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# Elucidating the Hierarchical Architecture of Polymer Spherulites via 4D Scanning Transmission Electron Microscopy

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## ABSTRACT

Hierarchical spherulite structures are ubiquitous in semicrystalline polymers and impact their properties. Elucidating these delicate and complex structures, which span from molecular-level chain folding to mesoscale spherulites, however, presents a formidable challenge. Here, we showcase low-dose four-dimensional (4D) scanning transmission electron microscopy (STEM) as a powerful technique for investigating the multiscale hierarchical structures of polymer spherulites. Applying it to poly( $\epsilon$ -caprolactone) and polyethylene (PE) spherulite films, we reveal the preferential orientation and growth direction of lamellar crystals, as well as the twisted lamella structure in PE banded spherulites. Notably, our observations reveal a non-radial twisting axis forming a spiral texture in PE films. With the enhanced spatial resolution of cryogenic 4D-STEM, we directly visualize individual lamellar crystals at the nanoscale, enabling the identification of chain tilt within a single lamella and the elucidation of lamella configurations at spherulite boundaries. These insights advance our understanding of polymer spherulite crystallization mechanisms and underscore low-dose 4D-STEM as a powerful tool for exploring the intriguing structures of soft materials.

## 1 | Introduction

As the foundation of the plastics and fiber industry, semicrystalline polymers possess unique structures involving repeated chain folding to form lamellar crystals  $\leq 20$  nm thick, connected by amorphous regions, which can further organize into remarkable higher-order formations such as spherulites, dendrites, and ‘shish kebab’ structures [1–7]. The way these hierarchical structures are built, including their constitution,

morphology, and configuration, directly impacts the mechanical, thermal, and electronic properties of polymer products [8–15]. As a typical example, the absence of a thermal gradient during polymer crystallization can make the lamellae grow radially and assemble with spherical symmetry, giving rise to the so-called spherulites. Spherulites can be confined in two dimensions (2D) in the form of disks or films with 2D circular symmetry. Twisting of lamellae around the growth direction may also occur, resulting in banded spherulites with periodic concentric rings

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about the spherulite centers [4, 7]. Such hierarchical spherulite structures can exist in nearly all semicrystalline polymers and profoundly affect their functionalities [1, 4–7]. Probing these complex structures from the molecular-level chain folding to the mesoscale spherulites, however, presents a grand challenge for material characterization. So far, the spherulite structure models have been largely developed based on indirect derivation from observations and measurements on the relatively coarse scale, such as X-ray diffraction (XRD) and spectroscopy, as well as morphological observation based on various microscopy techniques [3–7]. Despite the great success of current models on interpreting various experimental results, the lack of approaches to directly reveal multiscale hierarchical structure has led to ambiguities and uncertainties on the structure details, especially at nanoscale.

Recently four-dimensional scanning transmission electron microscopy (4D-STEM) has emerged as an effective technique to acquire spatially-resolved electron diffraction at each beam-scan position, generating the 4D dataset containing both real-space and reciprocal-space information [16]. The focusing power of STEM can form a nanometer or even atomic sized electron probe, with orders of magnitude enhancement on spatial resolution compared to microfocus XRD on synchrotron [17–25]. The advent of ultrafast direct electron detectors further enables 4D-STEM under a low electron dose, which can probe the pristine structure of organic solids with negligible irradiation damage. 4D-STEM has successfully revealed multiscale structure information on semicrystalline polymers, including phase identification [26], mapping morphology and orientation of polymer crystallites [27–32], orientation of molecular chains [33, 34], and even chain tilt inside individual lamellae [15]. More recently, Kanomi et al. have applied 4D-STEM on polystyrene spherulites to unveil the orientation distribution and branching structure of lamellar crystals [31]. Despite these pioneering studies, however, the potential of novel 4D-STEM on probing the highly hierarchical structure in polymer spherulites, especially the intricate banded spherulites, is yet to be adequately explored.

Here we harness the power of low-dose 4D-STEM to resolve manifold hierarchical structure of spherulite films. By using a small beam-forming aperture, we achieve sharper diffraction spots in 4D-STEM with enhanced signals at the low dose condition. The ultrafast detector allows us to map large areas with fine sampling, and thus to capture multiscale information of spherulite structure, ranging from the circular symmetry of entire spherulites ( $\sim 10\ \mu\text{m}$ ) to the molecular chain orientation inside individual lamellar crystals ( $< 10\ \text{nm}$ ). For nanoscale 4D-STEM with an increased electron dose, the liquid-nitrogen-cooling holder has been adopted to provide cryo-protection on the pristine polymer spherulites. Such low-dose cryo-4D-STEM further enables us to image the morphology of individual lamellar crystals without staining, as well as their configurations at spherulite boundaries that have not been addressed by previous microfocus XRD and 4D-STEM works [17–34]. Moreover, we extend 4D-STEM to the banded spherulites for the first time and discover the structure of lamella twisting around the non-radial axis, which has never been reported before and may be driven by molecular chain tilt that is also detectable by 4D-STEM.

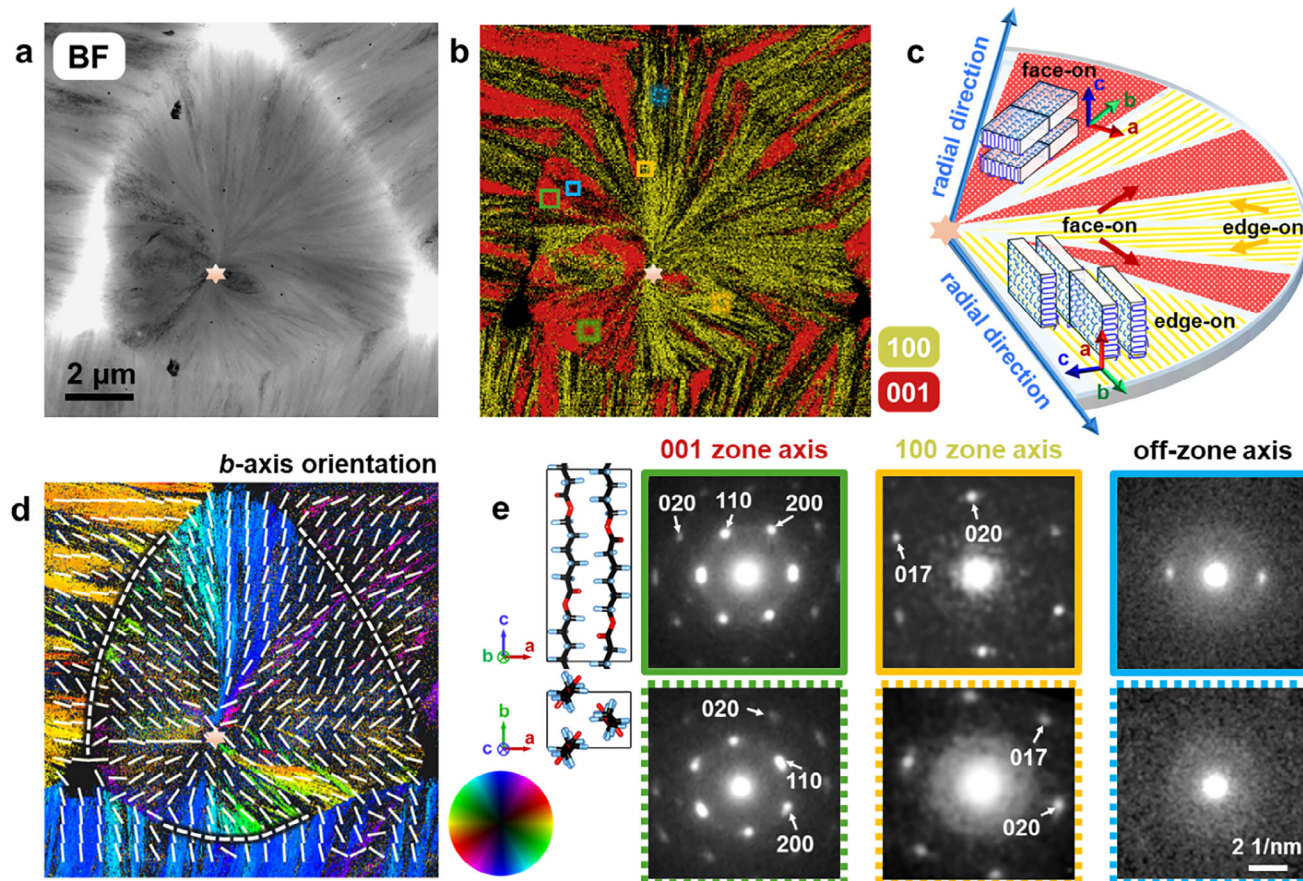
## 2 | Results and Discussion

We prepare poly( $\epsilon$ -caprolactone) (PCL) and polyethylene (PE) spherulite films 100–200 nm thick using solution-based spin coating on  $\text{SiO}_2$ -terminated silicon wafers (see Methods). By etching the  $\text{SiO}_2$  layers away, the freestanding films can be readily transferred to TEM grids for 4D-STEM characterization, avoiding artefacts from TEM sample preparation. Polarized light microscopy (PLM) reveals Maltese crosses on both films (Figures S1, S2), confirming their spherulite nature in 2D with approximate circular symmetry. Figure 1 displays the typical low-dose 4D-STEM (see Methods) result on a PCL spherulite film. A  $\sim 5\ \mu\text{m}$  beam-forming (condenser) aperture is used to form a 3–10 nm electron beam, giving rise to sharp spots in diffraction patterns (Figure 1e). This is in contrast to diffraction disks from larger apertures in most previous 4D-STEM studies, which suffer numerous drawbacks such as poorer reciprocal-space resolution, lower signal-to-noise ratio, and substantial overlap between different diffraction signals [35]. A smaller aperture also helps to reduce the electron dose to  $\sim 1.7\ \text{e}\ \text{\AA}^{-2}$  for acquiring 4D-STEM data. This is sufficiently low for reliably probing the pristine structure of polymers, as demonstrated by our previous work on the critical doses of polymer spherulites [36]. Notably, the peak positions and indexing in diffraction patterns from 4D-STEM data and selected area electron diffraction [36] are well aligned with the reported crystal structure.

### 2.1 | Lamella Orientation Mapping on PCL Spherulites

Figure 1a is a bright-field (BF) image reconstructed from 4D-STEM data, showing the clear morphology of a spherulite with radiating texture originating from the center, which reflects radiating crystallization as the characteristic of spherulites. Further quantitative peak fitting on diffraction signals (see Methods and Figures S3–S5) unveils a wealth of structure information such as lamellar crystal orientation. As shown in Figure 1b,e, the whole spherulite exhibits two primary out-of-plane orientations along [100] and [001] zone axes, corresponding to edge-on and face-on lamellae, respectively (Figure 1c). Regions with off-zone-axis orientation are also identified (Figure 1e) and appear dark in Figure 1b. All these regions seem to primarily nucleate near the spherulite center, grow radially outwards and intermix with each other as illustrated in Figure 1c. Interestingly, the continuous distribution of lamella orientation across PCL spherulite boundaries can be seen (also see Figure S6), which may reflect some clamping effect at the boundaries where different spherulites merge.

Both edge-on and face-on PCL lamellae contain in-plane b axis, which can be used to map the in-plane orientation in spherulites. As shown in Figure 1d, the local b-axis direction is distributed with approximate circular symmetry following the color wheel, which is better visualized by the superimposed line segments for the average b-axis directions: They align primarily along the radial direction in the whole grain, independent of zone-axis switching, demonstrating radial continuity as the essential attribute of spherulites [37, 38]. It reflects radiating crystallization



**FIGURE 1** | 4D-STEM showing the microstructure of PCL spherulites. (a) Reconstructed BF image of PCL spherulites. (b) Zone-axis orientation map in which yellow and red colors indicate regions oriented along the [100] and [001] zone axes, respectively, while all other regions are shown in black. (c) Schematic representation of the PCL spherulite structure, showing radial extension of regions with the two distinct zone axes. (d) Orientation map for the local b-axis derived from 020 diffraction. The colors are determined by mapping the orientation and intensity of 020 diffraction onto the color wheel (inset) within the HSV (Hue, Saturation, Value) color space. Dark colors indicate weak diffraction intensity. The overlaid white lines show the average b-axis orientations over  $\sim 157 \times 157 \text{ nm}^2$  local area. The dash and dotted lines mark the low- and high-angle boundaries, respectively. (e) Diffraction patterns from the squared regions with the same frame styles in (b). Inset displays the projected crystalline structures of PCL along the [010] and [001] directions, with carbon atoms shown in black, hydrogen in blue, and oxygen in red. The electron dose is  $\sim 1.7 \text{ e} \text{ \AA}^{-2}$  for each scanning position, and the scan step is 22.5 nm.

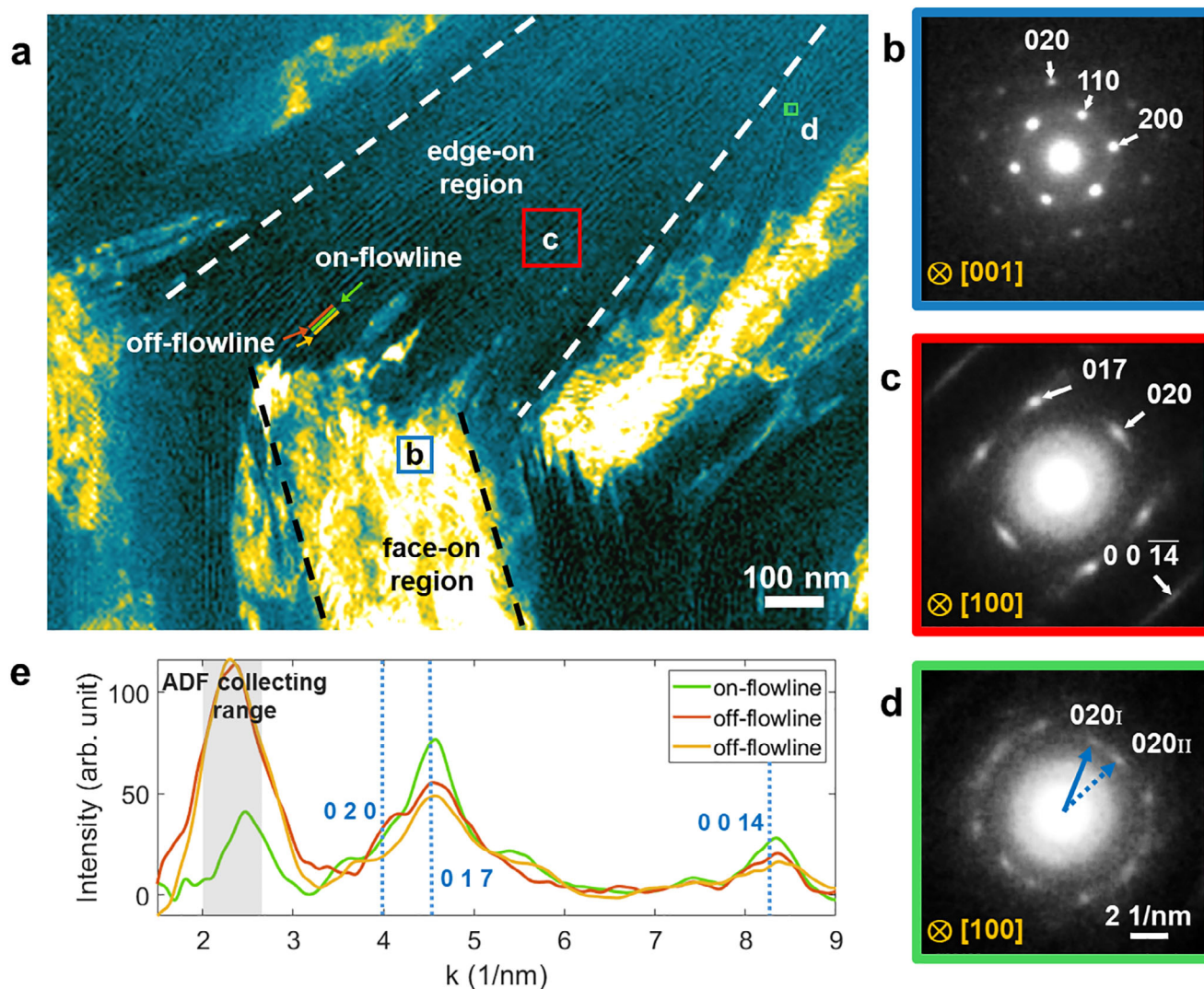
preferentially along b axis to form spherulites. The resultant circular symmetry is broken at the spherulite boundaries, giving rise to low- and high-angle boundaries as indicated by white dashed and dotted lines, respectively.

## 2.2 | PCL Lamellae and Spherulite Boundary Structure Revealed by cryo-4D-STEM

To further probe the nanoscale structure of PCL spherulites, we have carried out high-resolution 4D-STEM with the smaller scanning step size (6.2 nm in Figure 2 vs. 22.5 nm in Figure 1). It leads to the increased electron dose ( $6.0 \text{ e} \text{ \AA}^{-2}$  in Figure 2 vs.  $1.7 \text{ e} \text{ \AA}^{-2}$  in Figure 1) with potentially enhanced beam effects. We thus adopt a liquid-nitrogen-cooling holder to provide cryo-protection on the pristine structure of spherulites, whose effect has been demonstrated in our previous work [36]. Such high-resolution cryo-4D-STEM can resolve individual PCL lamellar crystals at the edge-on orientation, as well as their configuration at spherulite boundaries: The reconstructed annular dark-field

(ADF) image in Figure 2a displays a remarkable pattern consisting of densely packed dark lines  $\sim 8 \text{ nm}$  wide with  $\sim 14 \text{ nm}$  average spacing, consistent with the dimension of lamellar crystals derived from small-angle X-ray scattering (SAXS) measurement (Figure S7). The diffraction pattern taken from the dark line indeed corresponds to [100] zone axis (Figure 2c), confirming them as PCL lamellae at the edge-on orientation. A quantitative comparison of diffraction intensity profiles in Figure 2e further evinces that the dark lines (green curve) exhibit higher [100]-zone-axis diffraction signals (such as 0 1 7 and 0 0 14) but lower intensity in the ADF-integrated range (gray window for amorphous and 110 diffraction) compared to the inter-line regions (red and orange curves), explaining their dark contrast.

Owing to the larger dimension of PCL lamellae than the film thickness ( $\sim 200 \text{ nm}$ ), the edge-on lamellae are likely through the whole thickness without overlapping with other lamellae or amorphous phase, thus resulting in high contrast in Figure 2a. We note that it conventionally requires staining on TEM samples



**FIGURE 2** | 4D-STEM showing the nanoscale lamella structure at a PCL spherulite boundary. (a) Reconstructed ADF image showing a PCL spherulite boundary. (b,c) Diffraction patterns from the blue (b) and red (c) squared regions in (a), corresponding to [001] and [100] zone axes, respectively. (d) Diffraction pattern taken from the green squared region in (a) where two edge-on lamellae intersect. It shows the overlap of two [100] zone-axis diffraction patterns with an in-plane rotation in between, and 020 diffraction points along the dark-line directions. (e) Background-subtracted diffraction intensity profiles relative to reciprocal position  $k$ , obtained by azimuthally averaging the 2D diffraction patterns taken from the three colored lines in (a). This dataset is acquired under cryogenic conditions at  $\sim 100$  K with the electron dose of  $\sim 6.0$  e  $\text{\AA}^{-2}$  and the scan step of 6.2 nm.

to resolve such lamella morphology, while here we highlight the power of 4D-STEM to directly image lamellar crystals in their pristine state. These lamellae stack together and can even extend over 1  $\mu\text{m}$  primarily along the radial direction (Figures S8 and S9), which corresponds to  $b$  axis as the preferential growth direction (indicated by 020 diffraction in Figure 2c). Variations of the dark-line direction in rotated or curved edge-on lamellae are also noticeable, which are accompanied by the rotation of 020 diffraction vectors shown in Figure 2d. They may reflect small-angle branching or lamella splaying, the iteration of which eventually leads to the radial growth of lamellae into a spherulite [31, 39]. On the other hand, the  $c$ -axis, as the PCL chain direction, is always perpendicular to their faces (indicated by 0 0 14 diffraction in Figure 2c), evidencing that there is negligible chain tilt within the PCL lamellae, as reflected in the lamella model in Figure 1c.

Figure 2a also displays the lamella configuration across a spherulite boundary in the middle. Within the spherulite, the face-on lamellae appear the brightest (due to the higher 110 diffraction in the ADF-integrated range), which coexist with edge-on lamellae (dark lines) without evident transition regions, validating our schematic model in Figure 1c. Both edge-on and face-on lamellae can extend up to the boundary, reflecting continuous radiating growth until encountering another spherulite. The boundary is relatively sharp (Figure 1a, Figure S8), where lamellae with different orientations meet without obvious overlap or any amorphous boundary phase. Such sharp boundaries are widely present in PCL films, with another example shown in Figure S6. It is noted that the recent 4D-STEM work on isotactic polystyrene microtomed from a bulk spherulite has shown radiating lamellar crystals with mixed edge-on and face-on orientations similar to Figure 1 [31], suggesting that our observation is

not unique to spherulite thin films only, but also shared with some 3D spherulites. On the other hand, the well-resolved extended morphology of edge-on lamellar crystals and particularly their configuration at spherulite boundaries as shown in Figure 2 have not been reported before.

### 2.3 | Twisted Lamella Structure in PE Banded Spherulites

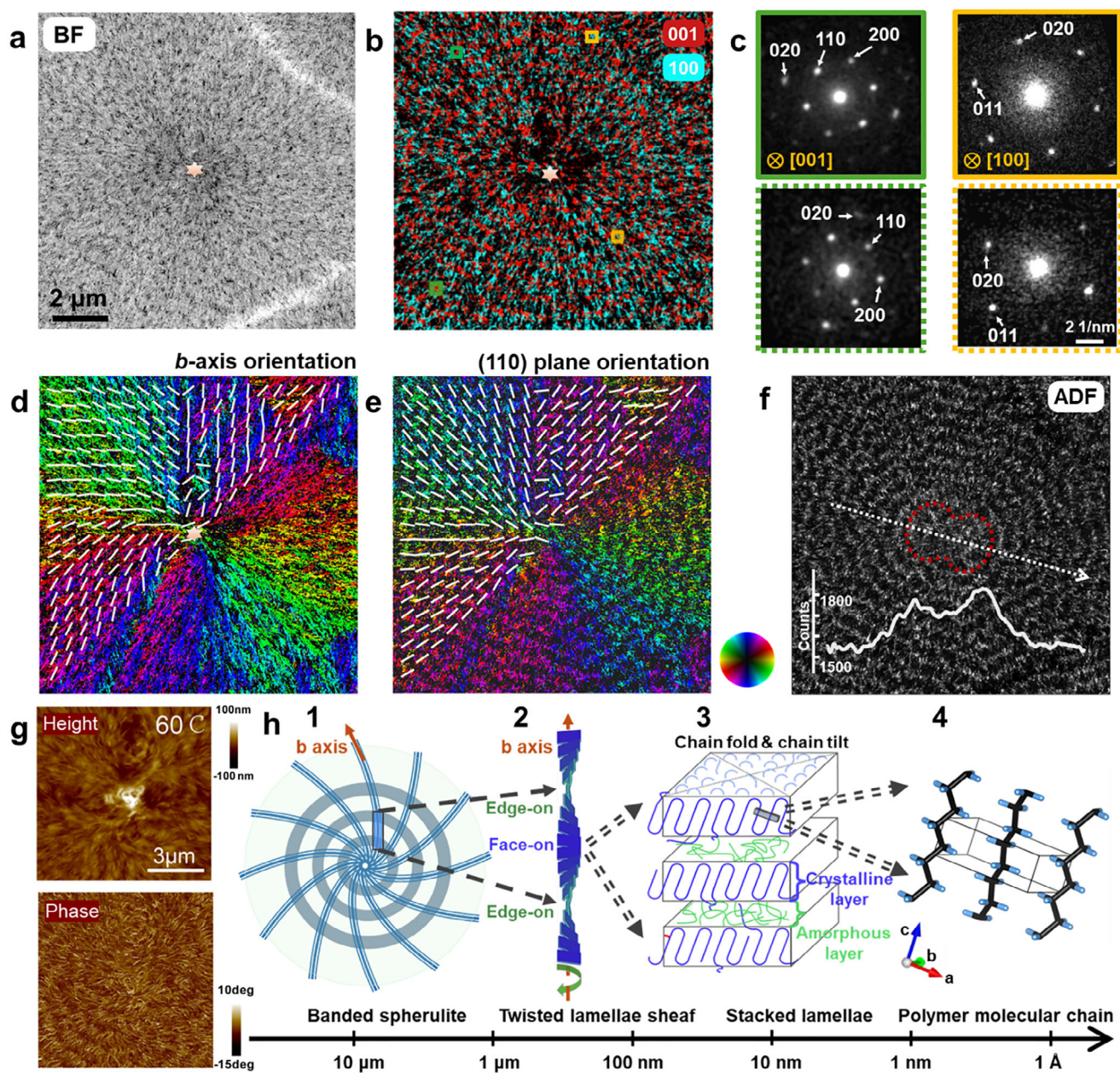
Besides visualizing the hierarchical structure in PCL spherulites, we further apply low-dose 4D-STEM to unravel the intricate structure in PE banded spherulites. As seen in Figure S2, besides the Maltese crosses, PE films also exhibit periodic concentric bands, which are the characteristic of banded spherulites with twisted lamella structure. Such concentric bands can also be detected by 4D-STEM in the reconstructed BF image (Figure 3a), which shows alternating bright and dark contrast surrounding the spherulite center and extending to the boundaries. Zone-axis mapping further reveals the origin of band contrast as the periodic switch between different zone axes of lamellar crystals (Figure 3b), reflecting synchronous lamella twisting periodically about b axis as illustrated in Figure 3h (panel 2). Such a twisting model has been explicitly validated by high-resolution 4D-STEM shown in Figure 4, which demonstrates the zone-axis rotation pathway from  $[001]$  to  $[102]$ ,  $[100]$ , and  $[10\bar{1}]$ , eventually reaching  $[00\bar{1}]$  zone axis. The twisting rate varies substantially between different zone axes (Figure 4f), with an average twisting rate  $\sim 276.3^\circ \mu\text{m}^{-1}$  and a faster twist for edge-on lamellae near  $[100]$  zone axis. The right-handed twisting sense is determined by the shift of the diffraction-pattern mass centers relative to the direct beam (Figure 4e), and has also been confirmed by the systematic tilting experiment (Supporting Information S1: Note 1). Similar characteristics have been identified in PE spherulites with a different band spacing (Figure S12), which can be tuned by controlling the concentration of the PE solution under well-controlled isothermal conditions (see Methods).

We note that the banded structure in PE spherulites can also be revealed by atomic-force microscopy (AFM) phase map, with the hardly detectable height variation associated with the lamella twist in the topography map (Figure 3g), further excluding the effect of height change on the observed band contrast. However, AFM cannot retrieve the structure and orientation information associated with the banded structure. Lamella twisting in banded spherulites has also been previously unveiled by micro-Fourier transform infrared spectroscopy (micro-FTIR), and electron microscopy imaging [40–47]. Nevertheless, pinpointing the detailed twisting pathway with varying crystallographic orientation requires spatially-resolved diffraction such as microfocus XRD, which has been adopted to analyze banded spherulites with microscale band spacing [22–25]. In contrast, 4D-STEM offers much enhanced spatial resolution and can readily resolve nanoscale lamella twisting, as demonstrated in Figures 3 and 4. On the other hand, our 4D-STEM characterization still cannot determine whether the identified twisting is intra- or inter-lamella.

For banded spherulites, the twisting axis is always assumed to be along the radius corresponding to the preferential growth

direction [48–54]. Surprisingly, the radial direction of our PE spherulites corresponds to the  $\{110\}$  plane normal, as unveiled by the in-plane orientation map in Figure 3e. This means that b axis as the twisting axis deviates from the radius by  $\sim 34^\circ$ , which has been explicitly confirmed in Figure S13. With this deviation, the b-axis orientation map displays a spiral pattern (see the superimposed line segments in Figure 3d and the schematic in Figure 3h), which maintains continuity across concentric bands with different zone axes. Despite the numerous literatures on curved lamellae in spherulites [38, 55, 56], such a spiral texture of twisting axis has never been reported before. With the  $\{110\}$  plane normal as the radial direction, it indicates that PE lamellae preferentially expose  $\{110\}$  faces at the growth fronts, which is actually consistent with the large  $\{110\}$  facets observed for PE single crystals, reflecting PE chain folding along  $\{110\}$  planes [57]. Indeed,  $\{110\}$  growth tips have been reported in PE spherulites [58], and sharper 110 diffraction has been detected by XRD, reflecting larger PE lamella size along  $\{110\}$  plane normals [24]. Our observation here further demonstrates that PE lamellae may preferentially grow with  $\{110\}$  faces at the spherulite growth fronts, instead of the long-assumed radial growth along the b axis. On the other hand, PCL spherulites have their b axis along the radius, which is also in accord with the large  $\{010\}$  facets exposed for PCL single crystals [59]. The identified radial directions in our spherulites are thus consistent with the single crystal growth behavior, and unveil a novel PE spherulite structure with the twisting axis consistently deviating from the radius.

It is noted that the mapped bands near the core region in Figure 3b are elongated with a dumbbell shape, which is widely observed in our PE spherulites (Figure S10). Such a shape is further illustrated in the ADF-STEM image in Figure 3f, and is in remarkable agreement with the sheaf-like structures reported in the early stage of spherulitic crystallization [38, 60, 61]. This indicates that the formation of a PE spherulite may initiate with a bundle of lamellar crystals growing preferentially along the lamella faces and much more slowly along the face normal directions. The subsequent radiating growth, enabled by diverging branching of lamellar crystals, gives rise to the sheaf-like structure with lamellae gradually curving back and eventually enveloping the slow-growth regions adjacent to the nucleation center [56]. It can form two cavities near the center of the dumbbell envelop, which may be filled later by residuals or small crystals from the solution, and also referred to as the “double leaf-shaped central bodies” [60, 61]. Indeed, 4D-STEM detects two peaks of the amorphous diffraction signal on both sides of the spherulite center (Figure 3f inset), which presumably correspond to the two cavities filled mostly by disordered residuals, as also revealed in Figure S10a. The identified two disordered regions at the double-leaf core is notably different from the single amorphous core reported in [31]. The revealed higher amorphous content at the spherulite center [62] is consistent with the thermal and chemical stability measurement on our PE spherulites: in situ heating atomic force microscopy (AFM) reveals the initiation of melting at the spherulite center, while ex situ AFM before and after solution etching identifies the initiation of dissolution at the spherulite center (see Supporting Information S1: Note 2). Both observations provide further validation of the spherulite growth mechanism described above.

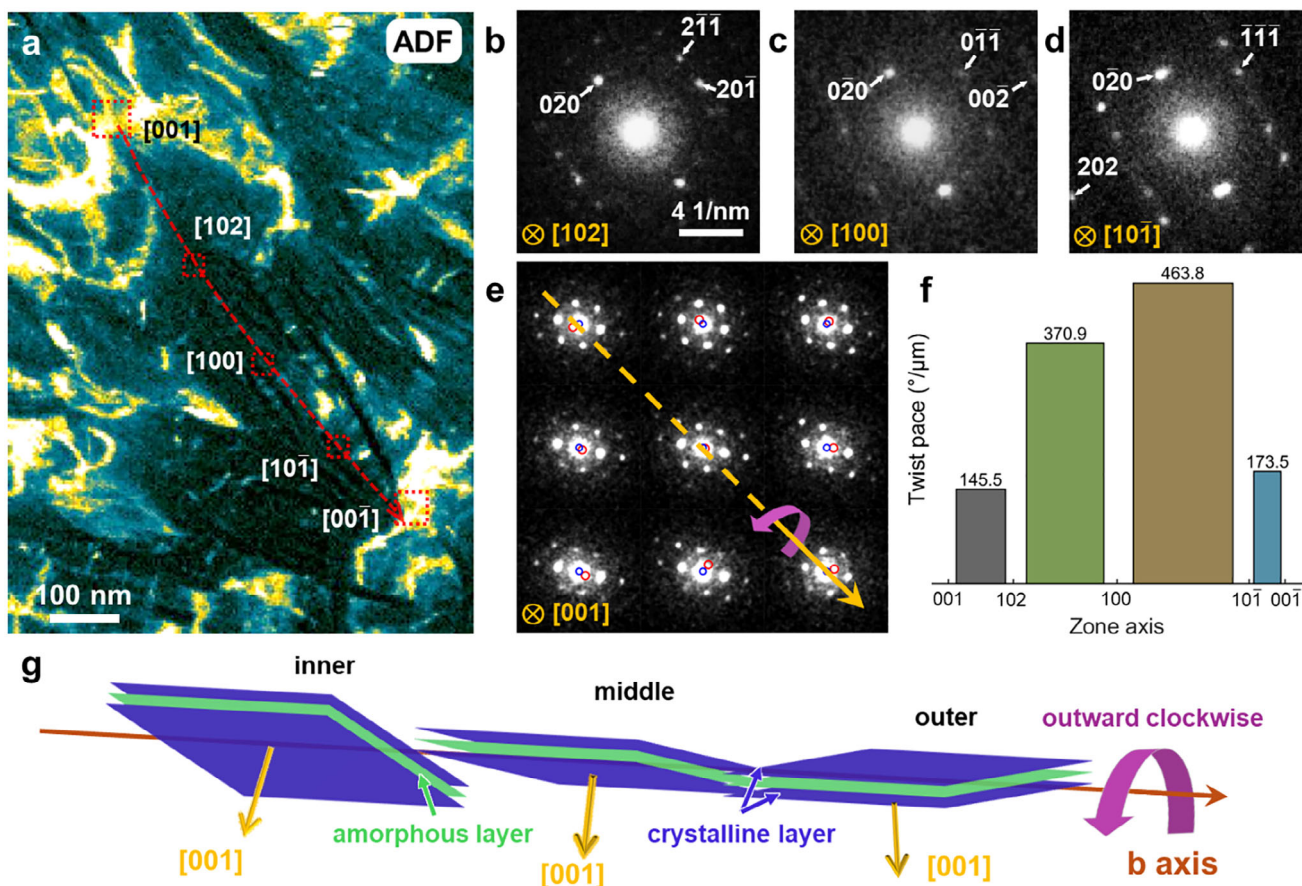


**FIGURE 3** | 4D-STEM showing the microstructure of PE banded spherulites. (a) Reconstructed BF image with the spherulite center indicated by a star. (b) Zone-axis orientation map in which green and cyan colors indicate regions oriented along [100] and [001] zone axes, respectively. (c) Diffraction patterns from the squared regions with the same frame styles in (b). (d,e) Orientation maps for the local b-axis (d) and (110) plane normal (e) derived from 020 and 110 diffraction, respectively. The color scheme is the same as Figure 1d, following the color wheel (inset). The concentric ring pattern in (e) results from the periodic disappearance of 110 diffraction caused by lamella twist. The overlaid white lines represent the average orientations over  $\sim 157 \times 157 \text{ nm}^2$  local area, which are along the radial direction in e but form a spiral pattern in (d). (f) Reconstructed ADF image overlaid with the intensity profile of amorphous diffraction along the dotted arrow, showing two peaks near the center (For amorphous signal analysis see Supporting Information S1: Note 3). The red dash line outlines the dumbbell shape of PE bands near the spherulite center. (g) AFM topography and phase maps of a PE spherulite at  $60^\circ\text{C}$ . (h) Schematic diagram illustrating the hierarchical structure of banded PE spherulites from micrometer scale to atom scale. The electron dose is  $\sim 3.0 \text{ e \AA}^{-2}$  for each scanning position, and the scan step is  $22.5 \text{ nm}$ .

## 2.4 | Individual PE Lamellae With Chain Tilt and Spherulite Boundary Structure

The reconstructed ADF image in Figure 4a also resolves individual PE lamellar crystals at the edge-on orientation, with similar dark-line contrast to PCL lamellae. This is further demonstrated in Figure 5 around the spherulite boundaries. The ADF image

in Figure 5b clearly resolves edge-on PE lamellae as individual dark lines  $\sim 9 \text{ nm}$  wide with  $\sim 18 \text{ nm}$  average spacing, whose dimensions are consistent with SAXS measurement (Figure S9). Such edge-on PE lamellae stack together to form concentric bands, separated by face-on lamellae near the [001] orientation (bright regions in Figure 5b). The periodic twisting limits the visible length of edge-on lamellae to  $\sim 200 \text{ nm}$ , much shorter

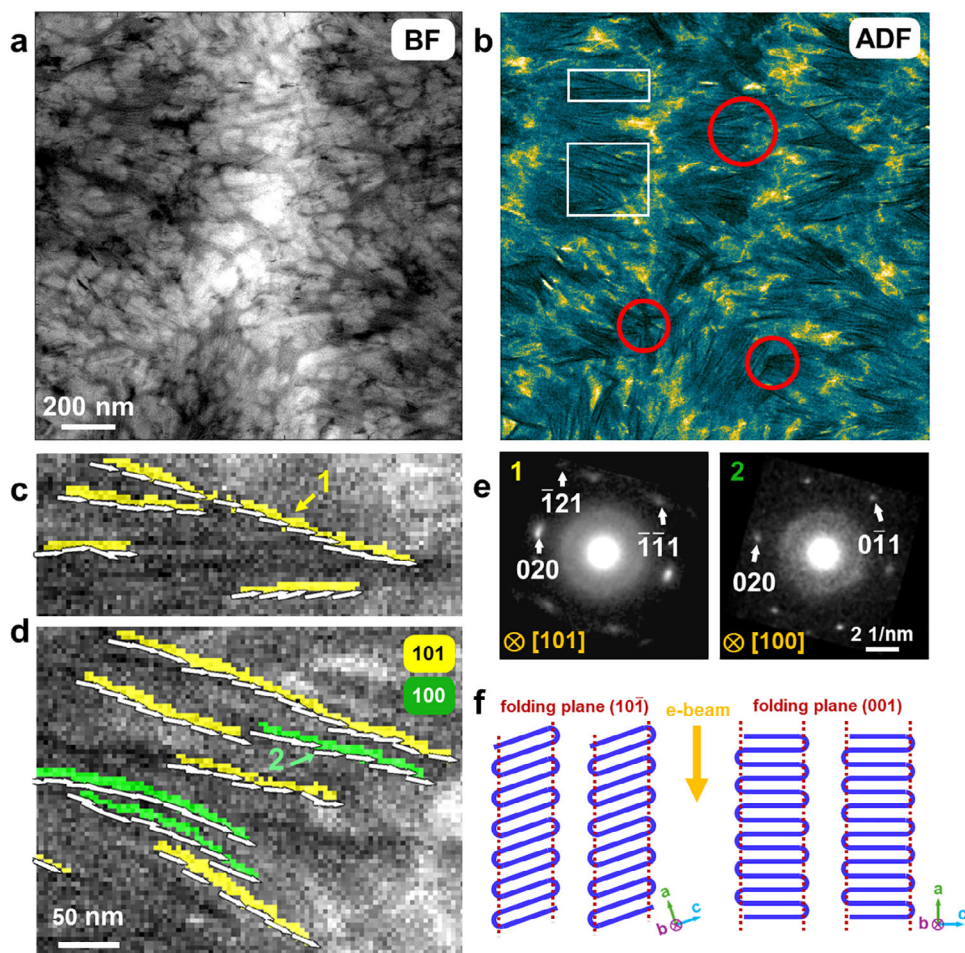


**FIGURE 4** | Twisting structure of banded PE spherulites. (a) Reconstructed ADF image showing two adjacent PE bands, with the detailed twisting pathway along the red arrow indicated by the labelled zone axes. (b–d) Diffraction patterns close to [102], [100], and [101] zone axes, respectively, extracted from the corresponding squared regions in (a). (e) Diffraction patterns close to [001] zone axis, extracted from the squared region in the upper-left part of (a). The blue and red circles represent the origins and centers of mass of the diffraction patterns, respectively. For visualization purposes, the distances between the two circles are magnified by a factor of 4. It shows that along b axis (yellow arrow) toward the periphery, the center of mass moves from left to right, indicating a clockwise rotation sense of PE lamellae from inside out. (f) Twisting rates between different zone axes, derived by dividing the angles between the adjacent zone axes by the distances in between, showing a faster twist for edge-on lamellae near [100] zone axis. (g) Schematic diagram illustrating PE lamella twisting. This dataset is acquired under cryogenic condition at  $\sim 100$  K with the electron dose of  $\sim 6.0$  e  $\text{\AA}^{-2}$  for each scanning position and the scan step of 3.7 nm.

than the well-extended lamellae in PCL spherulites (Figure 2a). The splaying of adjacent lamellae is visible, especially in the spherulite on the right-hand side, which may reflect the short-range repulsion between lamellae to cause radiating growth of spherulites as proposed by Bassett [39]. Similar to PCL, all edge-on PE lamellae are extended along b axis, as indicated by the superimposed arrows in Figure 5c,d. It even applies to the curved lamella in the middle of Figure 5d, suggesting that PE lamellae are still preferentially grown along the b axis. Given the spiral pattern of the b-axis direction revealed in Figure 3d, this implies an overall spiral growth for our PE spherulite films.

On the other hand, in contrast to PCL with the [100]-oriented edge-on lamellae only, PE edge-on lamellae possess two distinct zone axes [100] and [101] (green and yellow regions in Figure 5d,e), with the exposed {001} and {101} faces on the sides, respectively. Given that PE chains are always aligned along c axis, the [100]-oriented edge-on lamellae have their PE chains perpendicular to the exposed {001} faces, thus corresponding

to negligible chain tilt just like PCL lamellae. Whereas for the [101]-oriented edge-on lamellae, their PE chains are inclined to the exposed {101} faces by  $\sim 19^\circ$ . Detailed zone-axis mapping reveals that  $\sim 81\%$  edge-on lamellae are [101]-oriented in our PE spherulites (Figure S16), giving rise to the average chain-tilt angle  $\sim 16^\circ$ . The presence of chain tilt in PE crystals has been extensively reported and proposed as the driving force for lamella twisting, through generating distinct stress conditions at the two opposite fold surfaces [63]. It may explain the difference between PE and PCL spherulites, the latter of which show neither chain tilt nor lamella twisting. However, our measured chain-tilt angle is much smaller than the reported  $\sim 35^\circ$  chain tilt for PE crystals with the preferentially exposed {201} faces [55, 58, 64, 65], which is not observed in our PE spherulites. The recent 4D-STEM study on high-density PE also revealed smaller chain tilt in stacked lamellae ( $\leq 15^\circ$ ) compared to isolated lamellae ( $\sim 34^\circ$ ) [15]. This provides justification for our observation, given that most lamellar crystals observed in our PE spherulites are stacked with high density compared to those observed in [15].



**FIGURE 5** | PE spherulite boundaries and lamella structure. (a,b) Reconstructed BF (a) and ADF (b) images at a triple junction of PE spherulite boundaries. (c,d) Zoomed-in ADF images from the selected regions in (b), with the overlaid colors indicating the two zone axes ([101] and [100]) of individual lamellae (see more in Figure S16). The overlaid white arrows below the lamellae depict the b-axis direction. (e) Diffraction patterns from the designated edge-on lamellae in (c,d) reveal the [101] and [100] zone axes, respectively. (f) Schematic diagram illustrating two zone axes of edge-on lamellae caused by the chain tilt. The electron dose is  $\sim 6.8 \text{ e} \text{ \AA}^{-2}$  for each scanning position, and the scan step is 3.0 nm.

There are boundaries between three spherulites with bright contrast in the BF image (Figure 5a), reflecting reduced density or thickness at the boundaries (also see Figure S20). Figure 5b shows that the lamella bands merge at the boundaries with no enlarged spacing, indicating similar structure and thus density at the boundaries (also see Figures S21 and S22). Therefore, the observed bright contrast of boundaries in Figure 5a should reflect their considerably reduced thickness when the spherulites meet. This has been explicitly confirmed by AFM measurements (Supporting Information S1: Note 2), and may reflect the exhaustion of crystallizable material in the boundary regions [61]. On the other hand, the overlap and intertwining between lamellae can be identified at the boundaries (highlighted by circles), indicating the interpenetrating merging style among PE spherulites that is distinct from the sharp PCL boundaries (Figure 2a).

### 3 | Conclusion

In summary, we have demonstrated the power of low-dose 4D-STEM on probing the manifold hierarchical structure

of spherulite films, including orientation mapping of entire spherulites ( $\sim 10 \text{ }\mu\text{m}$ ), imaging individual lamellar crystals ( $< 10 \text{ nm}$ ), as well as deriving the molecular chain orientations inside them. Applying these capabilities on PE and PCL spherulites enables us to distinguish their characteristic structures, such as the radiating texture in PCL in contrast to the spiral texture in PE, as summarized in Table 1. In particular, the enhanced resolution of cryo-4D-STEM is able to directly resolve individual edge-on lamellar crystals  $< 10 \text{ nm}$  thick. It can not only reveal how the lamellae splay or twist to form the entire spherulites, but also unravel the configurations at spherulite boundaries that are difficult to access by other techniques. It further allows us to identify the preferential growth of PE lamellae with  $\{110\}$  faces at the spherulite growth fronts, in contrast to the long-assumed radial growth along b axis. We also derive the chain tilt within individual lamellae, which may be the driving force for lamella twisting to form banded spherulites. The achieved structure information not only showcases the effectiveness of low-dose 4D-STEM on interrogating the fascinating hierarchical structures of polymers, but also provides valuable insights into the crystallization mechanism of polymer spherulites.

**TABLE 1** | A comparison between PCL, PE, and the existing model for PE spherulites.

	Preferential growth direction	Radial direction	Lamella orientation	Chain tilt on average	Boundary structure
PCL	b axis	b axis	mostly [100] and [001]	negligible	non-penetrating type
PE	b axis	{110} plane normal	continuous twisting around b axis	~16°	interpenetrating merging type
existing model	b axis [48–54]	b axis [48–54]	twisted	~35° [55, 58]	unknown

## 4 | Methods

### 4.1 | Sample Preparation

16 and 20 mg of Polyethylene with a weight average molar mass ( $M_w$ ) of  $\sim 110 \text{ kg mol}^{-1}$ , a density of  $\sim 0.95 \text{ g cm}^{-3}$  (Alfa Aesar) was initially dissolved in 1 mL of decalin (Sigma-Aldrich) at a temperature of  $135^\circ\text{C}$  for 6 h to ensure thorough dissolution. Prior to spin-coating, the following steps were carried out to ensure precise temperature control during the entire process video [SI](#). (1) The spin coater was pre-heated to  $100^\circ\text{C}$  realized by the temperature controlled spin-coating instrument (LEBO SCIENCE), and a silicon substrate was placed on the heated chuck for at least 5 min. (2) The glass pipette used to transfer the PE/decalin solution was immersed in decalin maintained at  $135^\circ\text{C}$  on a hot stage, in order to avoid undesired crystallization induced by temperature differences between the pipette and the PE/decalin solution. After pre-heating, 3–4 drops of the PE/decalin solution were quickly transferred onto the heated substrate using the pre-heated pipette, and the preset spin-coating program (4000 rpm, 60 s) was immediately initiated. The entire solution-transfer process was completed within 30 s. During spin-coating, the airflow generated by rotation caused a slight decrease in temperature; however, due to the built-in temperature-control system, the temperature fluctuation did not exceed  $5^\circ\text{C}$ . After spin-coating, the substrate with the PE film was kept on the  $100^\circ\text{C}$  heated spin coater for an additional 5 min to ensure complete spherulitic growth. Poly( $\epsilon$ -caprolactone), weighing 80 mg, with a  $M_w$  of  $\sim 14 \text{ kg mol}^{-1}$  (Sigma-Aldrich), was dissolved in 1 mL of m-xylene (TCI) at room temperature for 1 h. It was then spin-coated onto wafers at 4000 rpm for 30 s under room temperature.

The prepared thin films were then transferred onto TEM grids: the thin films with underlying wafers were immersed in hydrofluoric acid aqueous solution to etch the  $\text{SiO}_2$  interlayers away. The thin films that floated to the surface of the solution were subsequently collected using TEM grids. Thin film thickness was measured with a profilometer (DEKTAK XT, BRUKER). PE films were  $\sim 100$  and  $140 \text{ nm}$  thick, while PCL films were about  $200 \text{ nm}$  thick. Higher PE solution concentration resulted in thicker films. Optical microscopy (OM) images, as presented in [Figure S1a,b](#), demonstrate the domains and boundaries of spherulites within the thin films. Additionally, images captured using polarized light microscopy (PLM), as shown in [Figure S1c,d](#) and [S2b](#), reveal the presence of Maltese cross patterns, indicative of polymer spherulite structure. Notably, the PLM image of PE ([Figure S2b](#)) displays an additional feature, extinction bands, suggesting a periodic twist in the lamellae. The contrasts observed in PLM images are attributed to the birefringent properties of polymer crystals, which can split an incident polarized light into two rays,

with different speeds and polarized directions. In the banded spherulite structure, the varying projections of the local indicatrix led to changes in the intensity of the emerging polarized light, as illustrated in the schematic in [Figure S2a](#). However, due to the limited resolution of PLM, it cannot directly observe the internal structures and local orientations of the spherulites.

### 4.2 | Transmission Electron Microscopy

The 4D-STEM datasets were acquired using an EMPAD, a high dynamic range direct electron detector enabling electron counting [66], on a ThermoFisher Scientific Spectra 300 microscope operated at 300 kV in STEM mode. The datasets utilized a small convergence semi-angle of  $110\text{--}270 \mu\text{rad}$ , a beam current of 2 pA (measured by EMPAD), an exposure time of 1 ms and a camera length of 1.15 m unless specified otherwise. The Gaussian-shaped probe size was measured to be  $9.6 \text{ nm}$  for  $110 \mu\text{rad}$  convergent semi-angle (similar to the predicted value based on the convergence angle [67]) at full-width at half-maximum (FWHM). Therefore, the instantaneous dose per probe position ( $\text{Dose} = 4I_{\text{Current}} \cdot t_{\text{exposure}} / \pi d_{\text{FWHM}}^2$ ) can be calculated by using these parameters, amounting to  $1.7 \text{ e } \text{\AA}^{-2}$ , permitting effective acquisition at room temperature.

However, when the scanning step was smaller than the probe size, which corresponds to over-sampling, one position is exposed repeatedly. In such instance, the average dose per area over the scan ( $\text{Dose} = I_{\text{Current}} \cdot t_{\text{exposure}} / L_{\text{scanning step}}^2$ ), calculated from the total electrons over scanning area, escalated to  $15.6 \text{ e } \text{\AA}^{-2}$  for a scanning step of  $2 \text{ nm}$ , which could rapidly destroy the polymer structure. Thus, the presented results were obtained with a scanning step greater than  $3 \text{ nm}$ , as presented in figure captions. To mitigate beam damage, a cryogenic environment with liquid  $\text{N}_2$  ( $\sim 100 \text{ K}$ ) was utilized [36, 68, 69] by a Gatan cryo-transfer holder, model 915. Cryogenic protection can enhance the critical dose of PE and PCL by  $\sim 7$  times compared to room temperature, which has been demonstrated in our previous work in [36]. The findings under cryogenic conditions were presented in [Figure 2](#) and [Figure S9](#) for PCL and [Figure 4](#) and [Figure S21](#) for PE.

The 4D STEM datasets were processed and analyzed using homemade MATLAB scripts. Diffraction patterns were first aligned using the central beam and subsequently transformed into a polar coordinate system, which facilitated the extraction of intensity profiles as functions of azimuthal angle and radial distance ([Figure S4](#)). The reconstructed BF and ADF images were generated by integrating the corresponding signals ( $<0.2$  and  $2.0\text{--}2.6 \text{ nm}^{-1}$ ) in diffraction patterns. The range between  $2.0$  and  $2.6 \text{ nm}^{-1}$  corresponds to the amorphous and  $110$  diffraction signals.

The b-axis orientation (Figures 1d and 3d) can be determined by fitting the 020 diffraction peak observed in the azimuthal intensity profile (with signals integrated over 3.9–4.1 nm<sup>-1</sup>). When multiple diffraction peaks are present, the local orientation is given by the azimuth of the strongest peak. If all diffraction intensity falls within three standard deviations of the mean intensity, the pattern is treated as no discernible diffraction peak. For the (110) plane orientation (Figure 3e), it is necessary to account for the symmetry-equivalent 110 diffraction spots. Zone-axis orientation (Figures 1b and 3b) is determined based on the strong diffraction signals: The zone axis is near [001] when both 110 and 020 diffraction intensity exceed 20% of the average intensity of the whole pattern. Likewise, the [100] zone axis is identified based on the (020) and (011)/(017) diffraction intensity for PE/PCL. The amorphous signal (profile in Figure 3f) can be fitted by a broad peak as demonstrated by our previous work [36] and Supporting Information S1: Note 3.

Electron energy loss spectroscopy (EELS) and energy-filtered transmission electron microscopy (EFTEM) were conducted on a Gatan GIF Continuum K3 System. EELS mapping is conducted in scanning mode with the convergence semi-angle of 28.7 mrad and collection semi-angle of 21.1 mrad.

Small-angle X-ray scattering (SAXS) measurements were performed on a Rigaku X-ray diffractometer (SmartLab 9 kW). The SAXS correlation auto-function analysis [70–72] is shown in Figure S7.

AFM measurements (Figures S14 and S15) were respectively conducted using a Nanoscope IIIa Multimode equipped with a temperature controlling sample stage (Digital Instruments) and a Veeco 3D, both operating in tapping mode. Silicon tips with a resonance frequency of ~303 kHz and a spring constant of about 40 N m<sup>-1</sup> were used. In the in situ heating experiments, prior to an AFM scanning at the new temperature, a 30-min temperature maintenance period was implemented to ensure the samples fully equilibrated at the desired temperature.

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## Conflicts of Interest

The authors declare no conflicts of interest.

## Data Availability Statement

One raw 4D-STEM dataset for Figure 1 is available in Zenodo a <https://doi.org/10.5281/zenodo.18067254> [73]. Additional datasets generated and analyzed during the current study are available from the corresponding author upon reasonable request.

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## Supporting Information

Additional supporting information can be found online in the Supporting Information section.

**Supporting File:** sml172608-sup-0001-SuppMat.docx.

**Supporting File:** sml172608-sup-0002-VideoS1.mp4.