

25 **Abstract**

26 Multiple Signal Classification (MUSIC) – a directional scanning and searching algorithm, has
27 gained its prominence in phased array-facilitated nondestructive evaluation. Nevertheless,
28 prevailing MUSIC algorithms are largely bound up with the use of a dense linear array, which fail
29 to access the full planar area of an inspected sample, leaving blind zones to which an array fails
30 to scan, along with the incapability of differentiating multiple damage sites that are close one from
31 another. To break above limitations, conventional MUSIC algorithm is ameliorated in this study,
32 by manipulating the signal representation matrix at each pixel using the excitation signal series,
33 instead of the scattered signal series, which enables the use of a sparse sensor network with
34 arbitrarily positioned transducers. In the ameliorated MUSIC (Am-MUSIC), the orthogonal
35 attributes between the signal subspace and noise subspace inherent in the signal representation
36 matrix is quantified, in terms of which the Am-MUSIC yields a full spatial spectrum of the
37 inspected sample, and damage, if any, can be visualized in the spectrum. Am-MUSIC is validated,
38 in both simulation and experiment, by evaluating single and multiple sites of damage in plate-like
39 waveguides with a sparse sensor network. Results verify that i) detectability of Am-MUSIC-
40 driven damage imaging is not limited by damage quantity; ii) Am-MUSIC has full access to a
41 sample, eliminating blind zones; and iii) the amelioration expands conventional MUSIC from
42 phased array-facilitated nondestructive evaluation to health monitoring using built-in sparse
43 sensor networks.

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45 **Keywords:** structural health monitoring; guided ultrasonic waves; multiple signal classification
46 (MUSIC); phased array; sparse sensor network

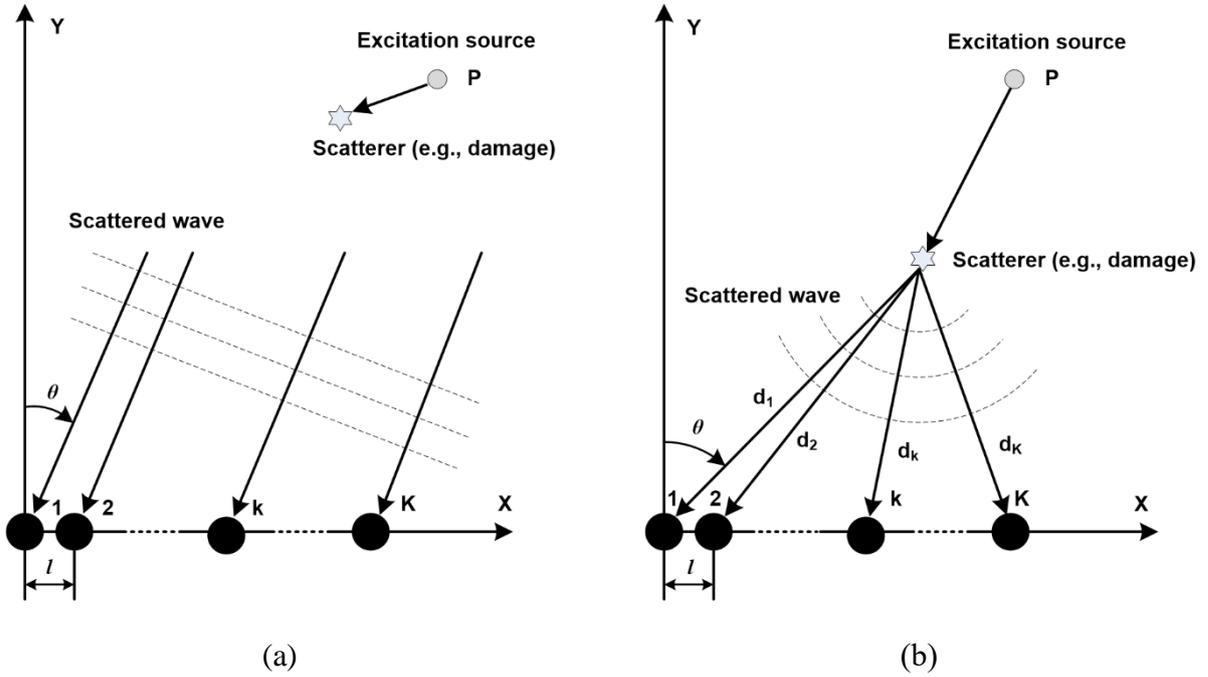
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48 **1. Introduction**

49 With the incentive to “visualise” hidden material defect or structural damage, continued effort has
50 been made to projecting identified results, by means of proper imaging algorithms, in synthetic
51 illustration, in which anomaly, if any, can be imaged intuitively [1]. Amongst various anomaly
52 imaging approaches, those by virtue of guided ultrasonic waves (GUWs) have demonstrated
53 eminent detectability, accuracy and precision [2-5], as represented by tomography-based imaging
54 [6], delay-and-sum imaging [7], time-reversal focusing imaging [8], probability-based imaging
55 [9], and array signal processing-based imaging [10-14], to name a few. Amongst them, the array
56 signal processing-based imaging can be implemented in various modalities, including sparse
57 reconstruction [10], minimum variance distortionless response method [11], subspace fitting [12],
58 maximum-likelihood method [13], and Multiple Signal Classification (MUSIC) [14].

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60 In particular, the MUSIC algorithm, with its theoretical framework shaped by Schmit [15] in 1981
61 for frequency estimation and radio direction finding, is a directional scanning and searching
62 method to unbiasedly estimate signal features in terms of the orthogonal attributes between signal
63 subspace and noise subspace. With a directional scanning ability, MUSIC has been extended to
64 various application domains such as radar positioning [16], sonar [17], seismic exploration [18],
65 biomedicine [19] and so forth. MUSIC has also proven effectiveness in GUW-based damage
66 imaging. Representatively, Stepinski and Engholm [20] revamped a conventional MUSIC
67 algorithm for estimating the *direction of arrival* (DOA) of incoming waves in plates with a
68 uniform circular array. Yang *et al* [21, 22] employed the MUSIC algorithm to calculate the arrival
69 times of impact-induced waves in conjunction with the use of wavelet transform, whereby to
70 predict the location of an impact to a plate. Majority of MUSIC-driven damage evaluation lies in
71 the *far-field hypothesis* that assumes a wave scatterer (e.g., damage) within the inspection region
72 is sufficiently far from the phased array, so that the waves emanating from the scatterer can be
73 considered as a plane wave when they arrive at the array, as illustrated schematically in **Fig. 1(a)**.

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Fig. 1. Use of phased array for evaluation of damage in a planar inspection region [23]: (a) far-field scenario; and (b) near-field scenario (d_k : the distance from the wave source through the damage then to the k_{th} array element).

Despite demonstrated effectiveness in numerous applications, the *far-field hypothesis* is often questioned in practical implementation, as the damage location is unknown *a priori*, and the wavefront scattered by the damage is naturally cylindrical rather than planar, provided the damage is in the near-field – the scenario in **Fig. 1(b)**. In recognition of such deficiency that prevailing MUSIC algorithms inherently has, enhancement has been made to improve conventional MUSIC-driven damage evaluation. Zhong *et al* [23, 24] developed a revamped MUSIC algorithm based on the Taylor expansion theory, applicable to detecting near-field damage in a composite oil tank. Extending this study, Yuan *et al* [25] took into account the anisotropy of composite structures and proposed a single-frequency component-based MUSIC algorithm able to improve the precision of locating a near-field impact site in composites. Zuo *et al* [26] calculated the cross-correlation function for the scattered signals received by a damage scattering model and the residual signals received in experiment, and applied the two-dimensional (2-D) MUSIC algorithm to identify damage in plate-like structures. Bao *et al* [27] combined the transmitter beamforming and weighted image fusion to enhance the conventional MUSIC, endowing it with the capability of localizing near-field corrosion in aluminum plates. Bao *et al* [28] further proposed a compensated MUSIC algorithm by considering the effect of both the localization error caused by structural anisotropy and the sensor position error, showing improved detection accuracy.

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Notwithstanding the forgoing, the prevailing MUSIC-based damage imaging approaches in general present the following limitations:

- i) *use of linear phased arrays*: the MUSIC algorithm, as an eigen-structure mathematic approach, is not inherently restricted to the use of linear arrays. When used for damage identification, a linear array that features a dense configuration of transmitter elements is usually adopted, and the element pitch (l) shall be uniform and small enough, ideally satisfying $l \leq \lambda / 2$ (λ : the wavelength of the wave generated by the array [29]), to facilitate construction of the signal representation matrix with the scattered signal series. In general, the conventional MUSIC algorithms do not sustain the use of a sparse sensor network with individual sensors at arbitrary locations;
- ii) *existence of blind zone*: the beamforming capacity of the algorithm degrades when the scanning angle is close to 0° or 180° . In most circumstances, those regions, where the scanning angles are in the range of $[0, 30]$ or $[150, 180]$, are deemed blind zones [30], in which damage, if any, may be overridden;
- iii) *ambiguous results due to mirror effect*: the identified damage using a MUSIC algorithm might be a mirrored dummy of the true damage which is located symmetrically with regard to the array surface [31]; and
- iv) *obscure multiple damage sites*: when MUSIC is used for imaging multiple damage sites, the number of scatterers shall be predicted beforehand. To this end, a threshold is selected, and the eigenvalues of the covariance matrix of the received signals that are larger than this threshold shall be counted as the number of the scatterers. However, selection of the threshold is a highly subjective manner at the discretion of individuals, and it is prone to contamination of measurement noise [32]. This results in inferior imaging resolution and makes it challenging to differentiate multiple damage sites that are close one from another.

Aimed at surmounting the above limitations that prevailing MUSIC-based damage imaging approaches may encounter, an ameliorated MUSIC (Am-MUSIC) algorithm is developed by manipulating the signal representation matrix at each image pixel using the excitation signal series instead of the scattered signal series. Thanks to that, Am-MUSIC algorithm does not necessarily entail the use of a linear phased array, and instead it is compatible with a sparse sensor network in which individual transducers can be positioned arbitrarily. At each image pixel, the orthogonal attributes between the signal subspace and noise subspace inherent in signal representation matrix is quantified, in terms of which Am-MUSIC yields a full spatial spectrum of the inspected sample, to visualize damage, regardless of its quantity in the sample. Both simulation and experiment are

141 performed to validate the Am-MUSIC algorithm, by evaluating single and multiple sites of
 142 damage in plate-like waveguides with a sparse sensor network comprising only a handful of
 143 miniaturized lead zirconate titanate (PZT) wafers.

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 145 The rest of this paper is organized as follows. The conventional MUSIC-driven damage imaging,
 146 based on the near-field hypothesis, is briefed in Section 2, on which basis the Am-MUSIC
 147 algorithm is developed, with key amelioration detailed in Section 3. Numerical verification of
 148 Am-MUSIC is illustrated in Section 4, followed with experimental validation in Section 5.

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150 2. Near-Field MUSIC Algorithm

151 GUWs guided by a plate-like waveguide, a.k.a. *Lamb waves*, are of a multimodal and dispersive
 152 nature. At a given frequency, Lamb waves feature a multitude of wave modes which can be
 153 classified as the symmetric and antisymmetric modes. We consider a pure, monochromatic Lamb
 154 wave mode in the waveform of a toneburst, as the excitation signal $s(t)$. $s(t)$ is defined in a
 155 complex domain as

$$156 \quad s(t) = u(t) \exp^{i\omega_0 t}, \quad (1)$$

157 where $u(t)$ denotes a window function to regulate the toneburst, t the time, i the imaginary unit,
 158 and ω_0 the central frequency of the toneburst. With the attenuation in magnitude as wave
 159 propagation in consideration, the Lamb wave, $R(\omega)$, after travelling the distance d can be
 160 represented, in the frequency domain, as

$$161 \quad R(\omega) = \frac{d_0}{\sqrt{d}} S(\omega) \exp^{-ikd}. \quad (2)$$

162 In the above, d_0 signifies an initial distance with regard to which the wave attenuation is calibrated;

163 $S(\omega)$ is the corresponding Fourier representation of $s(t)$; $k = \frac{\omega_0}{c}$, where k denotes the
 164 wavenumber and c represents the propagation velocity of the considered monochromatic Lamb
 165 wave mode.

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 167 Substituting Eqs. (1) into (2), the Lamb wave $r(t)$ when it arrives at the distance d can be
 168 yielded, in the time domain, as

$$169 \quad r(t) = \frac{d_0}{\sqrt{d}} F^{-1} \left\{ s(\omega) \exp^{-i\frac{\omega_0 d}{c}} \right\} = \frac{d_0}{\sqrt{d}} s\left(t - \frac{d}{c}\right) = \frac{d_0}{\sqrt{d}} u\left(t - \frac{d}{c}\right) \exp^{i\omega_0\left(t - \frac{d}{c}\right)}, \quad (3)$$

170 where $r(t)$ is the inverse Fourier transform of $R(\omega)$ and F^{-1} is the inverse Fourier transform.

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 172 For an intact waveguide, the captured wave signal, denoted with $r^{\text{measured-intact}}(t)$, is the direct
 173 arrival wave $r^{\text{direct}}(t)$, boundary-reflection wave $r^{\text{boundary-reflection}}(t)$ with incoherent noise
 174 $w^{\text{measured-intact}}(t)$, as

$$175 \quad r^{\text{measured-intact}}(t) = r^{\text{direct}}(t) + r^{\text{boundary-reflection}}(t) + w^{\text{measured-intact}}(t), \quad (4)$$

176 where $r^{\text{direct}}(t)$ is the arrival wave propagating along the path from the wave source to the wave
 177 receiver. Provided damage is present at an unknown location in the waveguide, the damage can
 178 be modeled as a secondary wave source to scatter the incoming Lamb waves. Ignoring mode
 179 conversion that is fairly weak in magnitude, the measured signal $r^{\text{measured-damage}}(t)$ comprises the
 180 direct arrival wave $r^{\text{direct}}(t)$, boundary-reflection wave $r^{\text{boundary-reflection}}(t)$, additional scattered wave
 181 from the damage $r^{\text{scattered}}(t)$, and the incoherent noise $w^{\text{measured-damage}}(t)$, as

$$182 \quad r^{\text{measured-damage}}(t) = r^{\text{direct}}(t) + r^{\text{boundary-reflection}}(t) + r^{\text{scattered}}(t) + w^{\text{measured-damage}}(t), \quad (5)$$

183 where $r^{\text{scattered}}(t)$ is the arrival wave propagating along a scattered path (namely, the path from the
 184 wave source to the damage and then to the wave receiver). Suppose that the direct waves and
 185 boundary-reflection waves are the same at $r^{\text{measured-intact}}(t)$ and $r^{\text{measured-damage}}(t)$, $r^{\text{scattered}}(t)$ which
 186 carries information pertaining to the damage location can be obtained through benchmarking
 187 reference signals obtained from the intact status, as

$$188 \quad r^{\text{measured-damage}}(t) - r^{\text{measured-intact}}(t) = r^{\text{scattered}}(t) + w(t) = r^{\text{residual}}(t), \quad (6)$$

189 where $w(t)$ is the difference between the two noise terms $w^{\text{measured-intact}}(t)$ and $w^{\text{measured-damage}}(t)$ in
 190 the intact and current statuses. Here, for convenience of discussion, the terms of $r^{\text{scattered}}(t) + w(t)$
 191 is referred to as the *residual signal*.

192
 193 With the near-field assumption, as schematically illustrated in **Fig. 1(b)**, Lamb wave is excited at
 194 a foreknown position P , scattered by the damage, and then received by a linear sensor array that
 195 consisting of K transducing elements with a uniform element spacing l . According to Eq. (3), the
 196 scattered signal received by the first array element, $r_1^{\text{scattered}}(t)$, is

$$197 \quad r_1^{\text{scattered}}(t) = \frac{d_0}{\sqrt{d_1}} s\left(t - \frac{d_1}{c}\right) = \frac{d_0}{\sqrt{d_1}} u\left(t - \frac{d_1}{c}\right) \exp^{i\alpha_0\left(t - \frac{d_1}{c}\right)}, \quad (7)$$

198 where d_1 signifies the distance from the wave source through the damage then to the first array
 199 element. Let $\tau_k = \frac{d_1 - d_k}{c}$ (*i.e.*, the time delay between two arrival signals captured by the first and

200 the k^{th} ($k=1, 2, \dots, K$) element in the array), and then the scattered wave signal received by the
 201 k^{th} element, $r_k^{\text{scattered}}(t)$, can be expressed as

$$202 \quad r_k^{\text{scattered}}(t) = \frac{d_0}{\sqrt{d_k}} s\left(t - \frac{d_k}{c}\right) = \frac{d_0}{\sqrt{d_k}} u\left(t - \frac{d_k}{c}\right) \exp^{i\omega_0\left(t - \frac{d_k}{c}\right)} = \frac{d_0}{\sqrt{d_k}} u\left(t - \frac{d_1}{c} + \tau_k\right) \exp^{i\omega_0\left(t - \frac{d_1}{c} + \tau_k\right)} .$$

(8)

203
 204 With the assumption that the array element spacing l is sufficiently small (namely, $l \leq \lambda / 2$,
 205 where λ is the wavelength of wave signal), $r_k^{\text{scattered}}(t)$ can be obtained based on the first element
 206 scattered signal $r_1^{\text{scattered}}(t)$ (defined in Eq. (7)) as

$$207 \quad \begin{aligned} r_k^{\text{scattered}}(t) &= \frac{d_0}{\sqrt{d_k}} u\left(t - \frac{d_1}{c} + \tau_k\right) \exp^{i\omega_0\left(t - \frac{d_1}{c} + \tau_k\right)} \\ &\approx \frac{d_0}{\sqrt{d_k}} u\left(t - \frac{d_1}{c}\right) \exp^{i\omega_0\left(t - \frac{d_1}{c} + \tau_k\right)} \\ &= \sqrt{\frac{d_1}{d_k}} r_1^{\text{scattered}}(t) \exp^{i\omega_0\tau_k} . \end{aligned} \quad (9)$$

208 According to the cosine theorem [33] and second-order Taylor expansion [34], τ_k can be re-
 209 written as

$$210 \quad \begin{aligned} \tau_k &= \frac{d_1 - d_k}{c} = \frac{d_1 - \sqrt{d_1^2 + (k-1)^2 l^2 - 2d_1(k-1)l \cos(90^\circ - \theta)}}{c} \\ &= \frac{-l \sin \theta}{c} (k-1) + \left(\frac{-l^2}{cd_1} \cos^2 \theta\right) (k-1)^2 + O\left(\frac{l^2}{d_1^2}\right) , \end{aligned} \quad (10)$$

211 where $O\left(\frac{l^2}{d_1^2}\right)$ denotes those terms, the order of which is greater than or equal to $\frac{l^2}{d_1^2}$. Using the
 212 second-order Taylor series approximation, the scattered wave signal received by the k^{th} element
 213 retreats to

$$214 \quad r_k^{\text{scattered}}(t) = \sqrt{\frac{d_1}{d_k}} r_1^{\text{scattered}}(t) \exp^{i\omega_0\tau_k} = \sqrt{\frac{d_1}{d_k}} r_1^{\text{scattered}}(t) \exp^{i\omega_0\left(\frac{-l \sin \theta}{c}(k-1) + \left(\frac{-l^2}{cd_1} \cos^2 \theta\right)(k-1)^2\right)} . \quad (11)$$

215 Letting $b_k(d, \theta) = \sqrt{\frac{d_1}{d_k}} \exp^{i\omega_0\left(\frac{-l \sin \theta}{c}(k-1) + \left(\frac{-l^2}{cd_1} \cos^2 \theta\right)(k-1)^2\right)}$, as the array steering factor for the k^{th}

216 scattered signal, and recalling the noise term in Eq. (6), the k^{th} residual signal, $r_k^{\text{residual}}(t)$, can be
 217 expressed as

$$218 \quad r_k^{\text{residual}}(t) = b_k(d, \theta) r_1^{\text{scattered}}(t) + w_k(t) . \quad (12)$$

219 For the linear array with K elements, the residual signal vector $\mathbf{R}^{\text{residual}}(t)$ can thus be obtained
 220 and expressed in a signal representation matrix, which reads

$$221 \quad \mathbf{R}^{\text{residual}}(t) = \mathbf{B}(d, \theta) \mathbf{r}_1^{\text{scattered}}(t) + \mathbf{W}(t), \quad (13)$$

222 where

$$223 \quad \mathbf{R}^{\text{residual}}(t) = [r_1^{\text{residual}}(t), L, r_k^{\text{residual}}(t), L, r_K^{\text{residual}}(t)]^T,$$

$$224 \quad \mathbf{B}(d, \theta) = [b_1(d, \theta), L, b_k(d, \theta), L, b_K(d, \theta)]^T = \begin{bmatrix} 1 \\ \mathbf{M} \\ \sqrt{\frac{d_1}{d_k}} \exp \left[i\omega_0 \left(\frac{-l \sin \theta}{c} (k-1) + \left(\frac{-l^2}{cd_1} \cos^2 \theta \right) (k-1)^2 \right) \right] \\ \mathbf{M} \\ \sqrt{\frac{d_1}{d_K}} \exp \left[i\omega_0 \left(\frac{-l \sin \theta}{c} (K-1) + \left(\frac{-l^2}{cd_1} \cos^2 \theta \right) (K-1)^2 \right) \right] \end{bmatrix},$$

$$225 \quad \mathbf{W}(t) = [w_1(t), L, w_k(t), L, w_K(t)]^T.$$

226
 227 Prevailing MUSIC-based damage imaging approaches have been developed by virtue of the signal
 228 representation matrix as defined in Eq. (13). They, in general, present the following limitations
 229 during practical implementation, as preliminarily commented in the preceding section:

- 230 i) In Eq. (9), the operation of approximation, $u(t - \frac{d_1}{c} + \tau_k) \approx u(t - \frac{d_1}{c})$, lies in the premise that
 231 τ_k is negligibly small. To accommodate such a pre-requisite, the element spacing in the
 232 phased array must be sufficiently small ($l \leq \lambda / 2$), leading to a dense configuration of the
 233 transducing elements; and
 234 ii) In Eq. (11), the steering vector is approximated using the second-order Taylor approximation,
 235 and the range error introduced by such approximation is remarkable when the damage is close
 236 to the array. For a range that is smaller than twice the array length (i.e., the length from the
 237 first element to the k^{th} element), such error could be 10% or above due to such approximation
 238 [35]. In addition, the steering vectors at the scanning angles θ and $180^\circ - \theta$ have the same
 239 value in Eq. (11), resulting in ambiguous results due to mirror effect.

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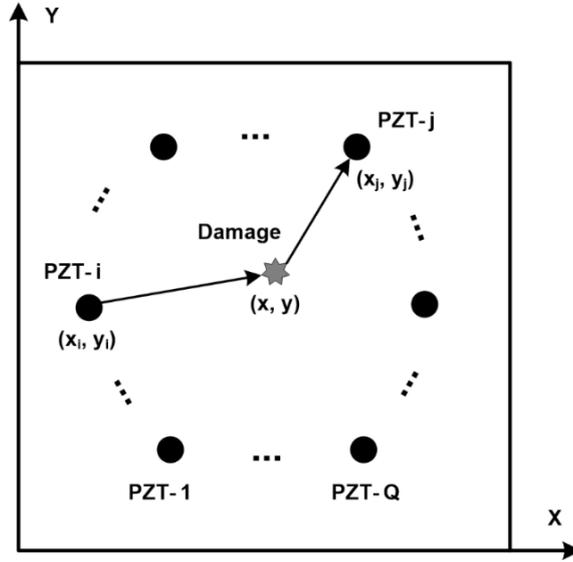
241 3. Am-MUSIC with A Sparse Sensor Network

242 Aimed at circumventing the above key limitations that conventional MUSIC-based damage

243 imaging possesses, the original MUSIC algorithm is revamped. Different from the use of a linear
 244 phased array, we allow a sparse sensor network with individual transducers that are randomly
 245 positioned. Without loss of generality, consider a sparse sensor network comprising Q PZT wafers
 246 (labelled as PZT-1, PZT-2, ..., PZT- j , ..., PZT- Q), as shown in **Fig. 2**. Positioned at an arbitrary
 247 location within the inspection region, each PZT wafer acts as either a wave transmitter or a wave
 248 receiver, leading to $M = Q(Q-1)/2$ transmitter–receiver pairs in the sensor network. Provided
 249 damage exists at pixel (x, y) within the inspection area, the propagation distance, d_{mxy} , for a
 250 Lamb wave, which is generated by the i^{th} transmitter at (x_i, y_i) , scattered by damage at (x, y)
 251 and then propagates to the j^{th} receiver at (x_j, y_j) , is

$$d_{mxy} = \sqrt{(x-x_i)^2 + (y-y_i)^2} + \sqrt{(x-x_j)^2 + (y-y_j)^2} = c \cdot t_{mxy}, \quad (14)$$

252 and t_{mxy} is the time for the wave traveling along the scattered path.



254
 255 **Fig. 2.** A plate waveguide with a sparse sensor network of Q PZT wafers.

256 Therefore, the scattered signal received by the m^{th} transmitter–receiver pair, $r_m^{\text{scattered}}(t)$, can be
 257 written according to Eq. (3) as

$$r_m^{\text{scattered}}(t) = \frac{d_0}{\sqrt{d_{mxy}}} s\left(t - \frac{d_{mxy}}{c}\right) \quad (m = 1, 2, L, M). \quad (15)$$

258 Equation (15) argues that for M transmitter–receiver pairs rendered by the sensor network,
 259 different scattering paths feature different degrees of time delay. A time shift, t_{mxy} , is then applied
 260 to the m^{th} scattered signal $r_m^{\text{scattered}}(t)$ in Eq. (15), as

262
$$r_m^{\text{scattered}}(t + t_{mxy}) = \frac{d_0}{\sqrt{d_{mxy}}} s\left(t - \frac{d_{mxy}}{c} + t_{mxy}\right) = \frac{d_0}{\sqrt{d_{mxy}}} s(t). \quad (16)$$

263 Letting $a_{mxy} = \frac{d_0}{\sqrt{d_{mxy}}}$ (a_{mxy} is referred to as the *array steering factor* for the m^{th} scattered signal
264 in what follows), Eq. (16) can be rewritten as

265
$$r_m^{\text{scattered}}(t + t_{mxy}) = a_{mxy} s(t). \quad (17)$$

266 With the noise term ($w(t)$ in Eq. (6)) in consideration, the residual signal vector for a total of M
267 received signals which are respectively scattered by the damage at pixel (x, y) , $\mathbf{R}_{xy}^{\text{residual}}(t)$, can be
268 expressed as the signal representation matrix

269
$$\mathbf{R}_{xy}^{\text{residual}}(t) = \mathbf{A}_{xy} s(t) + \mathbf{W}(t), \quad (18)$$

270 where

271
$$\mathbf{R}_{xy}^{\text{residual}}(t) = [r_1^{\text{residual}}(t + t_{1xy}), L, r_m^{\text{residual}}(t + t_{mxy}), L, r_M^{\text{residual}}(t + t_{Mxy})]^T,$$

272
$$\mathbf{A}_{xy} = [a_{1xy}, L, a_{mxy}, L, a_{Mxy}]^T,$$

273
$$\mathbf{W}(t) = [w_1(t + t_{1xy}), L, w_m(t + t_{mxy}), L, w_M(t + t_{Mxy})]^T.$$

274 Equation (18) implies that after compensating for the time delay to each residual signal, the
275 residual signal vector can be defined using the excitation signal series, instead of using the
276 scattered signal series as a conventional MUSIC algorithm does (Eq. (13)). It is such a merit of
277 the ameliorated MUSIC (Am-MUSIC) algorithm that enables the use of a sparse sensor network
278 with arbitrarily positioned transducers.

279
280 Recalling the MUSIC algorithm, the covariance matrix \mathbf{C} of the residual signal vector at pixel
281 (x, y) within the inspection region yields as

282
$$\begin{aligned} \mathbf{C} &= E[\mathbf{R}_{xy}^{\text{residual}}(t) \mathbf{R}_{xy}^{\text{residual}}(t)^H] \\ &= \mathbf{A}_{xy} E[s(t) g s(t)^H] \mathbf{A}_{xy}^H + \mathbf{A}_{xy} E[s(t) g \mathbf{W}(t)^H] + E[\mathbf{W}(t) g s(t)^H] \mathbf{A}_{xy}^H + E[\mathbf{W}(t) g \mathbf{W}(t)^H], \end{aligned} \quad (19)$$

284 where $E[\]$ denotes covariance computation, and superscript H represents the complex conjugate
285 transpose.

286
287 As the source signal and noise signal are uncorrelated and mutually independent, the covariance
288 matrix \mathbf{C} can be simplified as

289
$$\mathbf{C} = \mathbf{A}_{xy} \mathbf{R}_s \mathbf{A}_{xy}^H + \sigma^2 \mathbf{I}, \quad (20)$$

290 where $\mathbf{R}_s = E[s(t)gs(t)^H]$, and it signifies the covariance matrix of the source signal. σ^2 is noise
 291 power and \mathbf{I} the covariance matrix of the noise signal. The covariance matrix \mathbf{C} can be
 292 decomposed into two parts: namely a signal-related part and a noise-related part, as

$$293 \quad \mathbf{C} = \mathbf{U}\mathbf{\Sigma}\mathbf{U}^H = \mathbf{U}_s\mathbf{\Sigma}\mathbf{U}_s^H + \mathbf{U}_N\mathbf{\Sigma}\mathbf{U}_N^H, \quad (21)$$

294 where $\mathbf{U} = [\mu_1, \mu_2, \dots, \mu_M]$, and the columns of \mathbf{U} are the singular vectors; $\mathbf{\Sigma}$ is a diagonal
 295 matrix with singular values arranged in a descending order of magnitudes. Considering that \mathbf{A}_{xy}
 296 is the steering vector at pixel (x, y) with the dimension of $M \times 1$ and $\mathbf{A}_{xy}\mathbf{R}_s\mathbf{A}_{xy}^H$ in Eq. (20) is
 297 decomposed as $\mathbf{U}_s\mathbf{\Sigma}\mathbf{U}_s^H$ in Eq. (21), $\mathbf{U}_s = [\mu_1]$ denoting the signal subspace spanned by the
 298 eigenvectors corresponding to the first largest eigenvalue; and $\mathbf{U}_N = [\mu_2, \mu_3, \dots, \mu_M]$,
 299 representing the noise subspace spanned by the eigenvectors corresponding to the remaining $M -$
 300 1 eigenvalues.

301
 302 Based on Eqs. (20) and (21), the following expression can be obtained after multiplying
 303 covariance matrix \mathbf{C} with the noise subspace \mathbf{U}_N

$$304 \quad \mathbf{A}_{xy}\mathbf{R}_s\mathbf{A}_{xy}^H\mathbf{U}_N = \mathbf{0}. \quad (22)$$

305 As \mathbf{R}_s is a full rank matrix, Eq. (22) is further simplified as

$$306 \quad \mathbf{A}_{xy}^H\mathbf{U}_N = \mathbf{0}. \quad (23)$$

307 Equation (23) argues that the steering vector \mathbf{A}_{xy} at the position of damage is orthogonal with the
 308 noise subspace \mathbf{U}_N . This characteristic makes it possible for the Am-MUSIC to calculate the
 309 steering vector at each pixel across the entire inspection region and calibrate the degree of
 310 orthogonality between the steering vector and the noise subspace with the squared norm of vector
 311 $\mathbf{A}_{xy}^H\mathbf{U}_N$ as

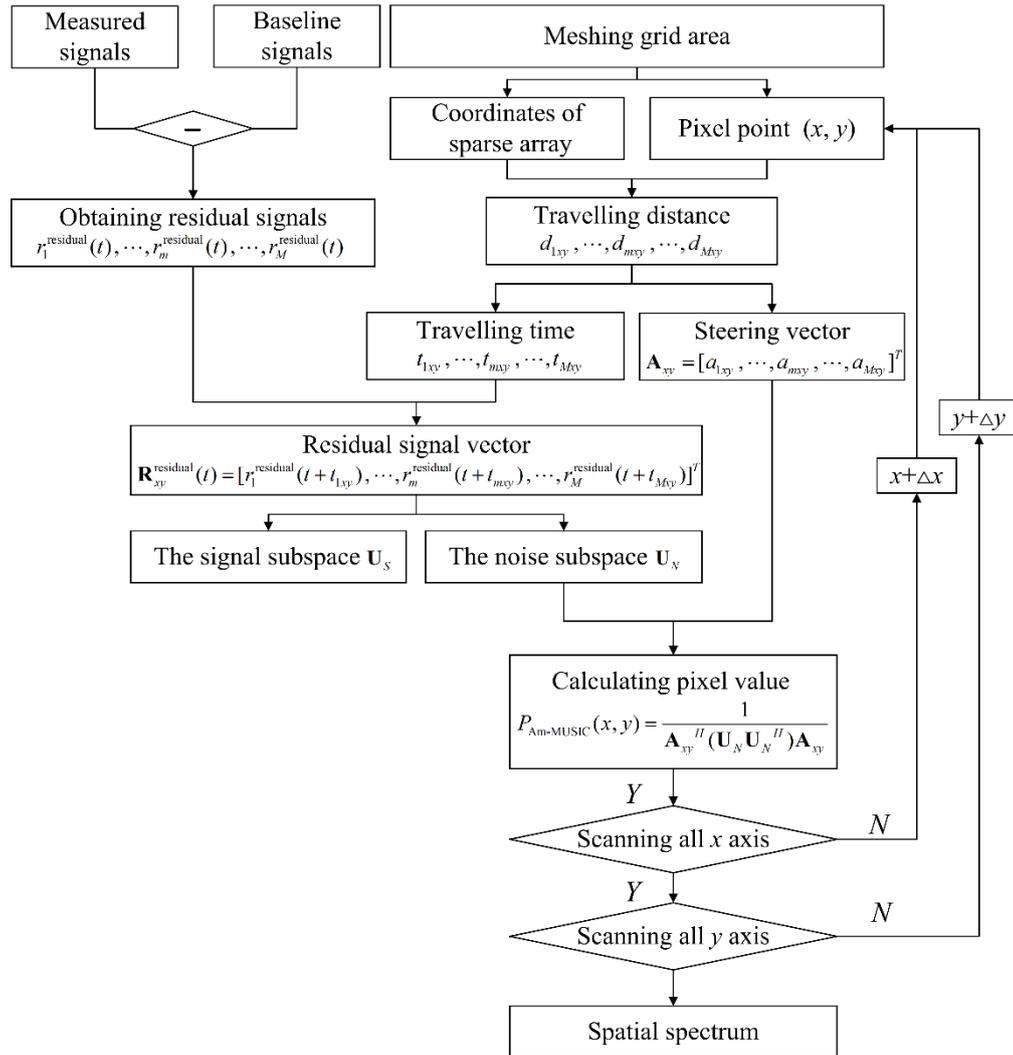
$$312 \quad \beta^2 = \|\mathbf{A}_{xy}^H\mathbf{U}_N\|^2 = \mathbf{A}_{xy}^H(\mathbf{U}_N\mathbf{U}_N^H)\mathbf{A}_{xy}. \quad (24)$$

313 Taking a reciprocal of the squared norm expression creates a peak in the spatial spectrum that
 314 corresponds to the damage location. Am-MUSIC algorithm defines the pixel value
 315 ($P_{\text{Am-MUSIC}}(x, y)$) within the inspection region as

$$316 \quad P_{\text{Am-MUSIC}}(x, y) = \frac{1}{\mathbf{A}_{xy}^H(\mathbf{U}_N\mathbf{U}_N^H)\mathbf{A}_{xy}}. \quad (25)$$

317 Equation (24) yields a full spatial spectrum for the inspection region, in which $P_{\text{Am-MUSIC}}(x, y)$
 318 culminates at the damage location.

319 In summary, the complete procedure of the proposed Am-MUSIC algorithm is flowcharted in a
 320 nutshell in **Fig. 3**. Notably, the Am-MUSIC algorithm calculates the signal representation matrix
 321 at each pixel throughout the entire inspection region, and it is therefore the number of scatterers
 322 (i.e., the number of multiple damage sites) is no longer required, in contrast to a conventional
 323 MUSIC algorithm in which the number of scatterers shall be predicted beforehand at a subjective
 324 discretion of individuals. It is also noteworthy that the computational cost of Am-MUSIC does
 325 not intend to increase compared with conventional MUSIC algorithms, regardless of the fact that
 326 the calculation is performed at every signal pixel – that is because the signal representation matrix
 327 can be formed cost-effectively and the steering factor in Eq. (18) can be calculated efficiently
 328 compared with that of conventional MUSIC algorithm in Eq. (13). In addition, it is such a merit
 329 of the Am-MUSIC algorithm that makes it possible to gauge the local region of interest (RoI) only
 330 – the vicinity in the sample where damage may exist, rather than scanning the entire sample. Such
 331 a merit remarkably lowers the computational cost and unburdens computing hardware when the
 332 inspection region has of large dimensions.

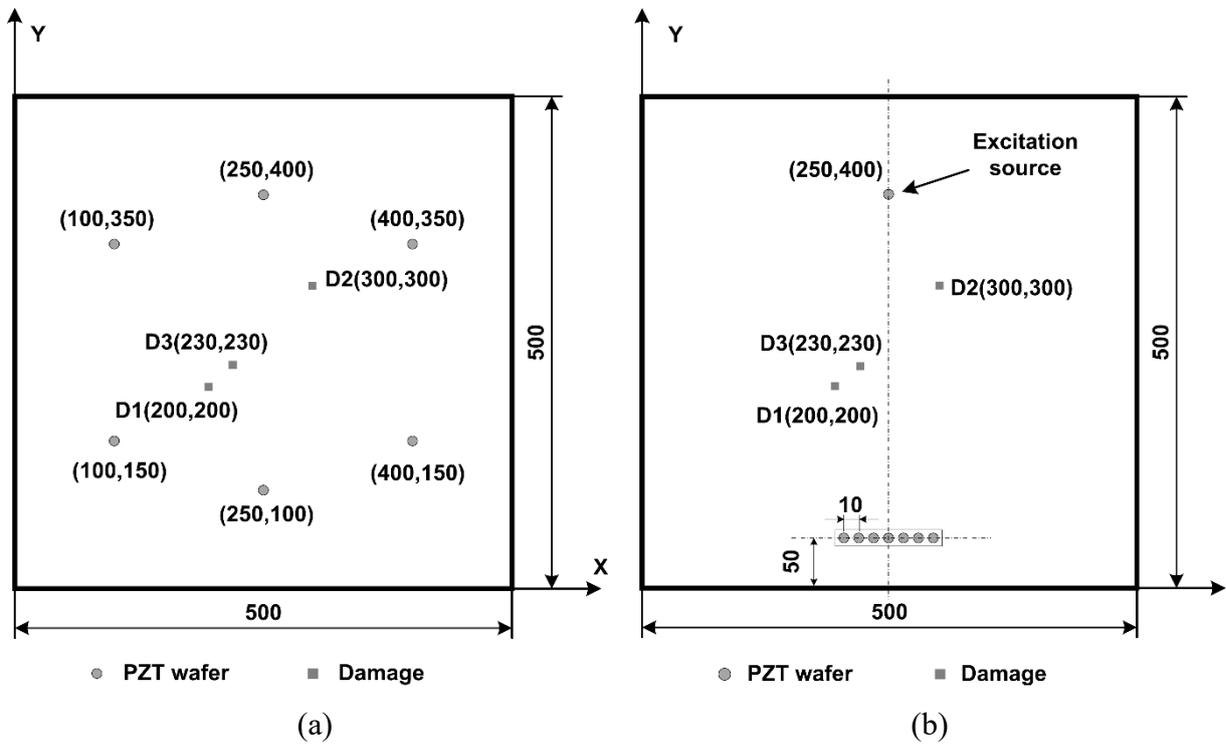


333
 334 **Fig. 3.** Key steps of Am-MUSIC algorithm.

335 **4. Numerical Validation**

336 To validate the developed Am-MUSIC algorithm for damage imaging, numerical simulation is
 337 implemented first. Consider a homogeneous, isotropic plate-like waveguide (density: $\rho=2,700$
 338 kg/m^3 ; Young modulus: $E=71 \text{ GPa}$; Poisson's ratio $\nu=0.33$), measuring $500 \text{ mm} \times 500 \text{ mm} \times 2$
 339 mm . Atop the waveguide there is a sparse sensor network with only six PZT wafers, as illustrated
 340 in **Fig. 4(a)**. Each PZT wafer functions as either a wave transmitter or a wave receiver, leading to
 341 15 transmitter-receiver pairs in the sensor network. For comparison against conventional MUSIC,
 342 another seven PZT wafers are arranged in a linear array as sensors, in **Fig. 4(b)**, along with an
 343 additional PZT wafer as wave actuator placed at the position (250 mm, 400 mm). In all cases, a
 344 3-cycle Hanning window tone burst with central frequency 200 kHz signal is selected as excitation
 345 signal to obtain S_0 wave mode, considering wave sensitivity and excitability. A total duration of
 346 $150 \mu\text{s}$ time length is analyzed for all numerical cases.

347



348 **Fig. 4.** Schematics of the plate waveguide in simulation (all dimensions in mm): (a) with a sparse sensor
 349 network for Am-MUSIC algorithm; and (b) with a linear array for conventional MUSIC algorithm.

350

351 Damage in the simulation is introduced to the waveguide by enforcing the material local stiffness
 352 to be zero. Three damage sites, labeled as D1-D3, are simulated in the waveguide, with respective
 353 positions highlighted in **Fig. 4** and summarized in **Table 1**. With these damage sites, three damage
 354 cases (C-I – C-III) are created by including different damage sites, **Table 2**.

355

Table 1 Three damage sites in simulation

Damage site	Position	
	x [mm]	y [mm]
D1	200	200
D2	300	300
D3	230	230

356

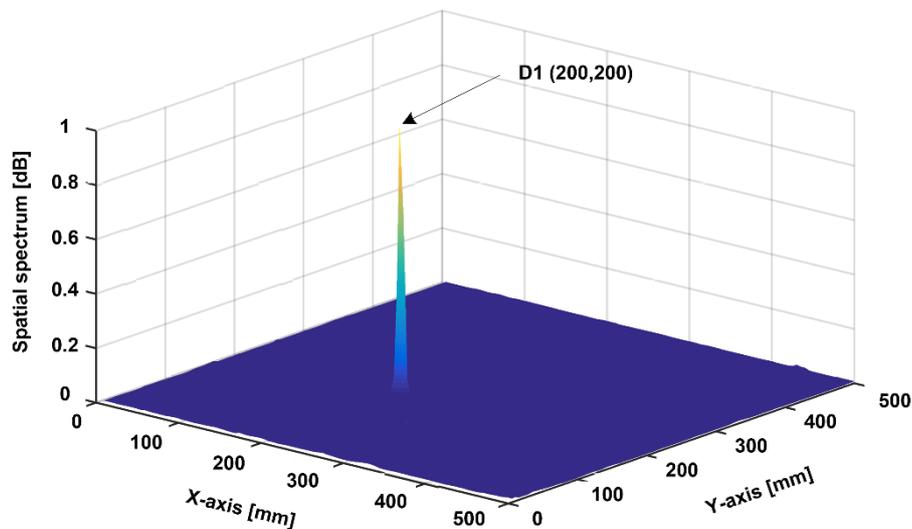
357

Table 2 Three damage cases in simulation

Damage case	Damage site included
C-I	D1
C-II	D1, D2
C-III	D1, D3

358

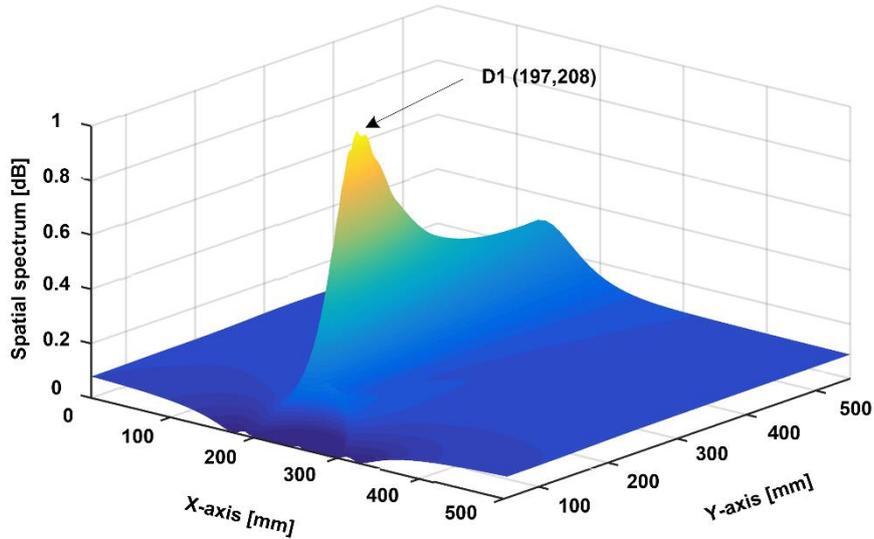
359 **Figure 5(a)** displays the spatial spectrum obtained using the Am-MUSIC algorithm, for C-I – the
 360 case with only a single damage site (D1), accurately pinpointing the damage location (200 mm,
 361 200 mm). For comparison, the image constructed using the conventional MUSIC algorithm is
 362 shown in **Fig. 5(b)**, indicating the damage location at (197 mm, 208 mm), which represents an
 363 error of (3 mm, 8 mm), in addition to an elongation artifact along the damage direction – a
 364 common deficiency for conventional MUSIC algorithms as illustrated elsewhere [21-25, 36]. The
 365 degree of such artifact depends on the signal-to-noise ratio and the point-spread function of the
 366 phased array at the location of the scatterer [37].



367

368

(a)

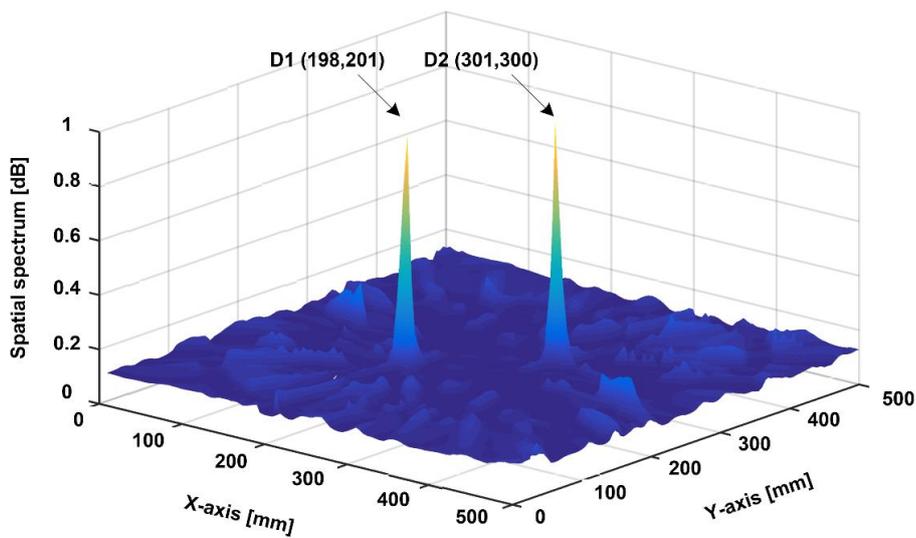


(b)

Fig. 5. Spatial spectra for C-I obtained by (a) Am-MUSIC algorithm; and (b) conventional MUSIC algorithm.

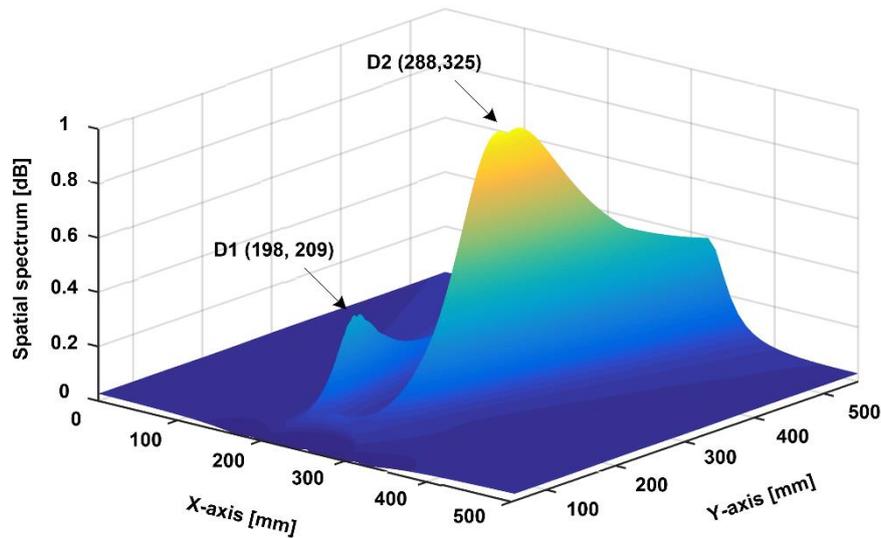
369
370
371
372
373

Figure 6(a) shows the spectrum for C-II obtained using the Am-MUSIC. Again, the identified results are observed to coincide exactly with actual damage sites, contrasting the spatial spectrum obtained using the conventional MUSIC algorithm in **Fig. 6(b)**, in which two damage sites are predicted with notable error. In addition, the peak at D1 location is much stronger than that at D2 as noted in **Fig. 6(b)**, contradicting the fact that D1 and D2 actually have the same degree of severity.



(a)

380
381

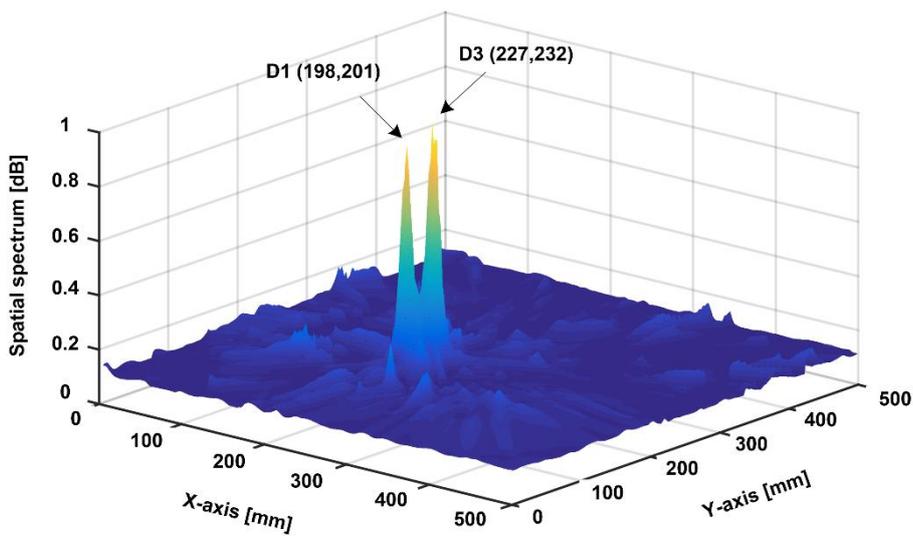


(b)

Fig. 6. Spatial spectra for C-II obtained by (a) Am-MUSIC algorithm; and (b) conventional MUSIC algorithm.

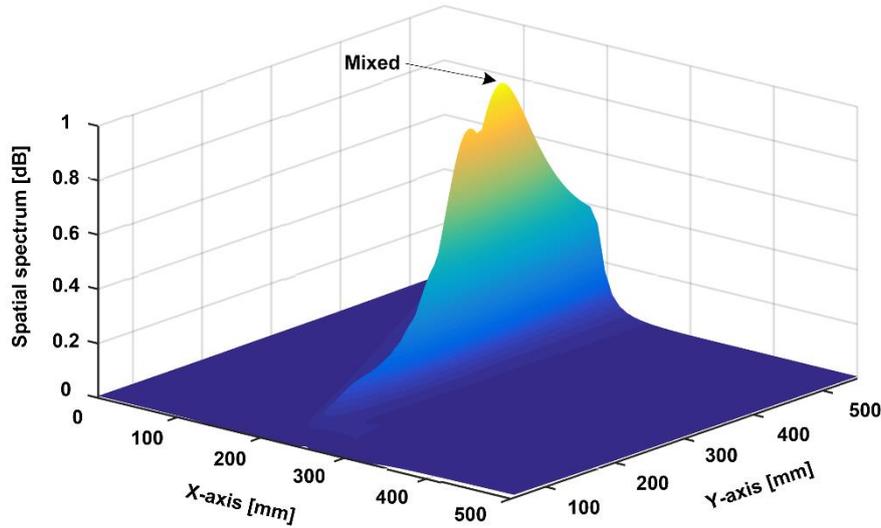
382
383
384
385
386

387 Provided that damage sites are in close proximity to each other – the case of C-III, the constructed
388 spatial spectra using the Am-MUSIC method and conventional MUSIC method are compared in
389 **Fig. 7.** In **Fig. 7(a)**, the two damage sites are localized precisely, in good agreement with the actual
390 positions; however, only one damage site is identified by the conventional MUSIC algorithm with
391 remarkable artifacts, in **Fig. 7(b)**, implying that the conventional MUSIC method may fail to
392 detect multiple damage sites which are close one to another.



(a)

393
394

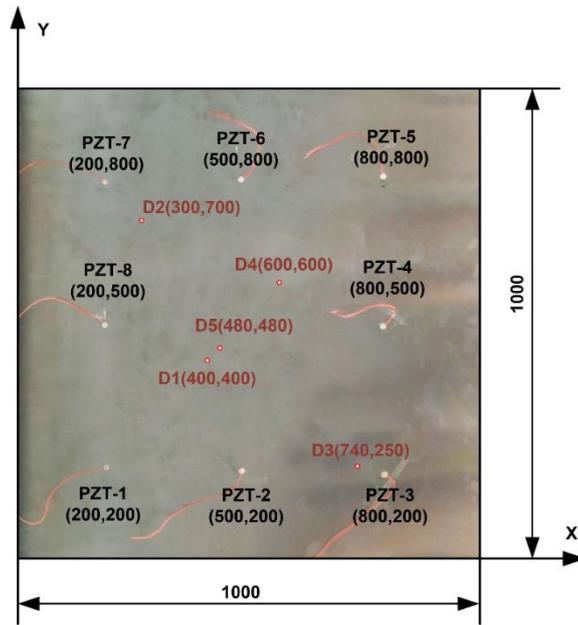


(b)

Fig. 7. Spatial spectra for C-III obtained by (a) Am-MUSIC algorithm; and (b) conventional MUSIC algorithm.

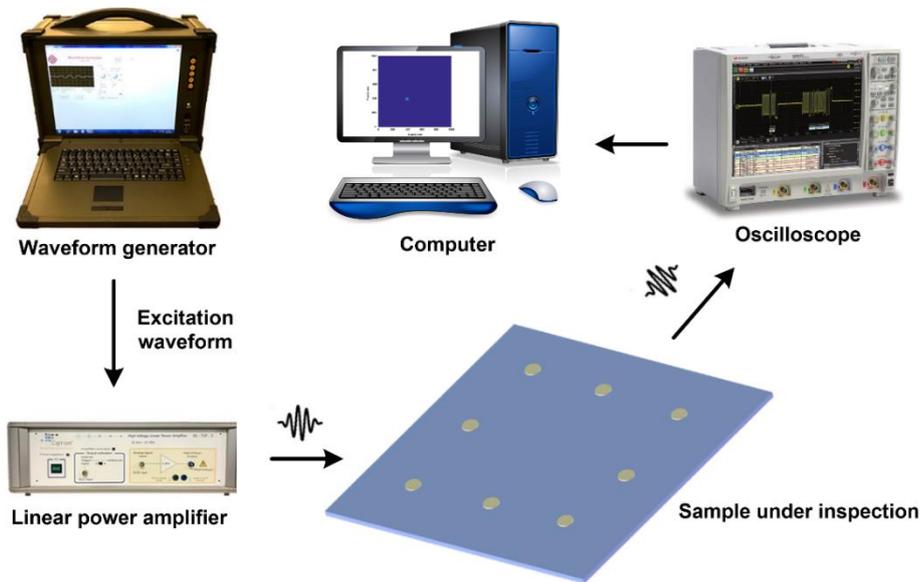
5. Experimental Validation

Subsequent to numerical simulation, effectiveness and accuracy of the Am-MUSIC-driven anomaly imaging is validated experimentally. A 2 mm-thick aluminum plate (dimensions: 1000 mm × 1000 mm × 2 mm; density: $\rho=2700 \text{ kg/m}^3$; Young modulus: $E=71 \text{ GPa}$; Poisson's ratio $\nu=0.33$) is prepared. A sparse sensor network, consisting of eight PZT wafers (labelled as PZT-1, PZT-2, ..., PZT-8), is surface-adhered on the plate, with the location of each wafer indicated in **Fig. 8(a)**. The experimental set-up is shown in **Fig. 8(b)**. The excitation signal – a Hanning-window-modulated 5-cycle toneburst at a central frequency of 200 kHz – is generated with an arbitrary waveform generator (NI[®] PXI-5412) and amplified by a linear power amplifier (Ciprian[®] US-TXP-3). The excitation signal is applied on each PZT wafer in turn to emit Lamb wave into the plate. S_0 mode Lamb wave signals, each in 300 μs , are acquired with a digital oscilloscope (NI[®] PXI-5105) at a sampling rate of 60 MHz.



(a)

413
414
415



(b)

416
417

Fig. 8. (a) An aluminum plate with a surface-adhered sparse sensor network consisting of eight PZT wafers in experiment (red ‘o’: actual damage and all dimensions in mm); and (b) Experimental set-up.

420

421 Similar to the simulation in Section 4, three categories of damage are considered in the experiment,
422 as recapped in **Table 3**.

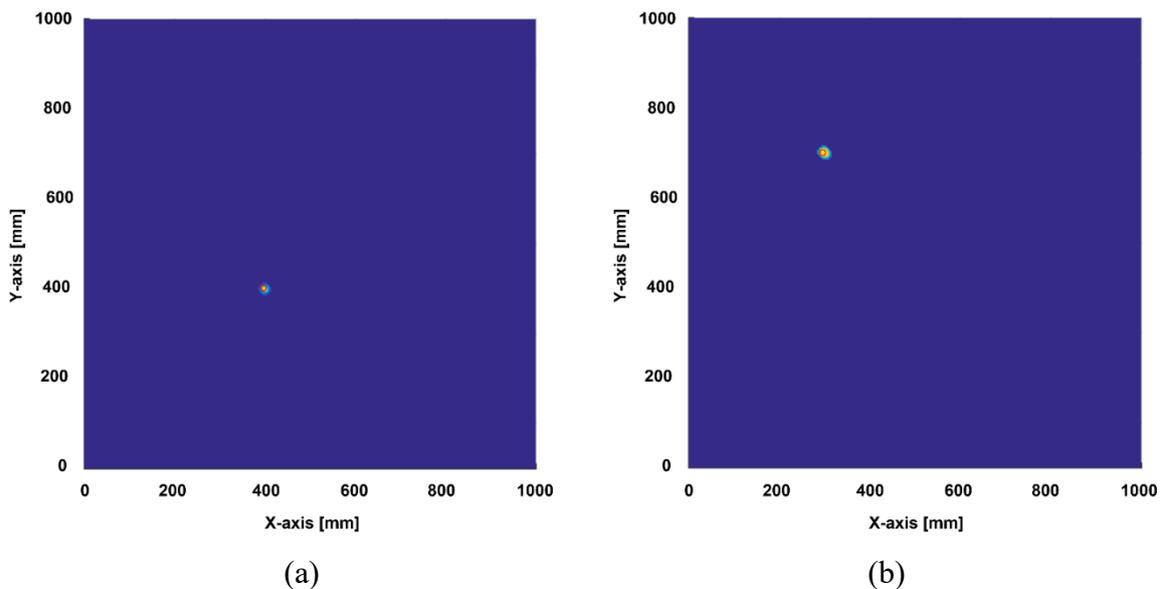
423

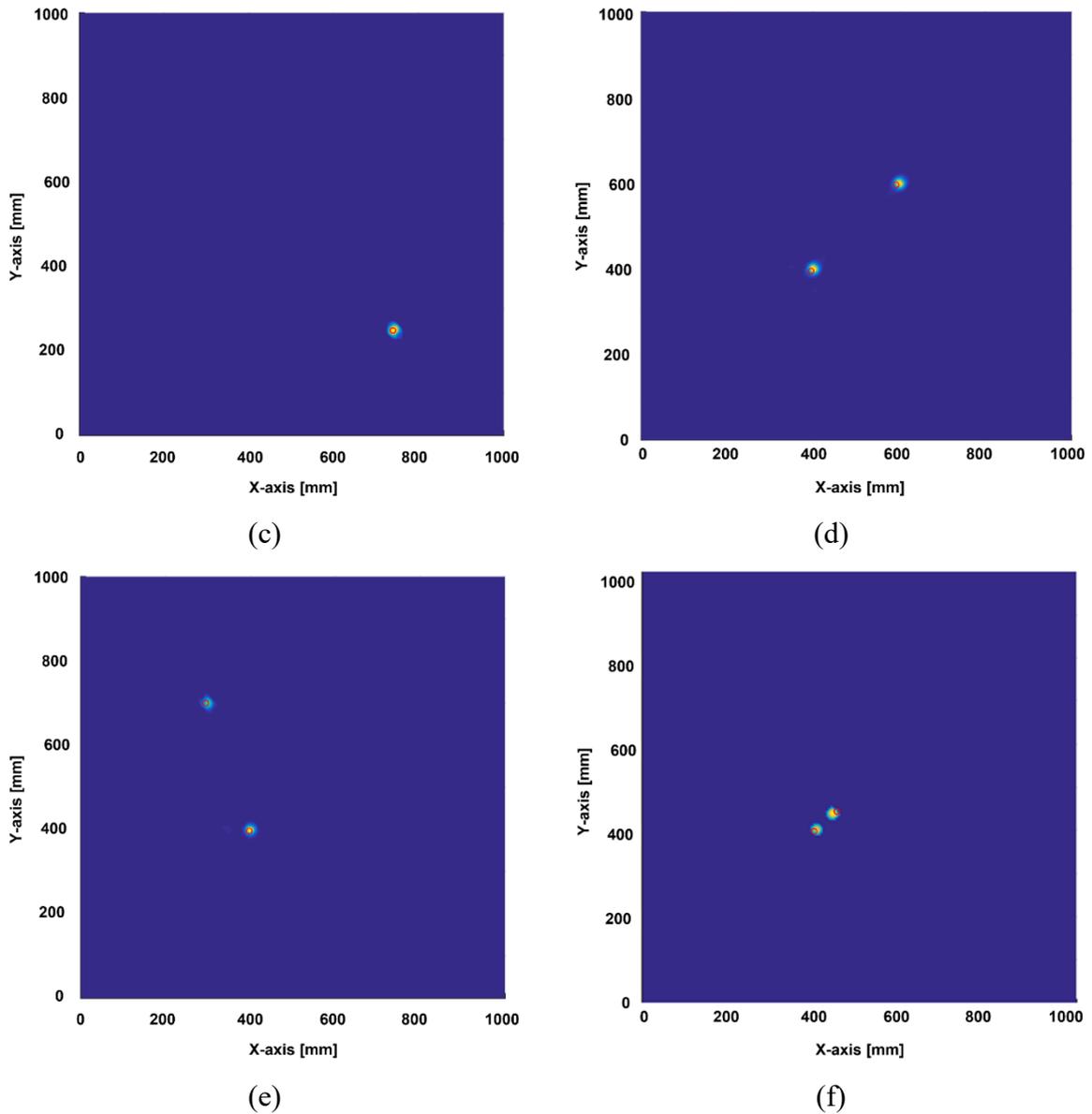
Table 3 Damage cases in experiment

Damage Category	Damage Case	Damage
Single damage	E-I	A through-hole D1 at (400 mm, 400 mm)

		(\O : 10 mm)
	E-II	A through-hole D2 at (300 mm, 700 mm) (\O : 10 mm)
	E-III	A through-hole D3 at (740 mm, 250 mm) (\O : 10 mm)
Double Damage (long distance apart)	E-IV	Two through-holes D1 at (400 mm, 400 mm) and D4 at (600 mm, 600 mm) respectively (\O : 10 mm)
	E-V	Two through-holes D1 at (400 mm, 400 mm) and D2 at (300 mm, 700 mm) respectively (\O : 10 mm)
Double Damage (short distance apart)	E-VI	Two through-holes D1 at (400 mm, 400 mm) and D5 at (480 mm, 480 mm) respectively (\O : 10 mm)

424
 425 The spatial spectra constructed using the Am-MUSIC algorithm for all the six damage cases are
 426 presented in **Fig. 9**, in which all damage sites are accurately located with precise depiction of the
 427 damage shape. In particular, for E-VI in which two damage sites are close one to the other, the
 428 algorithm still warrants high resolution and distinguishes individual damage sites.
 429

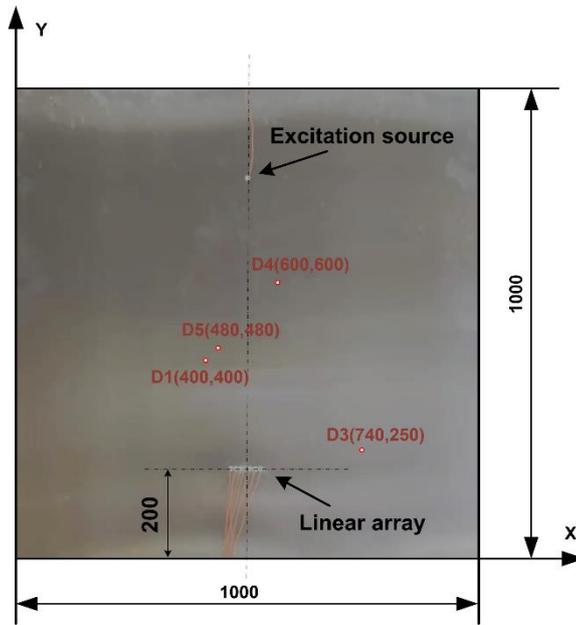




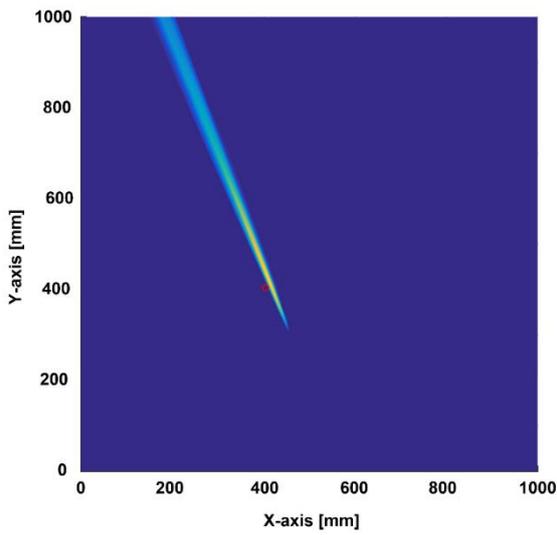
430 **Fig. 9.** Spatial spectra constructed using Am-MUSIC algorithm for damage case (a) E-I; (b) E-II; (c) E-
 431 III; (d) E-IV; (e) E-V; and (f) E-VI (red 'o': actual damage).

432
 433 To take a step further, conventional MUSIC algorithm in junction with the use of a linear array is
 434 recalled for comparison. Seven PZT wafers are configured in a linear array as receivers, in **Fig.**
 435 **10**, along with an additional PZT wafer as a wave actuator placed at the position (500 mm, 800
 436 mm). Four typical damage cases (E-I, E-III, E-IV, E-VI) are analysed. The spatial spectra
 437 constructed using the conventional MUSIC algorithm are shown in **Fig. 11**, showing inferior
 438 accuracy in damage localization and sizing; moreover, it fails to differentiate multiple damage
 439 sites in E-VI that are close one from the other, and also fails to identify damage D3 in E-III that
 440 is located in the blind zone for a conventional MUSIC algorithm.

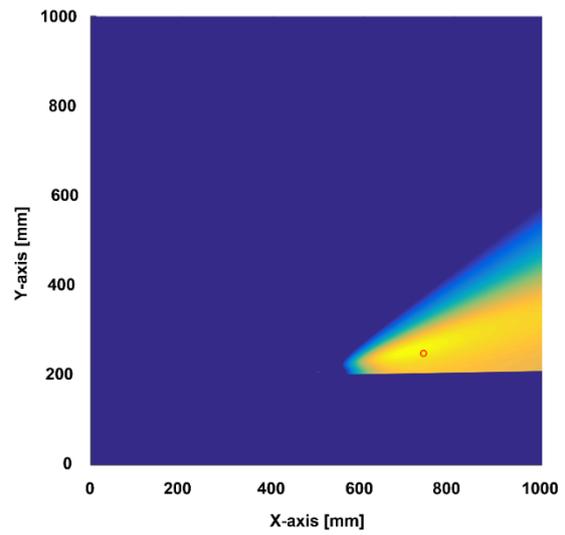
441



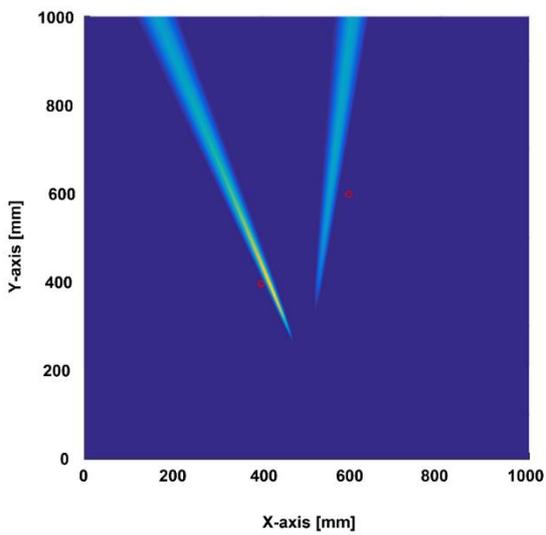
442
 443 **Fig. 10.** An aluminum plate with a surface-adhered linear array consisting of 7 PZT wafers in experiment
 444 (red 'o': actual damage and all dimensions in mm).



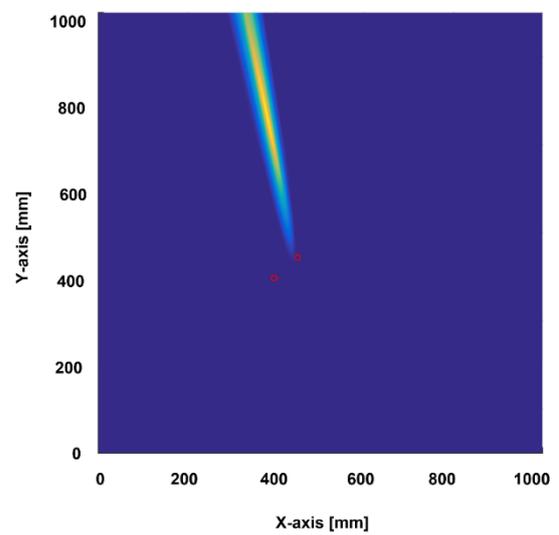
(a)



(b)



(c)



(d)

445 **Fig. 11.** Spatial spectra constructed using conventional MUSIC algorithm for damage case (a) E-I; (b) E-
446 III; (c) E-IV; and (d) E-VI (red 'o': actual damage).

448 **6. Concluding Remarks**

449 Aimed to circumvent some critical limitations of the conventional MUSIC algorithm-based
450 damage imaging, an ameliorated MUSIC algorithm is developed. In the Am-MUSIC algorithm,
451 the signal representation matrix at each pixel is manipulated by the excitation signal series, instead
452 of the scattered signal series, which enables the use of a sparse sensor network with arbitrarily
453 positioned transducers rather than a linear array featuring a dense configuration of transducing
454 elements with a uniform element pitch. By quantifying the orthogonal attributes between the
455 signal subspace and noise subspace inherent in the signal representation matrix, a full spatial
456 spectrum of the inspected sample can be generated, to visualize damage in the sample, irrespective
457 of the damage quantity. The effectiveness and accuracy of the Am-MUSIC algorithm are verified
458 in both simulation and experiment. Results show that compared with the conventional MUSIC
459 methods, the Am-MUSIC algorithm is capable of improving the detectability (in particular for
460 imaging of multiple damage sites that are close one to another) and eliminating blind zones. The
461 Am-MUSIC is conducive to expanding conventional MUSIC from phased array-facilitated
462 nondestructive evaluation to *in situ* health monitoring using built-in sparse sensor networks.

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