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Ultrasensitive self-powered UV PDs via depolarization and

heterojunction fields jointly enhanced carriers separation

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Abstract

In current semiconductor-based ultraviolet (UV) self-powered photodetectors (PDs), because the built-in electric field only exit in the space-charge region of p-n junction or schottky junction interfaces, the photo-generated carriers can't be separate efficiently and thus the photoelectric response is low. To solve this issue, in this study, we construct a new-type Pb,La(Zr,Ti)O₃ ferroelectrics/TiO₂ semiconductors heterojunction self-powered UV photodetector. By optimizing reasonably the TiO₂ layer number, the PDs with 2-layer TiO₂ thin film exhibit excellent self-powered ultraviolet photoelectric response performances with a large responsivity of 21 mA/W, high specific detectivity of 2.4×10^{11} Jones, and fast response speed (rise time of 2.4 ms, decay time of 3.2 ms) under a negative poling electric voltage. These values are far superior to those of all recently reported ferroelectric and semiconductor-based UV self-powered PDs. The huge improvement of the photoelectric response is primarily because the depolarization field cause by the high remnant polarization of Pb,La(Zr,Ti)O₃ and the built-in electric field resulting from the heterojunction interfaces of the Pb,La(Zr,Ti)O₃ and the TiO₂ thin films enhance jointly the separation efficiency of photo-generated electron and holes. More interestingly, under weak-light illumination (0.18 mW/cm²), the PDs exhibit a ultrahigh responsivity of 91.4 mA/W and large specific detectivity of 9.7×10^{11} Jones, which indicate its superior weak UV light detection ability. This work opens new avenue for designing and fabricating next-generation self-powered UV PDs by a simple, low-cost and easy to mass-produce way.

Keywords: Pb,La(Zr,Ti)O₃ ferroelectrics; TiO₂ semiconductors; Heterojunction; Self-powered; Ultraviolet photodetectors

1. INTRODUCTION

Ultraviolet (UV) photodetectors (PDs), which are prepared with wide band-gap materials (such as TiO₂, ZnO, SnO₂, Ga₂O₃ and GaN) [1-7], have been widely used in fire warning, wireless communication, missile tracking, environmental protection, remote control, and ozone hole detecting due to their superior photoelectric response properties and chemical thermal stability [8,9]. However, in order to convert UV light to electrical signal, these semiconductor-based UV PDs usually require an external driving force for the separation and transport of photo-generated carriers [10]. This causes a large amount of energy consumption and restricts the device miniaturization

and integration. Therefore, self-powered UV PDs drove by the built-in electric field and zero operating voltage have recently drawn more and more attention. At present, the built-in electric field is usually obtained by fabricating the junctions at the interface of two different materials, such as p-n junction between p- and n-type semiconductors, and schottky junction between semiconductors and metals. Nevertheless, these kinds of built-in electric fields exist only in the space-charge region of the junction interfaces, leading to low separation efficiency of photo-generated electrons and holes and thus weak photoelectric response. Therefore, the development of self-powered UV PDs fabricated by new materials is of great importance.

Different from conventional wide band-gap materials, ferroelectrics have large remnant polarization and thus high depolarization electric field throughout the whole bulk. This can cause efficient separation of photo-generated electron and hole even at zero bias and thus superior photoelectrical response. Besides, the wide band gap (> 3 eV) of ferroelectric materials is appropriate for preparing UV detectors. Based on above these advantages, ferroelectric based self-powered UV PDs have aroused extensive research interest [11-14]. Lately, we have fabricated Pb,La(Zr,Ti)O₃ ferroelectric based UV PDs, which exhibit the most superior photoelectric response performances in current ferroelectric based self-powered UV photodetectors attributed to large remnant polarization and thus high depolarized electric field of Pb,La(Zr,Ti)O₃ materials. However, due to the low carrier concentration of the ferroelectric materials, the photocurrent of this kind of photodetector is low, which hence requires high-precision measurement systems to detect the signal. Therefore, there is a significant and urgent need to explore new strategy for high-performance ferroelectric based self-powered UV PDs.

Recently, ferroelectric/semiconductor heterostructures have drawn increasing attention because they takes the advantages of both the semiconductor and the ferroelectric properties and exhibited promising applications in photoelectric devices [15-17]. In heterojunction devices, electron and hole pairs are created inside both the semiconductor and ferroelectric layers and thus high carrier concentration can be obtained. More importantly, the depolarization field throughout the ferroelectrics and the built-in electric field resulting from the ferroelectric/semiconductor heterojunctions enhance jointly the separation efficiency of photo-generated carriers [15,18]. These two advantages can greatly cause the improvement of the photoelectric response for the photoelectric devices. For example, through the insertion of a 0.5Ba(Zr_{0.2}Ti_{0.8})O₃-0.5(Ba_{0.7}Ca_{0.3})TiO₃ ferroelectric layer between n-Si and p-SnO_x, the values of short-circuit photocurrent density (J_{sc}) , open-circuit voltage (V_{oc}) , fill factor (FF) and power conversion efficiency (PCE) can improve from 12.6 mAcm⁻², 0.23 V, 27% and 8.3 % to 30.9 mAcm⁻², -2.0 V, 19% and 10.9 %, respectively [19]. However, until now, it has been still scarce to manipulate the ferroelectric effect for improving the photoelectric response of ferroelectric/semiconductor heterojunction self-powered UV PDs.

As stated above, we have already confirmed that Pb,La(Zr,Ti)O₃ ferroelectric thin films are promising materials for fabricating self-powered photoelectric devices

due to their large remnant polarization and hence high depolarization electric field, which can result in efficient and fast separation of the photo-generated electrons and holes [13,20-22]. Besides, TiO₂ semiconductor material has been regard as a good candidate for UV PDs due to its wide band-gap (~3.2 eV), high chemical and optical stability, non toxicity and low production cost. Based on above two factors, in this study, we design and fabricate Au/Pb_{0.93}La_{0.07}(Zr_{0.2}Ti_{0.8})_{0.9825}O₃ (PLZT)/TiO₂/FTO heterojunction self-powered UV PDs and discuss the internal mechanism of the influence of the TiO₂ layer numbers and the poling electric field on the photoelectric response of the device. At zero bias condition, the PD with 2-layers TiO₂ and -0.5 V poling electric voltage exhibits excellent photoelectric response performances with large responsivity of 21 mA/W, high detectivity of 2.4×10^{11} Jones, and fast response speed (rise time of 2.4 ms, decay time of 3.2 ms). These values are far superior to those of all recently reported self-powered UV detectors. This work provide a new way for designing and fabricating self-powered UV photodetectors with superior photoelectric response performances.

2. EXPERIMENTAL PROCEDURE

2.1. Fabrication of PLZT/TiO₂ heterojunction self-powered UV PDs

The TiO₂ precursor solutions were prepared via a sol-gel method. A certain amount of tetrabutyl titanate $[Ti(OC_4H_9)_4]$ were added into the mixed solution of acetylacetone (C₅H₈O₂), glacial acetic acid (CH₃COOH) and ethylene glycol $[(CH_2OH)_2]$ with stirring at room temperature for 2 h to obtain a clear and transparent yellow solution, and then aged for 24 h to form a transparent sol. In the following, the sol was smeared on FTO substrates by spin-coating way at the rate of 3000 rpm for 30 s to ensure the uniformity and the above processes were respectively repeated 1, 2, 3 times to prepare TiO₂ thin film with multiple thickness. Then, those films were annealed 600°C for 2 h crystallization. After at in air for that, $Pb_{0.93}La_{0.07}(Zr_{0.2}Ti_{0.8})_{0.9825}O_3$ thin films were prepared on the TiO₂ films with different layer numbers by a spin-coating method, as reported in our previous work [20,21]. Finally, translucent square Au electrodes with a side length of 0.05 cm and a thickness of approximately 50 nm were deposited on the top surface of the PLZT thin films by thermal evaporation to fabricate the Au/PLZT/TiO₂/FTO heterojunction self-powered UV PDs.

2.2. The material and device characterization

The crystal structure of the PLZT/TiO₂ heterojunction thin films was characterized by X-ray diffraction (XRD, D8 Advance, Bruker, Germany). The surface and cross-sectional images were examined via field-emission scanning electron microscopy (FE-SEM, SIGMA 500, Germany). The optical properties were obtained using a UV-Vis-NIR spectrophotometer (UV-Vis-NIR, UV-3600 Plus, Shimadzu, Japan). Polarization-electric field (*P-E*) hysteresis loops were measured by a ferroelectric test system (Precision LC II; Radiant Technologies Inc, USA) at 100 Hz. The photoelectric response properties of the photodetector at different wavelengths and poling voltages were carried out by using a Keithley 2634 source meter, and during the measuring process, a 150 W low-pressure ultraviolet enhanced xenon lamp was used as the light source, and monochromatic light was obtained through a grating spectrometer. The optical power was measured using a standard Si photodetector.

3. RESULTS AND DISCUSSION

Fig. 1 shows XRD patterns recorded in the 2θ range of 20° to 60° for PLZT/TiO₂ heterojunction thin films with different TiO₂ layer numbers. Obviously, all thin films have good crystal quality, no second or impurity phases are detected, and only the diffraction peaks from the PLZT and TiO₂ thin films are noted. Furthermore, It can be seen that there is a (101) reflection peak around 25.3° corresponding to anatase TiO₂ structure (PDF 71-1166) [23-26], and with increasing TiO₂ thickness, the intensity of the (101) diffraction peak gradually improves. All above these results confirm the successful preparation of PLZT/TiO₂ heterojunction thin films.

Fig. 2(a)-(c) display the FE-SEM surface morphology images of the PLZT/TiO₂ heterojunction thin films with different number of TiO₂ layers. It can be clearly observed that the thin films with 1-layer and 3-layer TiO₂ have many pores and the surface are uneven. Different from above these films, the films with 2-layer TiO₂ exhibit good uniformity, smooth and compactness, and no obvious macroscopic cracks are observed on the surface of the film, which is beneficial for obtaining superior photoelectric response. To investigate the interfaces of the PLZT/TiO₂ heterojunction thin films with different TiO₂ layer numbers, cross-sectional images are characterized and given in Fig. 2(d)-(f). As shown, for these three kinds of thin films,

the thickness of the TiO_2 layer are respectively 65 nm, 101 nm and 145 nm, and the PLZT layer thickness is about 300 nm. Besides, a well-interacted interface is formed between PLZT and TiO_2 layers, which is helpful to transfer carriers at the interface.

Fig. 3(a) shows the optical transmission spectra of the TiO₂ thin films, and PLZT/TiO₂ heterojunction thin films with different number of TiO₂ layers in the wavelength range of 295-800 nm. Clearly, all films are nearly transparent in the visible region (400-800 nm) and exhibit strong absorption for the UV light, which are favorable for fabricating UV PDs. In addition, as the TiO₂ thin film layer number increases, the transmittance of the TiO₂ thin films and PLZT/TiO₂ heterojunction thin film in the ultraviolet region gradually decreases, which means that the absorption ability of the thin films for the ultraviolet light is enhanced and thus more superior photoelectric response may be achieved. Here, the optical band gap (E_g) of the TiO₂ thin films and PLZT/TiO₂ heterojunction thin films can be determined by Tauc's Law [27,28],

$$(\alpha h v)^2 = A(h v - E_{\sigma}) \tag{1}$$

where α , hv, and A are the absorption coefficient, photon energy, and constant, respectively. E_g can be calculated by extrapolating the linear portion of the $(\alpha hv)^2$ versus hv curve to zero, as displayed in the $(\alpha hv)^2$ vs. (hv) plots (Fig. 3(b)). Fig. 3(c) shows how the E_g of TiO₂ thin films and PLZT/TiO₂ heterojunction thin films varies with the TiO₂ layer number. Obviously, the E_g slightly decreases due to the influence of quantum size effect as the number of TiO₂ thin film layers increases [27]. Figure 3(d) illustrates room-temperature *P*-*E* hysteresis loops of the PLZT/TiO₂ heterojunction thin films. As seen, when the TiO₂ layer number is 1 and 2, square *P*-*E* loops, which is the typical character of ferroelectrics, are observed in the thin films, and the remnant polarization are respectively as high as 45.8 μ C/cm² and 31.4 μ C/cm², which is good for fabricating large depolarization electric field and thus the efficient separation of photo-generated electron-hole pairs. However, with increasing the TiO₂ layer number to 3, the square *P*-*E* loop is not obvious and the remnant polarization decreases to 14.8 μ C/cm², which is mainly attributed to large leakage current caused by the TiO₂ semiconductor materials without spontaneous polarization.

In order to better discuss the origin of the photoelectric response of the PLZT/TiO₂ heterojunction PDs, we prepare TiO₂ UV PD with the Au/TiO₂/FTO structure, as shown in Fig. 4(a), and study the variation of the photocurrent of the devices with the TiO₂ layer number under intermittent UV illumination with 315 nm light at which the photocurrent is the largest, as given in Fig. 4(b)-(d). Clearly, the device with 2-layer TiO₂ exhibit the largest photocurrent (~ 90 nA), which is further higher than those with 1-and 3-layer TiO₂ because of its superior ultraviolet light absorption and photo-generated carriers separation ability. For the 1-layer TiO₂ thin film, as observed in Fig. 3(a), it has the weakest ultraviolet light absorption ability attributed to its thin thickness, and thus the concentration of the photo-generated carriers is low, which leads to small photocurrent. However, when the TiO₂ thickness is increased to 3 layer, although the concentration of the photo-generated carriers is high, the carrier recombination probability is improved, and thus the photocurrent decreases.

Fig. 5(a) presents the schematic diagram of the PLZT/TiO₂ heterojunction PDs with the Au/PLZT/TiO₂/FTO structure. The devices are poled by a Keithley 2634 source meter, and the negative poling direction is defined as the positive electric field applied to the FTO electrode. In order to more accurately investigate the effects of the poling voltages on the photoelectric response performances, after the devices are poled, a 5 min break is introduced before taking the measurement. Figure 5(b)-(d) illustrate current-time (I-t) curves of PLZT/TiO₂ heterojunction PDs with different number of TiO₂ layers under zero bias. Those devices are measured under the ultraviolet light with 340 nm wavelength at which their photocurrent is the largest. As seen, all devices exhibit superior self-powered performances and their photocurrent obviously improve with the increase of the poling electric field. In comparison with other devices, the PD with 2-layer TiO₂ has the maximum photocurrent of 180 nA under the poling electric voltage of -0.5 V, which is attributed to its good microstructure, superior ultraviolet light absorption ability and high photo-generated carriers separation efficiency. More interestingly, the photocurrent values of the heterojunction devices are further higher than those of TiO₂ PDs (Fig. 4) and PLZT PDs with the same layers which indicates the reasonability of this kind of device structure for fabricating self-powered PDs with high photoelectric response [22].

To further understand the internal mechanism of the poling voltage enhanced photoelectric responses in these Au/PLZT/TiO₂/FTO heterojunction devices, we provide schematic energy band diagrams of the devices at no poling and negative poling voltages under light illumination, as given in Fig. 6. It can be seen from Fig. 6(a) that without poling electric field, the dominant driving force for photo-generated carriers in Au/PLZT/TiO₂/FTO devices is the heterojunction field (E_j) formed at the interface between PLZT and TiO₂ thin films, which is attributed to the difference of the work function for these two kind of thin films [29,30]. When a negative electric field is applied, a depolarization electric field (E_{dp}) is generated inside the PLZT thin film and its direction is the same as that of the E_j . As a result, at this point, the total built-in electric field (E_{tot}) is equal to $E_j + E_{dp}$, which is more beneficial for the separation of the photo-generated electron and hole pairs and the improvement of the carriers mobility, and thus the photocurrent obviously increases with the application of the negative poling voltage.

For photodetectors, the spectral responsivity (R) and specific detectivity (D^*) are two critical parameters to evaluate the performance of the devices. The responsivity R, which represents the efficiency of a detector responding to optical signals, is calculated as follow [31-33]:

$$R_{\lambda} = \frac{I_{light} - I_{dark}}{P_{\lambda}S} \tag{2}$$

where I_{light} is the photocurrent, I_{dark} is the dark current, P_{λ} is the light intensity at each wavelength, and S is the effective illumination area ($S = 0.05 \times 0.05 \text{ cm}^2$ in this work). D^* reflects the ability of a detector to detect signals from the noise environment and can be defined as Eq. (2) [34],

$$D^* = \frac{R_\lambda \sqrt{S}}{\sqrt{2eI_{dark}}} \tag{3}$$

where e is the charge of an electron ($e = 1.602 \times 10^{-19}$ C). Fig. 7(a)-(c) give the R and

 D^* of the PDs with different number of TiO₂ layers at the negative poling voltages. In those devices, because the device with 2-layer TiO₂ has the largest photocurrent of 180 nA, as observed in Fig. 5, it exhibit the highest responsivity of 21 mA/W and largest detectivity of 2.4×10^{11} Jones. In order to characterize the detection ability of the PDs on the light signal with different intensities, Fig. 8 gives the dependence of the photoelectric response on the light power. As seen, with decreasing light intensity, although the photocurrent of all PDs reduce due to the decrease of the concentration of photo-generated electron and holes, the *R* and *D** gradually increase [18,34-36]. Especially, under the weak-ultraviolet light illumination (0.18 mW/cm², 340 nm), the Au/PLZT/TiO₂/FTO heterojunction PDs with 2-layer TiO₂ exhibit an ultrahigh *R* of 91.4 mA/W and *D** of 9.7×10^{11} Jones, respectively, which indicates that the PDs can be applicable in a wide field.

Besides *R* and *D**, another important parameter for characterizing the photoelectric response performance of photodetectors is the response speed. The rise time is the time required to convert a UV signal into an electrical signal, while the decay time represents the time required for the current to be restored to its original state, and thus they are both expected to be as fast as possible [22]. The rise time (t_r) and the decay time (t_d) are respectively the time of the photocurrent increasing from 10% to 90% as the light turns on, and the photocurrent decreasing from 90% to 10% as the light turn off [10,34]. Fig. 9(a)-(c) display respective the *I*-*t* curve of the Au/PLZT/TiO₂/FTO heterojunction PDs with different number of TiO₂ layers under a single UV illumination period. As observed, the device with 2-layer TiO₂ exhibits

the faster response speed with a rise/decay time of 2.4/3.2 ms, which is resulted from its good surface morphology, large depolarization electric field and thus the most superior photo-generated carries separation ability.

The key parameters of the PLZT/TiO₂ heterojunction self-powered UV PDs with different TiO₂ layer numbers, are summarized in Table 1. Obviously, the PD with 2-layer TiO_2 exhibits the most superior photoelectric response characteristics with a fast response and recover time of 2.4/3.2 ms, large R of 21 mA/W and high D^* of 2.4 \times 10¹¹ Jones. In addition, we compare the photoelectric response performances of recently reported ferroelectric and semiconductor-based self-powered UV PDs, as presented in Table [10,11,13,18,32,37-41]. Clearly, the self-powered UV PD with the Au/PLZT/TiO₂/FTO heterostructure obtained in this work exhibits the largest responsivity, highest specific detectivity, and fastest response speed among those self-powered UV PDs. This occurs primarily because the collaborative enhancement effect of the built-in electric field resulting from the heterojunction interfaces of the PLZT and the TiO₂ thin films and the depolarization field caused by the high remnant polarization of PLZT thin films can lead to photo-generated electron-hole pairs separation more effective and faster. This work provides an effective and simple way that can be scaled up for designing and fabricating self-powered UV PDs with excellent photoelectric response performances.

4. CONCLUSIONS

In summary, we have successfully prepared PLZT/TiO₂ heterojunction thin films

with different TiO₂ layer numbers through a facile spin-coating method and based on these films, constructed self-powered UV PDs with Au/PLZT/TiO₂/FTO structure. In comparison with PLZT and TiO2-based PDs, because of the increase of the photo-generated carriers and the total built-in electric field, the photoelectric response performances of the PLZT/TiO₂ heterojunction PDs obviously improve. In addition, we have also found that the TiO₂ layer number has obvious effective on the photoelectric response properties of the PDs. When the TiO₂ layer number is 2, the exhibits superior microstructure, UV light absorption properties and PD photo-generated carriers separation efficiency, and thus the best self-powered UV light detection ability with a high responsivity of 21 mA/W, large specific detectivity of 2.4×10^{11} Jones, and fast response speed (rise time of 2.4 ms and decay time of 3.2 ms). Furthermore, under weak-light illumination (0.18 mW/cm²), the R and D^* can be respectively as high as 91.4 mA/W and 9.7×10^{11} Jones, which greatly broaden the applicable range of this kind of self-powered UV PDs. The ferroelectrics/ semiconductors heterojunction strategy proposed in this work provide references for the designing and fabrication next-generation high-performances self-powered UV photodetectors.

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Figure captions:

Fig. 1. XRD patterns of PLZT/TiO₂ heterojunction thin films with different TiO₂ layer numbers.

Fig. 2. (a)-(c) FE-SEM images and (d)-(f) FE-SEM cross-sectional images of the $PLZT/TiO_2$ heterojunction thin films with different TiO₂ layer numbers.

Fig. 3. (a) Optical transmittance spectra; (b) $(\alpha hv)^2$ versus hv plots; (c) band gap values of TiO₂ thin films and PLZT/TiO₂ heterojunction thin films; and (d) *P*-*E* loops of PLZT/TiO₂ heterojunction thin films with different TiO₂ layer numbers.

Fig. 4. (a) Schematic structure diagram of TiO₂ thin film UV photodetectors. (b)-(d) Time-dependent photocurrent response of TiO₂ photodetectors with various TiO₂ layer numbers: (b) TiO₂-1 layers; (c) TiO₂-2 layers; (d) TiO₂-3 layers.

Fig. 5. (a) Schematic structure diagram of PLZT/TiO₂ heterojunction UV photodetectors. (b)-(d) Time-dependent photocurrent response of PLZT/TiO₂ heterojunction photodetectors with various TiO₂ layer numbers: (b) PLZT/TiO₂-1 layers; (c) PLZT/TiO₂-2 layers; (d) PLZT/TiO₂-3 layers at the negative poling voltages under switched-on/off UV light illumination.

Fig. 6. The schematic energy band diagrams of the photodetector with Au/PLZT/TiO₂/FTO structure under different polarization states: (a) without poling;

(b) negative poling.

Fig. 7. The responsivity and specific detectivity of PLZT/TiO₂ heterojunction photodetectors with various TiO₂ layer numbers: (a) PLZT/TiO₂-1 layers; (b) PLZT/TiO₂-2 layers; (c) PLZT/TiO₂-3 layers at the negative poling voltages.

Fig. 8. *I-t* curves of PLZT/TiO₂ heterojunction photodetectors with different TiO₂ layer numbers at the negative poling voltages and 340 nm light irradiation with various intensities: (a) PLZT/TiO₂-1 layers; (b) PLZT/TiO₂-2 layers; (c) PLZT/TiO₂-3 layers. Light power density-dependent *R*, and *D** with the TiO₂ layer number: (d) PLZT/TiO₂-1 layers; (e) PLZT/TiO₂-2 layers; (f) PLZT/TiO₂-3 layers.

Fig. 9. Single-cycle, time-dependent photocurrent of PLZT/TiO₂ heterojunction photodetectors with various TiO₂ layer numbers: (a) PLZT/TiO₂-1 layers; (b) PLZT/TiO₂-2 layers; (c) PLZT/TiO₂-3 layers at the negative poling voltages and under switched-on/off UV light illumination.

Fig. 1.

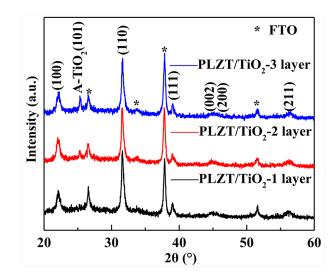


Fig. 2.

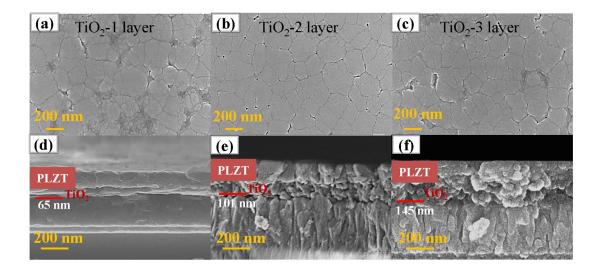


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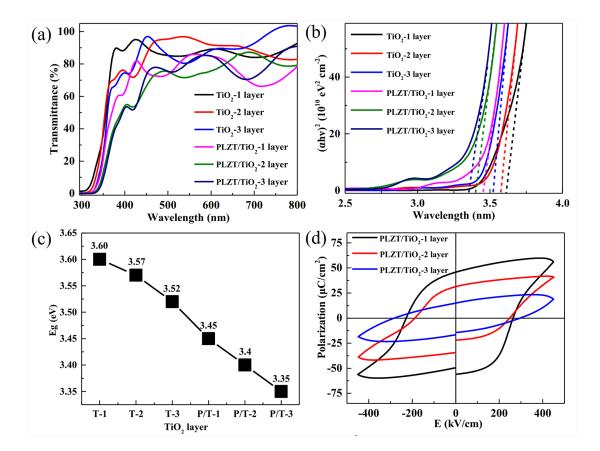


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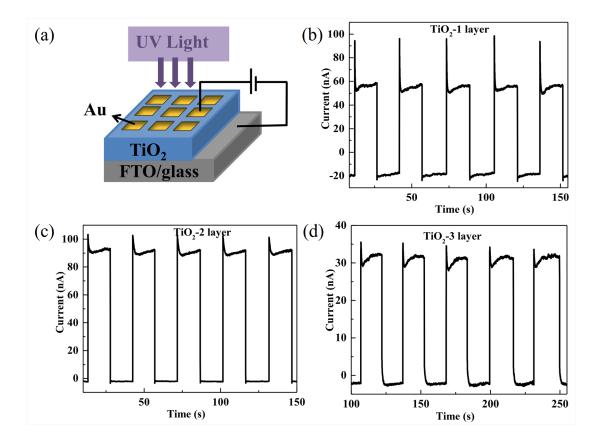


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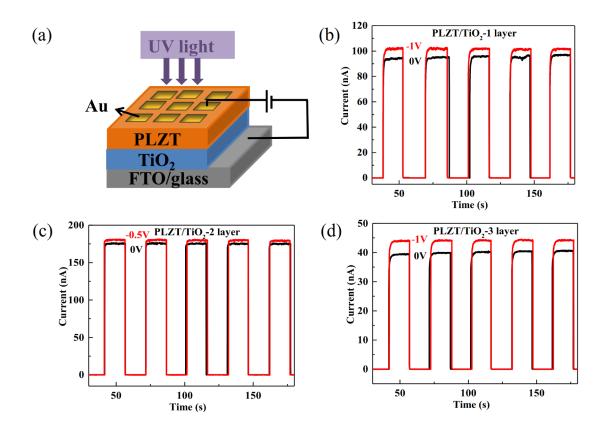


Fig. 6.

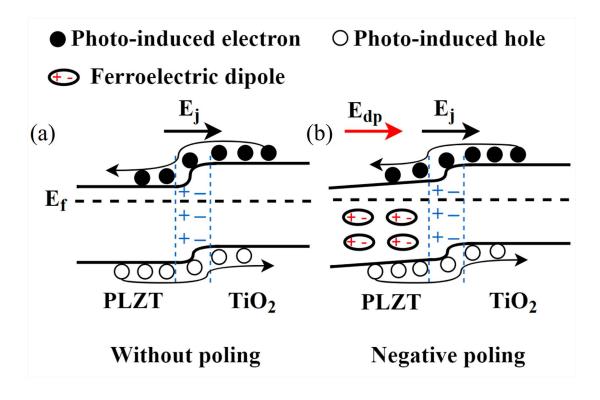


Fig. 7.

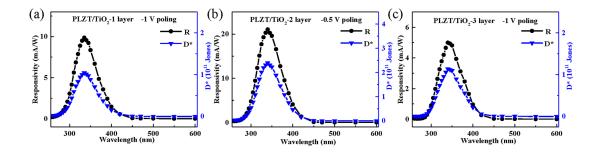


Fig. 8.

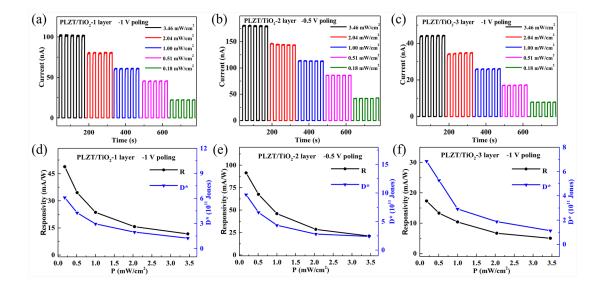


Fig. 9.

