

# Enabling Robust N-Type Perovskite Field-Effect Transistors Through an TiO<sub>2</sub> Interlayer Strategy

Jiangnan Xia, Xincan Qiu, Ping-An Chen, Yu Liu, Jiaqi Ding, Yu Zhang, Huan Wei, Zhenqi Gong, Chengyuan Peng, Wenpei Shi, Shuanglong Wang, Chen Chen,\* and Yuanyuan Hu\*

Metal halide perovskites (MHPs) show tremendous potential for field-effect transistors (FETs), but N-type Pb-based MHP FETs have been hindered by critical challenges, including high defect densities, ion migration, and poor reproducibility. In this work, a simple yet powerful ultrathin TiO<sub>2</sub> interlayer strategy is introduced that fundamentally transforms the fabrication of Pb-based MHP FETs. By pre-depositing an ultrathin TiO<sub>2</sub> layer before perovskite film deposition, reproducible and operationally stable MAPbI<sub>3</sub> FETs with remarkable performance are achieved. Comprehensive characterizations reveal that the TiO<sub>2</sub> interlayer enhances precursor wetting, promotes larger and more uniform grain formation, reduces defect density, and effectively suppresses non-radiative recombination and ion migration. The universality of this approach is demonstrated by successfully extending it to 2D Dion-Jacobson phase perovskites, including PDAPbI<sub>4</sub> and its derivatives. The fabricated devices exhibit excellent electrical characteristics, including high on/off ratios, low hysteresis, and impressive stability. As a proof of concept, a complementary inverter is constructed using perovskite-only components, showcasing the potential for integrated logic circuits. This work provides a robust fabrication method for high-performance Pb-based perovskite FETs with broad applicability.

## 1. Introduction

Metal halide perovskite (MHP) semiconductors have emerged as highly promising candidates for field-effect transistors (FETs) applications, owing to their exceptional solution processability and high intrinsic charge carrier mobility.<sup>[1-3]</sup> In recent years, significant advancements have been achieved in P-type FETs based on Sn-based MHP semiconductors, with hole mobilities exceeding 50 cm<sup>2</sup>/V·s, demonstrating remarkable performance potential.<sup>[4-7]</sup> Despite significant advancements in the development of P-type MHP FETs, their N-type counterparts, particularly Pb-based MHP FETs, have exhibited markedly slower progress.

Pb-based MHP semiconductors, known for their ambipolar charge transport characteristics, have been widely utilized in N-type FET devices. However, the performance of Pb-based MHP FETs is often compromised by inherent material

J. Xia, P.-A. Chen, Y. Zhang, Z. Gong, C. Peng, W. Shi, Y. Hu  
International Science and Technology Innovation Cooperation Base for  
Advanced Display Technologies of Hunan Province  
College of Semiconductors (College of Integrated Circuits)  
Hunan University  
Changsha 410082, China  
E-mail: [yhu@hnu.edu.cn](mailto:yhu@hnu.edu.cn)

J. Xia, J. Ding, H. Wei, Y. Hu  
Key Laboratory for Micro/Nano Optoelectronic Devices of Ministry of  
Education  
School of Physics and Electronics  
Hunan University  
Changsha 410082, China

X. Qiu, Y. Liu, Y. Hu  
Key Laboratory of Hunan Province for 3D Scene Visualization and  
Intelligence Education  
School of Electronic Information  
Hunan First Normal University  
Changsha 410205, China

C. Chen  
Science and Technology on Advanced Ceramic Fibers and Composites  
Laboratory  
College of Aerospace Science and Engineering  
National University of Defense Technology  
Changsha 410000, China  
E-mail: [angela\\_chen@nudt.edu.cn](mailto:angela_chen@nudt.edu.cn)

S. Wang  
Department of Applied Physics  
The Hong Kong Polytechnic University  
Hong Kong 999077, China

 The ORCID identification number(s) for the author(s) of this article can be found under <https://doi.org/10.1002/adv.202516610>

© 2025 The Author(s). Advanced Science published by Wiley-VCH GmbH. This is an open access article under the terms of the [Creative Commons Attribution](#) License, which permits use, distribution and reproduction in any medium, provided the original work is properly cited.

DOI: 10.1002/adv.202516610

challenges, including high carrier defect densities, structural phase transitions, and ion migration, which collectively manifest in severely hysteretic transfer characteristics and, in extreme cases, complete loss of gate modulation capability.<sup>[8,9]</sup> Although MAPbI<sub>3</sub> (MA = Methylammonium) FETs capable of operating at room temperature have been reported through strategies such as solvent passivation,<sup>[10]</sup> additive engineering,<sup>[11,12]</sup> and single-crystal synthesis,<sup>[13,14]</sup> these approaches often involve intricate experimental procedures and high operational complexity, limiting their reproducibility. This limitation significantly hinders the progress in optimizing the performance of Pb-based MHP FETs. Therefore, the development of a simple, efficient, and highly reproducible fabrication method is of paramount importance for advancing the performance of Pb-based MHP FET devices.

In this work, we propose a simple, robust interlayer strategy for fabricating Pb-based FETs. Using MAPbI<sub>3</sub> as a model system, an ultrathin TiO<sub>2</sub> interlayer is pre-deposited on the substrate prior to perovskite film formation. This modification yields reproducible and operationally stable MAPbI<sub>3</sub> FETs. Structural, morphological, and electrical characterizations reveal that the adoption of the TiO<sub>2</sub> interlayer benefits the wetting of the perovskite precursor solution, improving the film morphology, effectively reducing defects in the MAPbI<sub>3</sub>, suppressing non-radiative recombination, and dramatically mitigating ion migration, all of which lead to the enhancement of device performance. We further demonstrate the generality of the approach: it is effective not only for 3D Pb-based perovskites but also for other Pb-perovskite compositions. Notably, using this method, we have, for the first time, fabricated large-area thin-film FET arrays based on a 2D Dion–Jacobson (DJ) phase Pb perovskite (PDAPbI<sub>4</sub>) that operate at room temperature. Moreover, we built a complementary inverter entirely from perovskite FETs by combining an N-type PDAPbI<sub>4</sub> device with a P-type PEA<sub>2</sub>SnI<sub>4</sub> device, demonstrating basic circuit functionality. Our work therefore provides a simple, effective, and reproducible fabrication route that lays a practical foundation for advancing Pb-based N-type MHP FETs and their integration into functional circuits.

## 2. Results and Discussion

### 2.1. TiO<sub>2</sub> Interlayer Strategy for MAPbI<sub>3</sub> FET Fabrication

Traditional fabrication of MAPbI<sub>3</sub> FETs faces two major challenges that hinder room-temperature operation and reproducibility: first, the requirement for complex defect passivation processes to improve film quality,<sup>[10–12]</sup> and second, the inherent limitations imposed by anti-solvent crystallization techniques on fabrication yield.<sup>[15]</sup> Consequently, the achievement of N-type MAPbI<sub>3</sub> FETs that can function well at room temperature remains a formidable challenge, significantly impeding the advancement of N-type perovskite FETs.

We streamlined the fabrication process of reference MAPbI<sub>3</sub> FETs. As illustrated in **Figure 1a**, the MAPbI<sub>3</sub> precursor solution was prepared using lead acetate trihydrate (PbAc<sub>2</sub>) and MAI as starting materials, then spin-coated directly onto pre-patterned source/drain (S/D) electrodes on a SiO<sub>2</sub>/Si substrate, forming the MAPbI<sub>3</sub> active layer in a single-step, anti-solvent-free process. **Figure S1** (Supporting Information) demonstrates that

the MAPbI<sub>3</sub> films prepared without anti-solvents exhibit excellent crystallinity; the primary X-ray diffraction (XRD) diffraction peaks at 14.6°, 28.9°, and 43.6° correspond to the (110), (220), and (314) planes of MAPbI<sub>3</sub>, respectively. The UV-vis absorption spectrum of MAPbI<sub>3</sub> shows a steep absorption edge at 760 nm, and the steady-state photoluminescence (PL) spectrum peaks at 765 nm, corresponding to an optical bandgap of ≈1.55 eV. These results are consistent with previous studies.<sup>[16,17]</sup> Despite the high film quality, the fabricated bottom-gate bottom-contact (BGBC) FETs show no gate modulation. As shown in **Figure 1c,d**, the transfer and output characteristic curves demonstrate a complete loss of gate modulation, with the channel current (I<sub>DS</sub>) showing no response to gate voltage variations. This absence of field-effect modulation is likely caused by the intrinsic ionic migration and defect-related trapping in the MAPbI<sub>3</sub> films, which screens the gate field and prevents accumulation/depletion of mobile carriers in the channel.

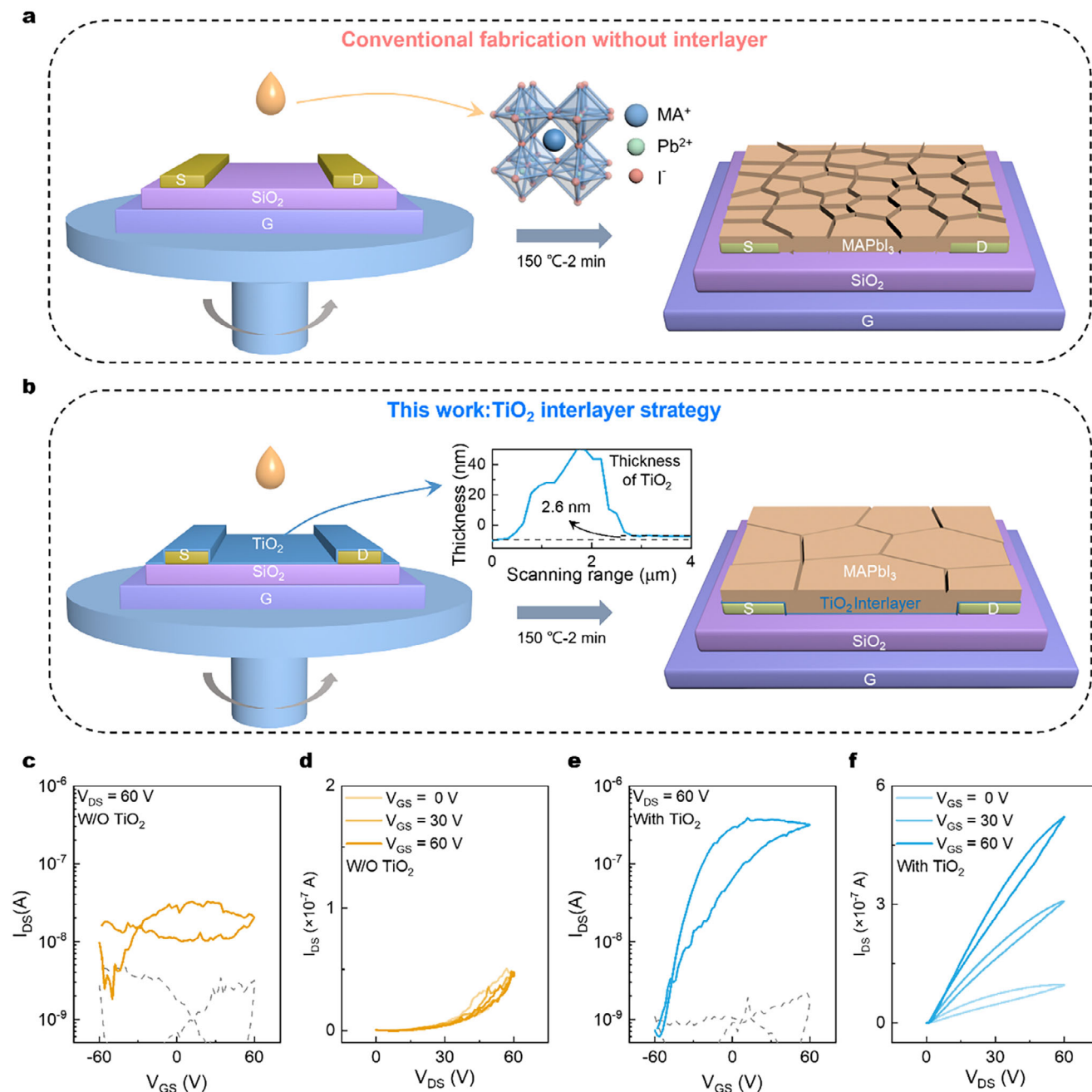
Notably, when an ultrathin TiO<sub>2</sub> interlayer (≈2.6 nm) was pre-deposited prior to spin-coating the MAPbI<sub>3</sub> active layer (**Figure 1b**; **Figure S3**, Supporting Information), the MAPbI<sub>3</sub> FETs performance exhibited remarkable enhancement. As shown in **Figure 1e,f**, MAPbI<sub>3</sub> FET devices with the TiO<sub>2</sub> interlayer demonstrated several orders of magnitude improvement in both on/off ratio and on-state current under identical gate voltage sweeps. The output characteristic curve now clearly exhibited gate-voltage modulation, displaying N-type FET behavior.

To rule out the possibility that the channel current originates from the TiO<sub>2</sub> interlayer itself, we fabricated BGBC FETs with containing TiO<sub>2</sub> only in the channel. As shown in **Figure S4** (Supporting Information), these TiO<sub>2</sub>-only FETs exhibit channel currents below 10<sup>−9</sup> A across the entire gate voltage sweep range, demonstrating insulating behavior. This result confirms that the measured current in the MAPbI<sub>3</sub> FET originates from charge transport in the MAPbI<sub>3</sub> active layer rather than conduction through the TiO<sub>2</sub>.<sup>[18]</sup>

### 2.2. Top-Gate MAPbI<sub>3</sub> FET Fabrication via TiO<sub>2</sub> Interlayer

We fabricated top-gate bottom-contact (TGBC) structure MAPbI<sub>3</sub> FETs using the TiO<sub>2</sub> interlayer. As illustrated in **Figure 2a**, we first selected common polymer dielectric Cytop as the dielectric for TGBC structure MAPbI<sub>3</sub> FETs with TiO<sub>2</sub> interlayer (**Figure 2b**). The transfer characteristics curves of the Cytop-based MAPbI<sub>3</sub> FETs are presented in **Figure 2c**, which exhibit typical N-channel transistor characteristics. Importantly, in such TGBC devices, the channel is formed at the top dielectric/MAPbI<sub>3</sub> interface rather than at the bottom TiO<sub>2</sub>/MAPbI<sub>3</sub> interface,<sup>[19]</sup> and thus the operation of the device again suggests the charge transport in the MAPbI<sub>3</sub> rather than in TiO<sub>2</sub>.

We subsequently employed ferroelectric polymers PVDF-HFP and PVDF-TrFE as dielectric materials for the TGBC devices. As shown in **Figure 2d,e**, devices with PVDF-HFP and PVDF-TrFE dielectrics exhibited superior performance in terms of smaller hysteresis than the Cytop devices, which is consistent with previous findings that ferroelectric dielectric is beneficial to Pb-perovskite devices.<sup>[20]</sup> In particular, MAPbI<sub>3</sub> FETs with PVDF-TrFE dielectrics demonstrated nearly hysteresis-free operation, with on/off ratios approaching 10<sup>4</sup> and low turn-on voltages.



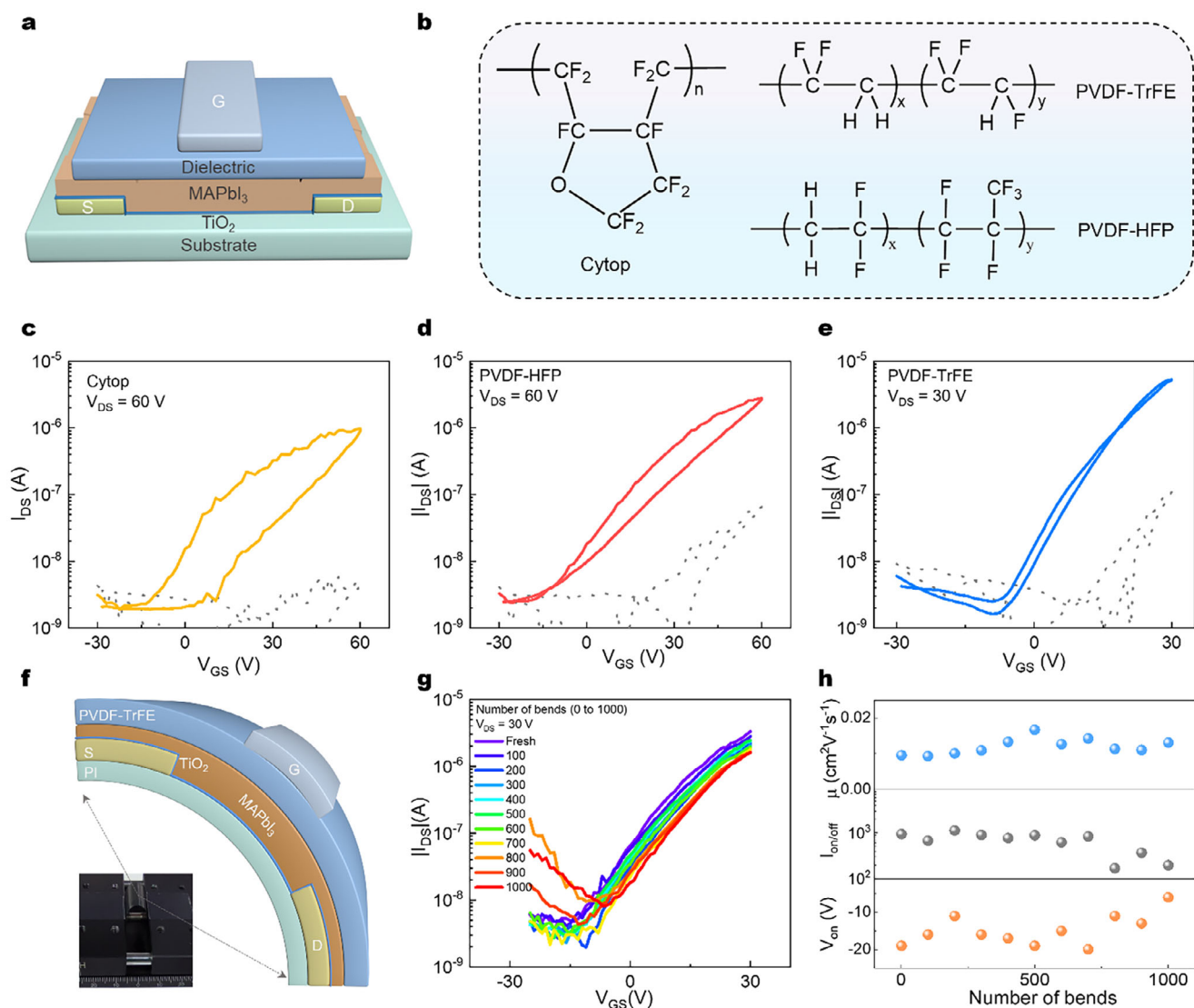
**Figure 1.** Enhanced device performance via TiO<sub>2</sub> interlayer engineering. Schematic of the MAPbI<sub>3</sub> FET device structure a) without and b) with TiO<sub>2</sub> interlayer. The c) transfer and d) output characteristic curves of the MAPbI<sub>3</sub> FET without TiO<sub>2</sub> interlayer. The e) transfer and f) output characteristic curves of the MAPbI<sub>3</sub> FET with TiO<sub>2</sub> interlayer.

Notably, the TGBC MAPbI<sub>3</sub> FETs with PVDF-TrFE dielectrics also display decent bias-stress stability: on-currents exhibit only a slight decrease with prolonged bias stress (Figure S5, Supporting Information).

Besides, we replaced the glass substrate with a PI substrate to fabricate a flexible TGBC MAPbI<sub>3</sub> FET device (Figure 2f), marking the first room-temperature operational flexible Pb-based MHP FET. After 1000 bending cycles, the transfer characteristics remain largely unchanged, with only a modest increase in

off-state current appearing after  $\approx 700$  cycles, likely due to structural damage in the PVDF-TrFE dielectric layer during bending. The stable trends in on/off ratio, threshold voltage, and mobility after 1000 cycles indicate excellent device stability (Figure 2g,h).

Following the development of this high-performance, room-temperature PVDF-TrFE-based MAPbI<sub>3</sub> FET, we evaluated its photodetection performance at 660 nm. Figure 3a displays the transfer characteristics under 660 nm illumination at varying optical power densities. The device exhibits pronounced photo-

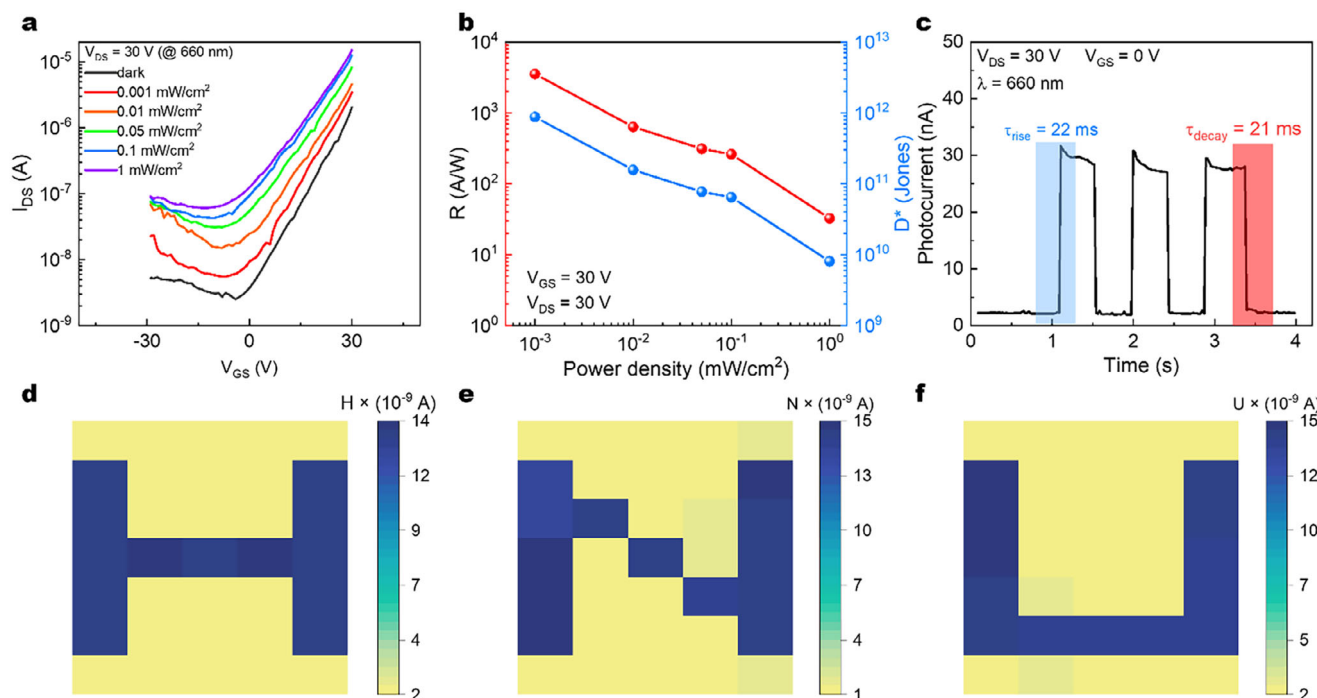


**Figure 2.** Top-gate MAPbI<sub>3</sub> FETs fabricated via TiO<sub>2</sub> interlayer strategy. a) Schematic diagram of the top gate structure MAPbI<sub>3</sub> FET through TiO<sub>2</sub> interlayer. b) Chemical structures of the three dielectric materials: Cytop, PVDF-HFP, and PVDF-TrFE. Transfer characteristics of top gate MAPbI<sub>3</sub> FETs using c) Cytop, d) PVDF-HFP, and e) PVDF-TrFE as dielectric layers, respectively. f) Schematic diagram of the flexible MAPbI<sub>3</sub> FET structure. Inset shows the illustration of the bending test for the flexible device. The bending radius is  $\approx 7$  mm. g) Changes in the transfer characteristic curves and h) corresponding performance parameters of the flexible device after 1000 bending cycles.

gating and photoconductive effects even at a low light intensity of  $0.001 \text{ mW cm}^{-2}$ . Figure 3b illustrates the dependence of responsivity (R) and specific detectivity ( $D^*$ ) on illumination intensity at gate voltage of 30 V. The peak values reach  $3 \times 10^3 \text{ A/W}$  for R and  $10^{12}$  Jones for  $D^*$ , demonstrating exceptional red-light detection capabilities. The dynamic response was quantified by measuring the photocurrent rise and decay times, yielding values of 22 ms and 21 ms, respectively (Figure 3c). To demonstrate practical imaging potential, we fabricated a device array consisting of  $5 \times 5$  devices. Leveraging the remarkable of the device to 660 nm light, the MAPbI<sub>3</sub> FET array successfully discriminates the letters “H,” “N,” and “U” with high contrast (Figure 3d,e), highlighting its promising application in photo imaging technologies.

### 2.3. Understanding the Role of TiO<sub>2</sub> Interlayer

To elucidate the beneficial role of the TiO<sub>2</sub> interlayer, we conducted comprehensive morphological and electronic characterizations. As shown in Figure 4a, when depositing MAPbI<sub>3</sub> precursor solution onto substrates, the substrates with the TiO<sub>2</sub> interlayer exhibited significantly smaller contact angles compared to bare substrates, which is crucial for forming uniform and compact perovskite films.<sup>[21,22]</sup> This conclusion was further corroborated by the SEM images in Figure 4b,c. The MAPbI<sub>3</sub> film without TiO<sub>2</sub> interlayer displayed relatively small grains ( $\approx 100$  nm) along with numerous particulates, indicative of poor film quality. In contrast, the film with TiO<sub>2</sub> interlayer showed larger, more homogeneous grains ( $\approx 200$  nm) with improved compactness and



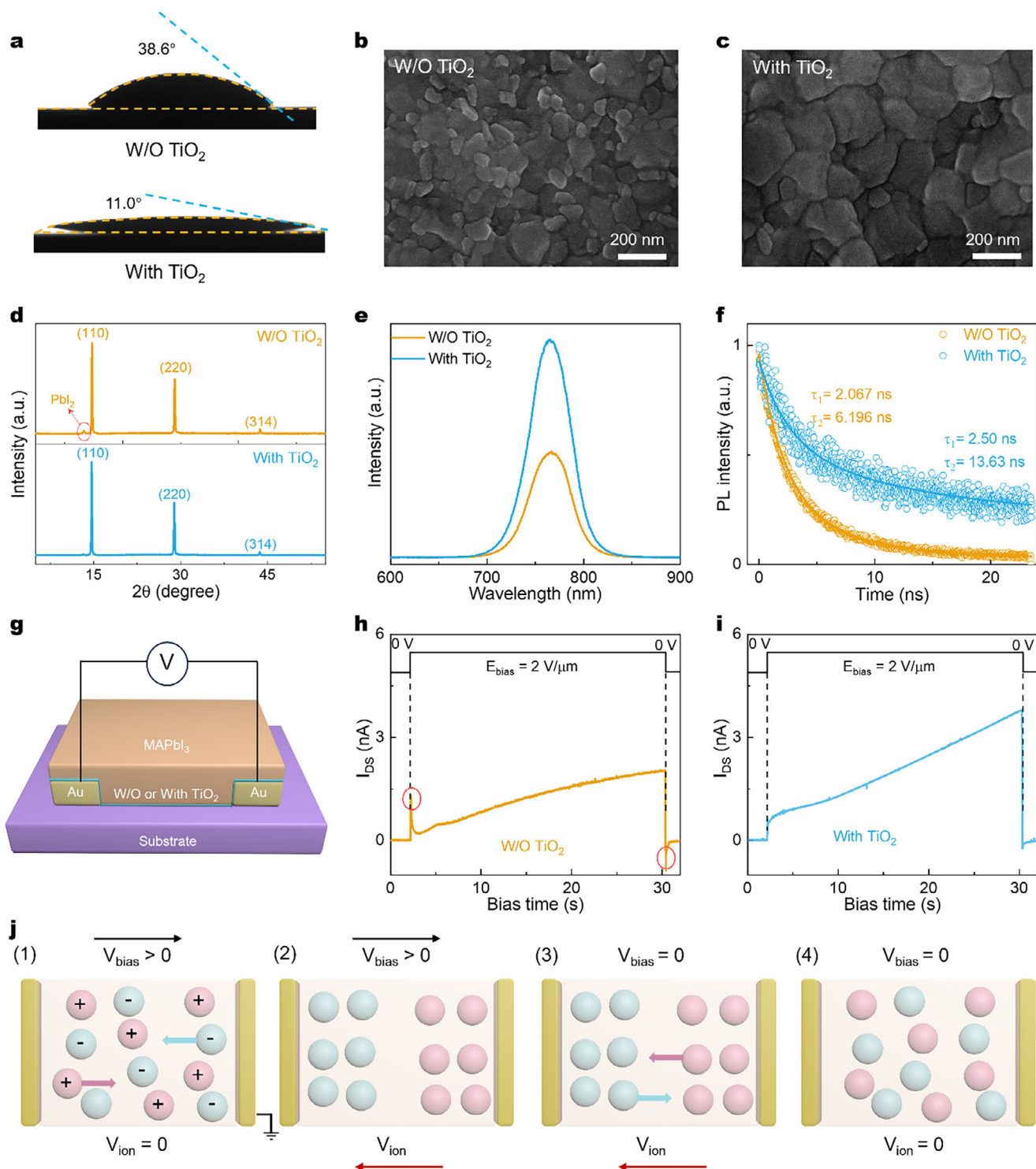
**Figure 3.** Photodetection performance of the TGBC MAPbI<sub>3</sub> FET with TiO<sub>2</sub> interlayer. a) Transfer characteristics of interlayer based MAPbI<sub>3</sub> device under various illumination power intensities ( $\lambda = 660$  nm,  $V_{DS} = 30$  V). b) Responsivity and detectivity values of the interlayer-based MAPbI<sub>3</sub> device as a function of power intensity at  $V_{GS}$  of 30 V. c) The time-dependent photocurrent measurements of the interlayer-based MAPbI<sub>3</sub> device under 660 nm light illumination. d) “H”, e) “N”, and f) “U” shaped mask (660 nm, power density = 0.01 mW cm<sup>-2</sup>,  $V_{GS} = 0$  V).

uniformity. The XRD patterns in Figure 4d reveal a weak PbI<sub>2</sub> diffraction signal at 13° for the MAPbI<sub>3</sub> film without TiO<sub>2</sub> interlayer, indicating a slight incomplete conversion of the perovskite precursor.<sup>[23]</sup> In contrast, the MAPbI<sub>3</sub> film with TiO<sub>2</sub> interlayer shows no such signal, and meanwhile, the full width at half maximum of the (110) diffraction peak becomes smaller (Figure S6, Supporting Information), suggesting that the TiO<sub>2</sub> interlayer effectively enhances the crystallinity of MAPbI<sub>3</sub> film. Figure 4e,f present the steady-state and transient PL spectra, respectively. The MAPbI<sub>3</sub> film with TiO<sub>2</sub> interlayer exhibits a higher PL intensity at the characteristic peak of 765 nm, suggesting lower defect density.<sup>[24]</sup> Furthermore, the transient PL decay times ( $\tau_1/\tau_2$ ) are 2.50 and 13.63 ns, respectively, which are longer than those of the MAPbI<sub>3</sub> film without TiO<sub>2</sub> interlayer, confirming effective suppression of non-radiative recombination.<sup>[25]</sup> These results collectively demonstrate that the incorporation of a TiO<sub>2</sub> interlayer not only improves the crystallographic quality of MAPbI<sub>3</sub> films but also reduces defect density and suppresses non-radiative recombination.

To further understand the electronic impact of the interlayer, we performed UPS measurements to determine the energy levels (Figure S7, Supporting Information). We observed two critical changes in the MAPbI<sub>3</sub> film with the TiO<sub>2</sub> interlayer. First, both the valence and conduction bands undergo a downward shift of  $\approx 0.34$  eV. This is likely due to the formation of an interface dipole at the TiO<sub>2</sub>/MAPbI<sub>3</sub> interface, which alters the local electrostatic potential. Second, the Fermi level shifts upward by 0.15 eV (from -4.20 to -4.05 eV), moving significantly closer to the conduction band. This upward shift is a clear signature of an n-doping ef-

fect, likely caused by electron transfer from the TiO<sub>2</sub> interlayer, which is known to have electron-donating oxygen vacancies. The n-doping is highly beneficial for N-type FET performance, as the increased carrier concentration enhances channel conductivity and can passivate electron trap states. Concurrently, the downward shift of the MAPbI<sub>3</sub> energy bands can reduce the electron injection barrier from the source/drain electrodes (Figure S7d, Supporting Information). These electronic benefits, synergistic with the improved film quality, collectively boost device performance.

Furthermore, we investigated the ion migration in the two films. In MHP FETs, ion migration within the lattice and along grain boundaries becomes particularly pronounced under high electric fields or elevated temperatures. Migrating ions accumulate near the gate, forming local electric fields that partially counteract the externally applied gate field, leading to significant hysteresis or even loss of gate modulation in the transfer characteristics of the FET devices.<sup>[26]</sup> Huang’s group previously demonstrated that ions within MHPs migrate under external electric fields ( $E$ ) exceeding  $1 \text{ V } \mu\text{m}^{-1}$ .<sup>[27,28]</sup> To investigate the ion migration in the perovskite films, we deposited MAPbI<sub>3</sub> films on substrates, followed by thermal evaporation of Au electrodes with a channel length of 20  $\mu\text{m}$ . A bias voltage of 40 V ( $E = 2 \text{ V}/\mu\text{m}$ ) was applied across the Au electrodes, and the current-voltage relationship over time was recorded to evaluate ion migration (Figure 4g). As shown in Figure 4h, for the MAPbI<sub>3</sub> device without a TiO<sub>2</sub> interlayer, the current initially increases upon bias application ( $t = 2$  s), then decreases significantly before slowly increasing again, forming a positive current peak. Upon bias removal ( $t = 30.4$



**Figure 4.** Impact of  $\text{TiO}_2$  on  $\text{MAPbI}_3$  films. a) Contact-angle measurements of the  $\text{MAPbI}_3$  precursor on substrates without and with the ultrathin  $\text{TiO}_2$  interlayer. b, c) SEM micrographs showing surface morphology of  $\text{MAPbI}_3$  films prepared (b) directly on the bare substrate and (c) on the  $\text{TiO}_2$ -coated substrate. d) XRD patterns comparing the crystallinity of  $\text{MAPbI}_3$  films with and without the  $\text{TiO}_2$  interlayer. e) Steady-state PL spectra and f) time-resolved PL decay traces for  $\text{MAPbI}_3$  films with and without the  $\text{TiO}_2$  interlayer, indicating improved optoelectronic quality in the presence of  $\text{TiO}_2$ . g) Schematic of the test structure used for ion-migration measurements. Current-time traces recorded under a constant bias of 40 V for  $\text{MAPbI}_3$  devices (h) without and i) with the  $\text{TiO}_2$  interlayer, showing pronounced ion-driven transient features in the untreated film and a fast, predominantly electronic response for the  $\text{TiO}_2$ -modified film. j) Schematic illustration of the dynamic ion behavior in  $\text{MAPbI}_3$  under an applied electric field: mobile ions migrate to the electrodes during bias, producing an internal ion-induced potential that opposes the external field and gives rise to transient current peaks at bias application and removal.

s), a reverse current peak is observed before the current gradually decays to zero, indicating a dynamic relaxation response. In contrast, for the MAPbI<sub>3</sub> device with TiO<sub>2</sub> interlayer, the current increases immediately upon bias application and returns to its original level instantly upon bias removal, demonstrating a rapid voltage-current response consistent with previous reports (Figure 4i).<sup>[29]</sup>

The dynamic relaxation behavior of the MAPbI<sub>3</sub> device can be explained using the schematic shown in Figure 4j. The two gold blocks represent the top electrodes, and the positively and negatively charged spheres represent mobile cations and anions within the MAPbI<sub>3</sub>, respectively. Initially, in the absence of an external electric field ( $V_{\text{bias}} = 0$  V), the ions are randomly distributed between the electrodes, maintaining electrical neutrality, and the internal ion-induced voltage is zero ( $V_{\text{ion}} = 0$  V). Upon application of an external bias ( $V_{\text{bias}} > 0$  V), ions migrate toward the perovskite/electrode interface under the electric field, inducing an internal ion-induced voltage opposite to the external bias. The  $V_{\text{ion}}$  potential partially counteracts the external bias voltage ( $V_{\text{bias}}$ ), leading to a reduction in current flow. This corresponds to the initial signal peak observed in Figure 4h. Upon removal of the external bias, the accumulated ions relax back toward equilibrium, producing a transient reverse current that decays as ionic concentration gradients dissipate (Figure 4h).<sup>[29,30]</sup> The positive and negative current peaks observed at bias application and removal are therefore signatures of ion migration in MAPbI<sub>3</sub>; the absence of such peaks in devices with a TiO<sub>2</sub> interlayer indicates that ion migration is effectively suppressed by the interlayer.

In summary, the TiO<sub>2</sub> interlayer initiates a cascade of improvements that fundamentally enhance device performance and stability. It first improves wetting of the perovskite precursor, which promotes larger, more uniform grains and improved film compactness, leading to better crystallinity and reduced defect density. Electronically, the interlayer introduces a beneficial n-doping effect that increases carrier concentration and simultaneously causes a favorable shift in the energy band alignment through interface dipole formation. Finally, the combination of improved crystallinity and a more compact microstructure substantially inhibits ion migration under an applied electric field. Together, these structural, electronic, and ionic improvements account for the markedly superior performance and electrical stability of the MAPbI<sub>3</sub> FETs incorporating the TiO<sub>2</sub> interlayer.

#### 2.4. Universality of the Interlayer Strategy

After confirming the role of the ultrathin TiO<sub>2</sub> interlayer, we explored the universality of this strategy. Recent advances in N-type MHP FETs have predominantly focused on 3D MHPs and 2D Ruddlesden-Popper (RP) phases<sup>[31,32]</sup>; N-type DJ phase MHP FETs remain unreported despite DJ phases offering enhanced environmental stability relative to 3D MHPs and superior out-of-plane charge transport compared with 2D RP perovskites.<sup>[24]</sup> Motivated by our success with MAPbI<sub>3</sub>, we extended the interlayer approach to 2D DJ-phase perovskite PDAPbI<sub>4</sub> (Figure 5a) and obtained room-temperature functional FETs.

Figure 5b shows the XRD of solution-processed PDAPbI<sub>4</sub> films, exhibiting sharp reflections at 8.6°, 14.1°, and 27.4° assignable to the (002), (210), and (006) planes, respectively.<sup>[33]</sup>

Atomic force microscopy (AFM) revealed uniformly distributed grains of  $\approx 100$  nm with dense, continuous morphology (Figure 5c). Optical absorption spectra showed a sharp edge at 450 nm, corresponding to a bandgap of 2.56 eV (Figure 5d). Notably, PDAPbI<sub>4</sub> demonstrates significantly superior environmental stability compared to MAPbI<sub>3</sub>. As evidenced in Figure 5e, while MAPbI<sub>3</sub> films underwent complete degradation (manifested as black-to-yellow transition) within 24 h under ambient conditions, PDAPbI<sub>4</sub> maintained its colors throughout the same period. This enhanced stability originates from the hydrophobic organic cations in the DJ structure that protect the inorganic sheets from moisture and other environmental agents.<sup>[24,34]</sup>

Using identical device fabrication protocols, we converted these films into FETs. PDAPbI<sub>4</sub> devices deposited on bare substrates could not work, exhibiting source-drain currents below 10<sup>-9</sup> A throughout gate voltage sweeps ( $\pm 60$  V). Remarkably, incorporation of the ultrathin TiO<sub>2</sub> interlayer enabled robust N-type FETs, with an on/off ratio greater than 10<sup>3</sup> (Figure 5f) and output curves showing clear gate modulation (Figure 5g). Comparative stability tests revealed PDAPbI<sub>4</sub> FETs maintained performance after 20 days in an Ar-filled glovebox, while MAPbI<sub>3</sub> FETs degraded substantially within just 6 days, highlighting the superior potential of 2D DJ-phase perovskites (Figure S8, Supporting Information).

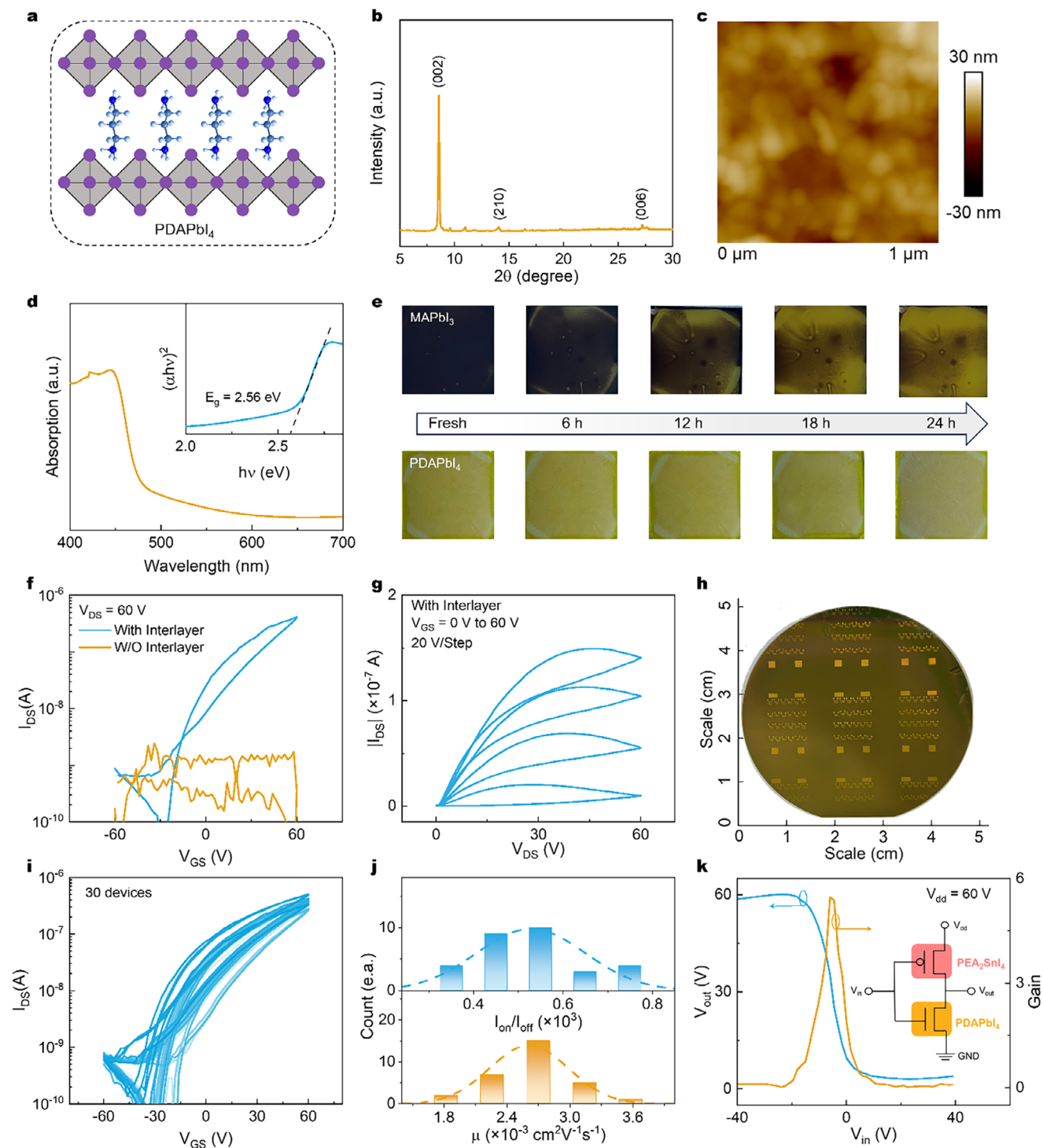
To demonstrate scalability and device uniformity, we fabricated a 2-inch wafer-scale PDAPbI<sub>4</sub> FET array using our anti-solvent-free, single-step spin-coating approach (Figure 5h). Statistical analysis of 30 devices showed uniform performance: on-state currents and threshold voltages remained stable, while mobility and on/off ratios followed normal distributions with peak values reaching  $3.6 \times 10^{-3}$  cm<sup>2</sup>V<sup>-1</sup>s<sup>-1</sup> and  $7.5 \times 10^2$ , respectively (Figure 5i,j). The consistent yield and stability of these N-type devices provide a practical basis for realizing large-area electronics.

As a proof-of-concept for integrated logic, we combined our TiO<sub>2</sub>-enabled N-type PDAPbI<sub>4</sub> FETs with P-type PEA<sub>2</sub>SnI<sub>4</sub> FETs (Figure S9, Supporting Information) to construct an all-MHP inverter. The inverter shows clear voltage inversion and a peak gain approaching 6 (Figure 5k), demonstrating the feasibility of perovskite-only logic elements.

Finally, we extended the interlayer strategy to additional DJ-phase perovskites by alloying MA<sup>+</sup> with PDA<sup>2+</sup> to form PDAMAPb<sub>2</sub>I<sub>7</sub> and PDAMA<sub>2</sub>Pb<sub>3</sub>I<sub>10</sub>. Films processed without the TiO<sub>2</sub> interlayer again yielded negligible transistor currents ( $I_{\text{DS}} < 10^{-9}$  A) under gate voltage sweeps (Figure S10, Supporting Information). In stark contrast, devices with the interlayer demonstrated standard N-type FETs behavior, featuring a near-zero threshold voltage, an on/off ratio of 10<sup>3</sup>, and well-modulated output characteristics (Figure S11, Supporting Information). These results confirm that the ultrathin TiO<sub>2</sub> interlayer is an effective, generalizable strategy to enable and stabilize N-type transport across a range of Pb-based perovskites.

### 3. Conclusion

In conclusion, this work introduces a simple yet powerful interlayer strategy that fundamentally addresses critical challenges in Pb-based perovskite FETs. By strategically introducing an ultrathin TiO<sub>2</sub> interlayer, we have demonstrated a reproducible approach to fabricating high-performance N-type MHP FETs that



**Figure 5.** Fabrication and characterization of Dion-Jacobson phase perovskite FETs. a) The chemical structure and b) XRD pattern of PDAPbI<sub>4</sub> film. c) AFM image and d) UV-vis absorbance of PDAPbI<sub>4</sub> film. e) Photographs showing the evolution of MAPbI<sub>3</sub> and PDAPbI<sub>4</sub> films in the atmospheric environment with time. f) Transfer and g) output characteristics of PDAPbI<sub>4</sub> FET. h) The photo of the 2-inch PDAPbI<sub>4</sub> FET arrays. i) Transfer curves of 30 randomly selected devices from a large-area PDAPbI<sub>4</sub> FET array prepared. j) Statistical distribution of the mobility and on/off ratio of the 30-device array. k) The output voltage and gain as a function of input voltage for the all-perovskite inverter.

overcomes longstanding limitations such as poor film quality, high defect densities, and significant ion migration. Our comprehensive investigations reveal that the TiO<sub>2</sub> interlayer enhances precursor wetting, promotes larger and more uniform grain formation, reduces defect density, and effectively suppresses non-radiative recombination and ion migration. Beyond MAPbI<sub>3</sub>, we have successfully extended this strategy to 2D Dion-Jacobson phase perovskites, including PDAPbI<sub>4</sub> and its derivatives, highlighting the universality of our approach. This work not only provides a straightforward fabrication method for high-performance Pb-based perovskite FETs but also offers fundamental insights into improving the electronic properties of metal halide perovskite semiconductors.

## 4. Experimental Section

**Materials:** Methylammonium iodide (MAI, ≥ 99.5%), propane-1,3-diammonium (II) iodide (PDAl<sub>2</sub>, ≥ 99.99%), and Lead (II) iodide (PbI<sub>2</sub>, ≥ 99.99%) were purchased from Xi'an Polymer Light Technology Corp. Poly(vinylidene fluoride-trifluoroethylene) (P(VDF-TrFE), 70:30 mol%) and poly(vinylidene fluoride-co-hexafluoropropylene) (PVDF-HFP) were purchased from Piezotech. Lead acetate trihydrate (PbAc<sub>2</sub>, 99.998%), n-butyl Acetate (nBA, 99.5%), 2-butanone (MEK, 98%), and N, N-dimethylformamide (DMF, anhydrous, 99.8%) were bought from Sigma-Aldrich. Cytop solution was prepared by adding Cytop solvent into Cytop with a volume ratio of 1:3. TiO<sub>2</sub> solution was obtained by mixing TiO<sub>2</sub> colloidal and anhydrous ethanolin. All materials were used as received without any additional purification.

**Preparation of MAPbI<sub>3</sub> Films:** For the preparation of MAPbI<sub>3</sub> precursor solutions, two distinct methods were employed. In the PbI<sub>2</sub> method, MAPbI<sub>3</sub> was synthesized from a precursor solution containing 0.75 M PbI<sub>2</sub> and MAI in a 1:1 molar ratio dissolved in DMF. In the PbAc<sub>2</sub> method, MAPbI<sub>3</sub> was formed by mixing 0.75 M PbAc<sub>2</sub> and MAI in a 1:3 molar ratio in DMF. Both precursor solutions were heated at 45 °C for 10 h within an Ar-filled glovebox to ensure complete dissolution. Prior to use, the solutions were filtered through 0.45 μm PTFE filters to remove any undissolved particulates. For film fabrication, the PbI<sub>2</sub> method involved spin-coating the precursor solution at 4000 rpm for 30 s. During the spin-coating process, a droplet of chlorobenzene (CB) was applied to the sample surface at the 5-s mark to promote uniform crystallization. In the PbAc<sub>2</sub> method, the precursor solution was spin-coated directly onto the substrates at 4000 rpm for 30 s without the addition of chlorobenzene. All perovskite films were subsequently annealed at 150 °C for 2 min in the Ar-filled glovebox to enhance crystallinity and film quality.

**Preparation of PDAMA<sub>n-1</sub>Pb<sub>n</sub>I<sub>3n+1</sub> Films:** The preparation process of PDAMA<sub>n-1</sub>Pb<sub>n</sub>I<sub>3n+1</sub> (n = 1–3) is as follows: PbI<sub>2</sub>, MAI, and PDAl<sub>2</sub> are weighed according to a specific molar ratio and dissolved in DMSO solvent to prepare a precursor solution with a concentration of 0.05 m. The spin-coating process is carried out at 4000 rpm for 30 s, and no anti-solvent is added during the spin-coating process. Finally, the film is annealed at 150 °C for 1 min to complete crystallization and film optimization.

**Fabrication of FET:** For the BGBC devices, Cr/Au (3/32 nm) electrodes were first pre-patterned on Si/SiO<sub>2</sub> substrates using photolithography. These substrates were then sequentially cleaned by ultrasonication in ultrapure water, acetone, and isopropanol for 5 min each, followed by UV/ozone treatment for 15 min. A TiO<sub>2</sub> precursor solution was spin-coated onto the substrates at 4000 rpm for 30 s and annealed at 150 °C for 30 min in ambient conditions. Subsequently, the perovskite film was directly deposited via spin-coating inside a glovebox. For the MAPbI<sub>3</sub> TGBC devices, the substrate used for fabrication was replaced with glass slides. The preparation methods for TiO<sub>2</sub> and MAPbI<sub>3</sub> remained the same as those for the BGBC devices. Additionally, a dielectric layer was spin-coated on top of the MAPbI<sub>3</sub> film. The channel length and width of the device are 20 and 1000 μm, respectively. Finally, a 50 nm Al gate electrode

was deposited by thermal evaporation through a shadow mask to complete the device fabrication. For the dielectric layers, either P(VDF-TrFE) (70 mg mL<sup>-1</sup>, dissolved in n-butyl acetate, nBA) or PVDF-HFP (100 mg mL<sup>-1</sup>, dissolved in methyl ethyl ketone, MEK) was spin-coated onto the perovskite film using a two-step program: the first step at 500 rpm for 3 s and the second step at 4000 rpm for 20 s. Post-annealing at 90 °C for 1 h was performed to enhance the film quality. The spin-coating processes for Cytop followed conventional protocols. The device fabrication of the PDAMA<sub>n-1</sub>Pb<sub>n</sub>I<sub>3n+1</sub> FET with the BGBC structure is the same as the above. The channel length and width of the device used are 60 and 1000 μm, respectively.

**Characterization of Film:** The film morphology was measured by SEM (TESCAN MIRA3) and atomic force microscopy (AFM) (Bioscope system). The X-ray diffraction (XRD) patterns were recorded by D/max 2550 (Rigaku) under Cu Kα (λ = 1.5406 Å) irradiation in the ambient atmosphere. PL spectra were recorded by a Thermo Scientific Lumina. The ultraviolet-visible-near-infrared (UV-vis-NIR) absorption spectra of perovskite films were characterized with UV-3600 PLUS (Shimadzu).

**Measurement of FET:** Electrical properties of FET devices were performed with a Keysight B2912A Precision Sources through a probe station in an Ar glovebox. All measurement of FETs is carried out under dark conditions. It should be noted that all the measurement procedures were measured in DC mode rather than pulse mode.

**Characterization of MAPbI<sub>3</sub> Phototransistor:** During light detection, the light source was placed in a dark glovebox to avoid influence by ambient light. The light at 660 nm was emitted using commercial light sources. The intensity of the incident light at different wavelengths was calibrated by a commercial silicon photodetector. The responsivity (R) is defined as  $R = \frac{I_{\text{light}} - I_{\text{dark}}}{P \times S}$ , where I<sub>light</sub> and I<sub>dark</sub> represent the current under illumination and in dark conditions, respectively. P is the incident light power intensity, and S is the effective active area of the MAPbI<sub>3</sub> FET device. If one assumes that the dark current is mainly originates from shot noise of dark current (I<sub>dark</sub>), the specific detectivity (D\*) can be written as:  $D^* = \frac{R}{\sqrt{2qI_{\text{dark}}S}}$ .

**Statistical Analysis:** The performance data of the perovskite FETs for n devices (n = 30) were summarized and presented as a distribution graph without pre-processing. The mean or standard deviation values were not calculated. The data processing was performed by Origin software.

## Supporting Information

Supporting Information is available from the Wiley Online Library or from the author.

## Acknowledgements

The authors thank the Center for Analysis and Testing of Hunan University for providing the AFM. Y.H. thanks the National Key Research and Development Program (2021YFA1200700), the National Natural Science Foundation of China (62222403, U21A20497). The study was supported by the "111 Center" (B25033), the Foundation of Science and Technology on Advanced Ceramic Fibers and Composites Laboratory (Grant No. 6142907240402), and the Scientific Research Fund of Hunan Provincial Education Department (23B0842) for financial support.

## Conflict of Interest

The authors declare no conflict of interest.

## Data Availability Statement

The data that support the findings of this study are available from the corresponding author upon reasonable request.

## Keywords

dion-jacobson (DJ) phase perovskite, field-effect transistor, metal halide perovskites, photodetector, TiO<sub>2</sub> interlayer

Received: August 27, 2025

Revised: September 24, 2025

Published online: October 7, 2025

- [1] S. P. Senanayak, K. Dey, R. Shivanna, W. Li, D. Ghosh, Y. Zhang, B. Roose, S. J. Zelewski, Z. Andaji-Garmaroudi, W. Wood, N. Tiwale, J. L. MacManus-Driscoll, R. H. Friend, S. D. Stranks, H. Siringhaus, *Nat. Mater.* **2023**, 22, 216;
- [2] M. Lee Michael, J. Teuscher, T. Miyasaka, N. Murakami Takurou, J. Snaith Henry, *Science* **2012**, 338, 643;
- [3] Z. Huihui, Y. Wonryeol, R. Youjin, Z. Guanhaojie, B. Sai, L. Ao, N. Yong-Young, *Nat. Electron.* **2023**, 6, 650.
- [4] X. Yang, Y. Liu, S. Yang, Y. Wu, Y. Lei, Y. Yang, A. Liu, J. Chu, W. Li, *Adv. Funct. Mater.* **2024**, 34, 2403917.
- [5] Z. Zhou, Q. Li, M. Chen, X. Zheng, X. Wu, X. Lu, S. Tao, N. Zhao, *ACS Energy Lett.* **2023**, 8, 4496.
- [6] I. H. Chao, Y.-T. Yang, M.-H. Yu, C.-H. Chen, C.-H. Liao, B.-H. Lin, I. C. Ni, W.-C. Chen, A. W. Y. Ho-Baillie, C.-C. Chueh, *Small* **2023**, 19, 2207734.
- [7] W. Yang, K. Zhang, W. Yuan, L. Zhang, C. Qin, H. Wang, *Adv. Mater.* **2024**, 36, 2313461.
- [8] H. Zai, Y. Ma, Q. Chen, H. Zhou, *J. Energy Chem.* **2021**, 63, 528.
- [9] Y. Zhang, A. Ummadisingu, R. Shivanna, D. H. L. Tjhe, H.-I. Un, M. Xiao, R. H. Friend, S. P. Senanayak, H. Siringhaus, *Small* **2023**, 19, 2302494.
- [10] X.-J. She, C. Chen, G. Divitini, B. Zhao, Y. Li, J. Wang, J. F. Orri, L. Cui, W. Xu, J. Peng, S. Wang, A. Sadhanala, H. Siringhaus, *Nat. Electron.* **2020**, 3, 694.
- [11] P. Senanayak Satyaprasad, M. Abdi-Jalebi, S. Kamboj Varun, R. Carey, R. Shivanna, T. Tian, G. Schweicher, J. Wang, N. Giesbrecht, D. Di Nuzzo, E. Beere Harvey, P. Docampo, A. Ritchie David, D. Fairen-Jimenez, H. Friend Richard, H. Siringhaus, *Sci. Adv.* **2020**, 6, aaz4948.
- [12] P. Senanayak Satyaprasad, B. Yang, H. Thomas Tudor, N. Giesbrecht, W. Huang, E. Gann, B. Nair, K. Goedel, S. Guha, X. Moya, R. M. C. Christopher, P. Docampo, A. Sadhanala, H. F. Richard, H. Siringhaus, *Sci. Adv.* **2017**, 3, 1601935.
- [13] D. Li, H. C. Cheng, Y. Wang, Z. Zhao, G. Wang, H. Wu, Q. He, Y. Huang, X. Duan, *Adv. Mater.* **2017**, 29, 1601959.
- [14] D. Li, G. Wang, H.-C. Cheng, C.-Y. Chen, H. Wu, Y. Liu, Y. Huang, X. Duan, *Nat. Commun.* **2016**, 7, 11330.
- [15] A. Bruno, D. Cortecchia, X. Y. Chin, K. Fu, P. P. Boix, S. Mhaisalkar, C. Soci, *Adv. Energy Mater.* **2017**, 7, 1700265.
- [16] S. Chen, Y. Liu, X. Xiao, Z. Yu, Y. Deng, X. Dai, Z. Ni, J. Huang, *Joule* **2020**, 4, 2661.
- [17] W. Zhang, M. Saliba, D. T. Moore, S. K. Pathak, M. T. Horantner, T. Stergiopoulos, S. D. Stranks, G. E. Eperon, J. A. Alexander-Webber, A. Abate, A. Sadhanala, S. Yao, Y. Chen, R. H. Friend, L. A. Estroff, U. Wiesner, H. J. Snaith, *Nat. Commun.* **2015**, 6, 6142.
- [18] J. Shao, S. Yang, Y. Liu, *ACS Appl. Mater. Interfaces* **2017**, 9, 16202.
- [19] J. Zaumseil, H. Siringhaus, *Chem. Rev.* **2007**, 107, 1296.
- [20] J. Xia, X. Qiu, Y. Liu, P.-A. Chen, J. Guo, H. Wei, J. Ding, H. Xie, Y. Lv, F. Li, W. Li, L. Liao, Y. Hu, *Adv. Sci.* **2023**, 10, 2300133.
- [21] X. Tong, L. Xie, J. Li, Z. Pu, S. Du, M. Yang, Y. Gao, M. He, S. Wu, Y. Mai, Z. Ge, *Adv. Mater.* **2024**, 36, 2407032;
- [22] W. Shen, U. Jung, Z. Xian, B. Jung, J. Park, *J. Alloys Compd.* **2022**, 929, 167329.
- [23] Y. Zhou, A. Najjar, J. Zhang, J. Feng, Y. Cao, Z. Li, X. Zhu, D. Yang, S. F. Liu, *ACS Appl. Mater. Interfaces* **2022**, 14, 28729.
- [24] X. Qiu, J. Xia, Y. Liu, P.-A. Chen, L. Huang, H. Wei, J. Ding, Z. Gong, X. Zeng, C. Peng, C. Chen, X. Wang, L. Jiang, L. Liao, Y. Hu, *Adv. Mater.* **2023**, 35, 2305648.
- [25] S. Sidhik, I. Metcalf, W. Li, T. Kodalle, C. J. Dolan, M. Khalili, J. Hou, F. Mandani, A. Torma, H. Zhang, R. Garai, J. Persaud, A. Marciel, I. A. Muro Puente, G. N. M. Reddy, A. Balvanz, M. A. Alam, C. Katan, E. Tsai, D. Ginger, D. P. Fenning, M. G. Kanatzidis, C. M. Sutter-Fella, J. Even, A. D. Mohite, *Science* **2024**, 384, 6701.
- [26] X. Qiu, Y. Liu, J. Xia, J. Guo, P.-A. Chen, H. Wei, J. Guo, X. Shi, C. Chen, Z. Zeng, H. Chen, L. Jiang, L. Liao, Y. Hu, *Cell Reports Physical Science* **2023**, 4, 101217.
- [27] Z. Xiao, Y. Yuan, Y. Shao, Q. Wang, Q. Dong, C. Bi, P. Sharma, A. Gruverman, J. Huang, *Nat. Mater.* **2015**, 14, 193.
- [28] Z. Xiao, Q. Dong, C. Bi, Y. Shao, Y. Yuan, J. Huang, *Adv. Mater.* **2014**, 26, 6503.
- [29] S. Zhou, G. Zhou, Y. Li, X. Xu, Y.-J. Hsu, J. Xu, N. Zhao, X. Lu, *ACS Energy Lett.* **2020**, 5, 2614.
- [30] D. Li, H. Wu, H.-C. Cheng, G. Wang, Y. Huang, X. Duan, *ACS Nano* **2016**, 10, 6933.
- [31] Y. Liu, P.-A. Chen, Y. Hu, *J. Mater. Chem. C* **2020**, 8, 16691.
- [32] X. Qiu, Y. Liu, W. Li, Y. Hu, *Nanoscale* **2020**, 12, 22425.
- [33] F. Zhang, S. Y. Park, C. Yao, H. Lu, S. P. Dunfield, C. Xiao, S. Uličná, X. Zhao, L. Du Hill, X. Chen, X. Wang, L. E. Mundt, K. H. Stone, L. T. Schelhas, G. Teeter, S. Parkin, E. L. Ratcliff, Y.-L. Loo, J. J. Berry, M. C. Beard, Y. Yan, B. W. Larson, K. Zhu, *Science* **2022**, 375, 6576.
- [34] W. Park, Y. Reo, W. Yang, H. Choi, S. Jeon, B. Lim, A. Liu, H. Zhu, Y.-Y. Noh, *ACS Energy Lett.* **2024**, 9, 2436.