_s \ 0 \ 0 \ 1 \ 0]$ , and  $m_c$  is the compensation slope given by  $m_c = V_p/T_s$ .

In the charging current control loop, the "PWM comparator2" compares the error signals  $v_c$  with the ramp signal:

$$v_{\rm ramp} = V_L + m_{\rm ramp}(t \bmod T_s) \tag{13}$$

where  $V_L$  and  $m_{\text{ramp}}$  are the lower threshold and rising slope of the ramp signal, respectively. Basically,  $Q_2$  is turned on if  $v_c > v_{\text{ramp}}$ , and is off otherwise. Then, we can also get a switching function for  $Q_2$ , as

$$s_{22} = v_c (d_2 T_s) - v_{\text{ramp}}$$
  
=  $C_2 x_{21} - m_{\text{ramp}} d_2 T_s - V_L$  (14)

where  $C_2$  is expressed as  $C_2 = [0 \ 0 \ 0 \ 0 \ 0 \ 1]$ .

Combining (11), (12) and (14), an exact discrete-time model is derived. Here, we will investigate the dynamical behavior of the system by examining the movement of the eigenvalues when some chosen parameters are varied. Suppose the equilibrium point is given as  $\boldsymbol{x} = \boldsymbol{x}_{e2}$ . By setting  $\boldsymbol{x}_{n+1} = \boldsymbol{x}_n = \boldsymbol{x}_{e2}, s_{21} = 0$  and  $s_{22} = 0$ , the equilibrium point  $\boldsymbol{x}_{e2}$  and steady-state duty cycles  $D_1$  and  $D_2$  can be obtained. Then, the Jacobian,  $\boldsymbol{J}_2(\boldsymbol{x}_{e2}) = \partial \boldsymbol{x}_{n+1}/\partial \boldsymbol{x}_n$ , evaluated at the equilibrium point can be derived as

$$\boldsymbol{J}_{2}(\boldsymbol{x}_{e2}) = \frac{\partial f}{\partial \boldsymbol{x}_{n}} + \frac{\partial f}{\partial d_{2}} \frac{\partial d_{2}}{\partial \boldsymbol{x}_{n}} + \frac{\partial f}{\partial d_{1}} \frac{\partial d_{1}}{\partial \boldsymbol{x}_{n}} \bigg|_{\boldsymbol{x}_{n} = \boldsymbol{x}_{e2}}$$
(15)

In (15), the relationship between  $d_i(i = 1, 2)$  and  $x_n$  is implicit in  $s_{21} = 0$  and  $s_{22} = 0$ . Thus, the Jacobian can be derived with implicit function theorem, giving

$$J_{2}(\boldsymbol{x}_{e2}) = \frac{\partial f}{\partial \boldsymbol{x}_{n}} - \frac{\partial f}{\partial d_{2}} \left[ \frac{\partial s_{22}}{\partial d_{2}} \right]^{-1} \frac{\partial s_{22}}{\partial \boldsymbol{x}_{n}}$$
(16)  
$$- \frac{\partial f}{\partial d_{1}} \left[ \frac{\partial s_{21}}{\partial d_{1}} \right]^{-1} \left[ \frac{\partial s_{21}}{\partial \boldsymbol{x}_{n}} + \frac{\partial s_{21}}{\partial d_{2}} \left[ \frac{\partial s_{22}}{\partial d_{2}} \right]^{-1} \frac{\partial s_{22}}{\partial \boldsymbol{x}_{n}} \right]_{\boldsymbol{x}_{n} = \boldsymbol{x}_{e2}}$$

For brevity, we omit the derivation of the Jacobians for operating modes  $M_1$  and  $M_3$ .



Fig. 4. Loci of eigenvalues for (a)  $V_{ocpv}$  is decreased from 34V to 20V with  $I_{mppt} = 0.3A$  in  $M_1$ ; (c) load current increases from 10% load to full load with  $I_{mppt} = 0.8A$  and  $V_{ocbat} = 38V$ ; (e) load current decreases from full load to 10% load with  $I_{mppt} = 0.8A$  and  $V_{ocbat} = 39.99V$ ; (b), (d) and (f) are enlarged views of (a), (c) and (e) showing movement of the pair complex eigenvalues crossing the unit circle. The blue locus corresponds to eigenvalues in  $M_1$ , black refers to  $M_2$  and the purple corresponds to  $M_3$ . Arrows indicate movement directions of eigenvalues.

Now, we solve the characteristic equation for  $J_2(x_{e2})$ :

$$\det[\lambda I - J_2(\boldsymbol{x}_{e2})] = 0 \tag{17}$$

where I is unit matrix. From (17), we can compute all the eigenvalues. If all the eigenvalues are inside the unit circle, the equilibrium state is stable. Any eigenvalue crossing the unit circle from interior to exterior indicates a bifurcation. In particular, if a negative real eigenvalue crosses the unit circle at (-1, 0) point, a fast-scale period-doubling bifurcation

(7)



Fig. 5. Stability boundaries in different parameter space. (a)–(d) for different PI parameter space and the bifurcation is slow-scale Neimark-Sacker bifurcation. (a) System operates in  $M_1$ , with  $V_{ocpv} = 28$  V,  $V_{ocbat} = 36$  V,  $k_{p3} = 10$ ,  $\tau_{f3} = 0.02$  ms; (b) system operates in  $M_2$ , with  $V_{ocbat} = 36$  V,  $k_{p3} = 10$ ,  $\tau_{f3} = 0.02$  ms;  $I_{mppt} = 2$  A; (c)–(d) system operates in  $M_3$  with  $V_{ocpv} = 28$  V,  $I_{mppt} = 2$  A; (c)  $k_{p3} = 10$ ,  $\tau_{f3} = 0.02$  ms; (d)  $k_{p1} = 1$ ,  $\tau_{f1} = 0.01$  ms. (e)–(f) for varying sunlight intensity,  $\tau_{f1}$  and  $V_p$ , with  $k_{p1} = 1$ ,  $k_{p3} = 10$ ,  $\tau_{f3} = 0.02$  ms and  $V_{ocbat} = 36$  V being held constant. (e)  $\tau_{f1} = 0.39$  ms,  $V_p = 0$  V; (f)  $\tau_{f1} = 0.47$  ms,  $V_p = 0.08$  V. Points are stability boundaries obtained from circuit simulations. Stable and unstable regions of operation are located above and below the boundary line, respectively.

occurs; and if a pair of complex eigenvalues move out of the unit circle smoothly, the system undergoes a slow-scale NS bifurcation. Here, when both conditions are satisfied, coexistence of the two types of bifurcation is observed [13]. Moreover, if any eigenvalue "jumps" across the unit circle, a non-smooth bifurcation occurs.

Numerical calculations of the eigenvalues are performed. Tables III to V show typical scenarios of the variation of the eigenvalues under the same condition in Figs. 2 and 3. From Table III, we observe the system loses stability via a smooth NS bifurcation as  $V_{ocpv}$  is decreased to 26V in  $M_1$ . As  $V_{ocpv}$ is further decreased, a negative real eigenvalue crosses the (-1, 0) point, and slow-scale and fast-scale oscillations emerge. From Table IV, it can be seen that the number of eigenvalues becomes 6 and the modulus of one complex conjugate pair is larger than unity when the system switches from  $M_2$  to  $M_1$ , indicating the occurrence of a non-smooth bifurcation. From Table V, it can be found that the system loses stability as long as the system switches from  $M_1$  to  $M_3$  as the load decreases, then this bifurcation is also a non-smooth one.

The corresponding loci for eigenvalues are plotted in Fig. 4. From Figs. 4(a) and (b), we observe a negative real eigenvalue and a pair of complex eigenvalues which moves through the unit circle smoothly. In Figs. 4(c) and (d), one complex conjugate pair "jumps" off the unit circle when  $R_L$  decreases to 39.1  $\Omega$ , indicating the occurrence of a non-smooth bifur-



Fig. 6. Stability boundary with  $V_{ocpv}$  and  $V_p$  serving as bifurcation parameters. Points are stability boundaries obtained from circuit simulations.

cation. From Figs. 4(e) and (f), it can be observed that one complex conjugate pair "jumps" off the unit circle as the system switches from  $M_1$  to  $M_3$ . These numerical results agree well with the bifurcation phenomena observed in the circuit simulations shown earlier in Figs. 2 and 3.

### IV. STABILITY BOUNDARIES IN PARAMETER SPACE

For most practical purposes, only period-1 operation is acceptable, and bifurcation is not permitted. Hence, stability boundaries are very useful design information for engineers

$V_{ocpv}(V)$	Mode			eigenvalues			Modulus	Remarks
32	$M_1$	-0.49288	0.36751	0.61298	0.98926	$0.99921 \pm j 0.03803$	0.99993	Stable
27	$M_1$	-0.77664	0.36793	0.61298	0.98926	$0.99927 \pm j 0.03801$	0.99999	Stable
26	$M_1$	-0.84601	0.36800	0.61298	0.98926	$0.99928 \pm j 0.03800$	1.00000	NS bifurcation
24	$M_1$	-1.00169	0.36813	0.61298	0.98926	$0.99930 \pm j 0.03800$	1.00000	Coexisting

TABLE III Eigenvalues for decreasing  $V_{ocpv}$  with  $I_{\rm mppt}=0.3{\rm A}$  in operating mode  $M_1$ 

TABLE IV	
EIGENVALUES FOR INCREASING LOAD CURRENT (REDUCING LOAD RESISTOR) WITH $I_{ m mppt}=0.8$ A and $V_{ocba}$	t = 38V.

•	/	PP-	000000

$R_L(\Omega)$	Mode			eigenvalues				Modulus	Remarks
48	$M_2$	-0.86349	$0.14803 \pm j 0.35684$	$0.99912 \pm j 0.00607$	0.36801	0.60789	0.38633	0.99914	Stable
40	$M_2$	-0.86489	$0.14803 \pm j 0.35684$	$0.99910 \pm j 0.00606$	0.36801	0.60789	0.38632	0.99912	Stable
39	$M_1$	-0.86427	$0.99998 \pm j 0.03874$	0.36801	0.61298	0.98721		1.00073	slow-scale

TABLE V

Eigenvalues for decreasing load current (increasing load resistor) with  $I_{\rm mppt}=0.8{\rm A}$  and  $V_{ocbat}=39.99{\rm V}$ .

$R_L(\Omega)$	Mode			eigenvalues				Modulus	Remarks
10	$M_1$	-0.86413	$0.99776 \pm j 0.03916$	0.36799	0.51886	0.99161		0.99853	Stable
13	$M_1$	-0.86413	$0.99776 \pm j 0.03924$	0.36799	0.51886	0.99161		0.99853	Stable
14	$M_3$	-0.86436	$0.96344 \pm j 0.29113$	$0.99886 \pm j 0.00478$	0.36798	0.50681	0.99887	1.00647	slow-scale

to identify how far or close the system is from the instability region and in which way the system would loss its stability. In this section, we will generate such stability boundary information. In general, the PI control parameters are important deciding parameters, as  $k_p$  need to be set sufficiently large to reduce the steady-state error while  $\tau_f$  should be small to improve the dynamic response. However, if  $k_p$  is too large or  $\tau_f$  is too small, the system will lose stability. Moreover, in the foregoing sections, we have mentioned that the external parameters, such as  $V_{ocpv}$ ,  $I_{mppt}$ ,  $V_{ocbat}$  and  $R_L$  are varying from time to time. It is therefore appropriate to take these parameters as the bifurcation parameters. Corresponding to the choice of component values given in Tables II, Fig. 5 shows the specific stability boundaries obtained through analysis and circuit simulations at full load. From Figs. 5 (a)-(d), the following observations are made:

- The variation of PI parameters will cause slow-scale NS bifurcation. As  $k_p$  becomes larger or  $\tau_f$  becomes smaller, the system loses stability more readily.
- Smooth NS bifurcation can possibly occur in all operating modes under variation of PI parameters, solar intensity, and input voltages.
- Parameters  $k_{p1}$  and  $\tau_{f1}$  affect the system's stability in all operating modes, and they have narrower operating ranges in  $M_1$ . Hence, they should be carefully designed in  $M_1$ . When the system operates in  $M_1$ , the variation of  $I_{mppt}$  may alter the stability region, and specifically as  $I_{mppt}$  decreases, the stable area shrinks.
- When the system operates in  $M_2$  and  $M_3$ , as parameters  $V_{ocpv}$  and  $V_{ocbat}$  increase, the system will be farther from the instability boundary.

Furthermore, from Figs. 5 (e)–(f), we may draw the following conclusions:

· Smooth slow-scale and fast-scale bifurcations can possi-



Fig. 7. Experimental prototype of multi-operating-mode PBHPS.

bly occur in  $M_1$  and  $M_2$  under variation of  $I_{mppt}$ , or  $V_{ocpv}$  or even compensation voltage  $V_p$ .

- Some parts of the stability boundaries coincide with operating mode boundaries, and they are non-smooth bifurcation boundaries where the system hops from stable to unstable operation as  $I_{mppt}$  or  $V_{ocpv}$  decreases. For example, referring to (e), when  $V_{ocpv}$  is set at 30 V and  $I_{mppt}$  at 2 A, the system operates stably in operating mode  $M_2$ . If  $I_{mppt}$  decreases to 1.3 A, the system will switch from  $M_2$  to  $M_1$  and become unstable, which is a non-smooth bifurcation.
- The input voltage  $V_{ocpv}$  not only affects the slow-scale stability boundaries and non-smooth bifurcation, but also has significant effect on fast-scale stability boundaries, and the fast-scale instability area can be effectively reduced with increasing compensation voltage  $V_p$ .

It should be noted that the above procedure generally permits stability boundaries to be generated with any given set of internal and external parameters serving as bifurcation parameters.



Fig. 8. Typical experimental waveforms in operating mode  $M_1$  for different values of  $V_{ocpv}$ , with  $V_{ocbat} = 38$  V,  $I_{mppt} = 0.3$  A and  $I_o = 3.3$  A kept constant. Alternating component of output voltage  $v_o$  is shown. (a) Stable period-1 operation with  $V_{ocpv} = 32$  V; (b) Neimark-Sacker bifurcation with slow-scale oscillation with  $V_{ocpv} = 26$  V; (c) coexisting fast-scale and slow-scale bifurcation with  $V_{ocpv} = 24.3$  V; (d) enlarged view of (c).

To avoid bifurcation and ensure the system operates within the desired region, we first find the worst-case operating condition for each control loop from the stability boundaries given in Fig. 5. Then, under the worst-case operating condition, we derive the stability boundaries with control parameters serving as the bifurcation parameters and select appropriate values to ensure stability of the control loop. Thus, the system will be stable under all operating conditions. We take the fast-scale bifurcation as an example here. From Figs. 5 (e)-(f), it can be observed that regardless of the operating mode, fast-scale instability is affected only by  $V_{ocpv}$  and  $V_p$ , i.e., the peak current control loop loses its stability when  $V_{ocpv}$  and  $V_p$ are inappropriate. Therefore, to avoid this type of bifurcation, the stability boundary in the  $(V_{ocpv}, V_p)$  parameter space can be derived, as shown in Fig. 6. Here we see that the worstcase operating condition for the peak current control loop is when  $V_{ocpv}$  assumes the smallest value. Then, the peak current control loop will remain stable under the worst-case operating condition if  $V_p > 0.1$ . Thus, as long as  $V_p > 0.1$ , the system exhibits no fast-scale bifurcation under variation of  $V_{ocpv}$ .

### V. EXPERIMENTAL VERIFICATION

A 240 W prototype, as shown in Fig. 7, was built to study the bifurcation phenomena identified and analyzed in the previous sections. The parameters of the prototype used are the same as those used for circuit simulation given in Table II. Fig. 8 shows the experimentally observed slow-scale oscillation and coexisting slow-scale and fast-scale instability in  $M_1$ . It can be found that the slow-scale oscillation mainly occurs in the buck-boost converter. This verifies the loss of stability of the output voltage control-loop through a smooth NS bifurcation. A fast-scale oscillation is found in  $i_{Lb1}$ , as a result of the peak current control-loop becoming unstable as  $V_{ocpv}$  is decreased. These phenomena agree with the simulation results given in Fig. 2 and verify the analysis results in Section III.

A step load change is applied to the PBHPS prototype. Under this test condition, the experimental waveforms are shown in Fig. 9. We can see that the system can properly respond to the step change and switches its operating mode correctly. However, the system becomes unstable while its operating mode is changed, verifying the non-smooth bifurcation as shown earlier in Fig. 3.

The excellent agreement among analysis, circuit simulations and experimental results verifies the ability of the discretetime mapping model in explaining the smooth bifurcation and non-smooth bifurcation, leading to slow-scale or fastscale oscillations, as well as coexisting oscillation. The results presented above reveal the salient phenomena. Others are omitted for brevity of presentation.



Fig. 9. Transient waveforms with  $V_{ocpv} = 26$  V and  $I_{mppt} = 0.8$  A. Alternating component of output voltage  $v_o$  and battery voltage  $v_{bat}$  are shown. (a) Load current steps up from 10% load to full load with  $V_{ocbat} = 38$  V and  $\tau_{f1} = 0.39$  ms kept constant. The system switches from operating mode  $M_2$  to  $M_1$ , and slow-scale bifurcation occurs in  $M_1$ . (b) Load current steps down from full load to 20% load with  $V_{ocbat} = 39.99$  V and  $\tau_{f1} = 0.6$  ms kept constant. The system switches from operating mode  $M_1$  to  $M_3$ , and slow-scale bifurcation occurs in  $M_3$ .

#### VI. CONCLUSIONS

Due to the intermittent nature of the availability of renewable energy, hybrid renewable power generation systems are usually designed to operate with multiple structures and multiple operating modes. The dynamical behavior of such systems is thus rather complex and the design for stable period-1 operation is non-trivial. A commonly used multioperating-mode stand-alone is studied in detail in this paper to illustrate the complex behavior of this kind system. The photovoltaic-battery hybrid power system is found to lose stability via a smooth Neimark-Sacker bifurcation, perioddoubling bifurcation, or a non-smooth bifurcation. These phenomena have been thoroughly analyzed with a discrete-time model. Moreover, stability boundaries are derived in a designoriented form, allowing critical operation conditions be readily extracted for identification of stable operating regions. Salient bifurcation phenomena were demonstrated in an experimental prototype. The smooth and non-smooth bifurcation observed in this paper are generic to the multi-structure and multioperating mode power system, and this line of research is highly relevant to the design of renewable energy systems that have inherent multiple structures and mandatorily operate with multiple modes.

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