

Enhanced magnetoelectric effect in longitudinal-transverse mode Terfenol-D/Pb(Mg_{1/3}Nb_{2/3})O₃-PbTiO₃ laminate composites with optimal crystal cut

Yaojin Wang,^{1,2,3} Siu Wing Or,^{2,a)} Helen Lai Wa Chan,² Xiangyong Zhao,¹ and Haosu Luo[†]

¹State Key Laboratory of High Performance Ceramics and Superfine Microstructure, Shanghai Institute of Ceramics, Chinese Academy of Sciences, 215 Chengbei Road, Jiading, Shanghai 201800, People's Republic of China

²Department of Applied Physics, The Hong Kong Polytechnic University, Hung Hom, Kowloon, Hong Kong

³Graduate School of the Chinese Academy of Sciences, Beijing 10039, People's Republic of China

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Magnetoelectric (ME) laminate composite consisting of optimal crystal cut thickness-polarized piezoelectric 0.7Pb(Mg_{1/3}Nb_{2/3})O₃-0.3PbTiO₃ (PMN-PT) single crystal and the length-magnetized magnetostrictive Tb_{0.3}Dy_{0.7}Fe_{1.92} (Terfenol-D) alloy has been fabricated. The cut optimization of PMN-PT crystal greatly enhances the longitudinally magnetized-transversely polarized (*L-T*) mode ME effect, which has a superior ME voltage coefficient α_E of ~ 3.02 V/cm Oe in low frequency band. Near the resonance frequency of 95 kHz, the coefficient dramatically increases and reaches the maximized value of 33.2 V/cm Oe, which is almost two times larger than the previously reported $\langle 001 \rangle$ -oriented PMN-PT crystal based *L-T* mode laminate composite. © 2008 American Institute of Physics. [DOI: 10.1063/1.2943267]

I. INTRODUCTION

The magnetoelectric (ME) effect is the polarization P response to an applied magnetic field H , and the converse ME effect is a magnetization M response to an applied electric field E .¹ In the past decades, considerable research efforts have been put on the ME effect, first in single-phase materials, then in two-phase bulk composites, and lately in two-/three-phase laminate composites.²⁻⁹ The magnetic-to-electric field conversion in ME composite is a magnetoelasto-electric coupling.³ The ME coupling in laminate composite is realized by a stress mediated interaction between the magnetostrictive phase (or ferromagnet) and the piezoelectric (or ferroelectric) phase, which is often referred to as a multiferroic composite.⁴ Some simple architecturally engineered nanostructured composites, such as CoFe₂O₄ nanopillars in a BaTiO₃ matrix, have also recently been reported.⁵ It seems that these nanocomposites should have potential for high ME coupling because the two phase materials have a more intimate contact in nanodimension. But these nanostructured composites suffer from the drawback that the magnetostrictive and piezoelectric effects are dramatically decreased in nanodimension due to substrate's clamping effect.⁶ To date, it is known that laminated composites of piezoelectric 0.7Pb(Mg_{1/3}Nb_{2/3})O₃-0.3PbTiO₃ (PMN-PT) single crystal and magnetostrictive Tb_{0.3}Dy_{0.7}Fe_{1.92} (Terfenol-D) alloy possess superior ME effect and ultrahigh magnetic field sensitivity due to their greater product effect of the piezoelectric effect and the magnetostrictive effect.⁷ However, the orientation of piezoelectric PMN-PT single crystal was out of consideration in previous studies. Actually, based on the

working principle of *L-T* mode laminate composite, the Terfenol-D mainly vibrates longitudinally when the applied magnetic field is along the length direction. Due to the mechanical coupling, the PMN-PT single crystal has to vibrate synchronously along the length direction and generates electric charge along the thickness direction. So the piezoelectric coefficient d_{31} and electromechanical coupling factor k_{31} should contribute greatly to the ME effect.⁷ Consequently, designing new optimal cut-types in the single crystal piezoelectric phase to produce optimized transverse piezoelectric performance is a new method to enhance the applications of ME effect.

In this paper, we have developed an *L-T* mode ME laminate composite by sandwiching a special-oriented piezoelectric PMN-PT single crystal between two longitudinally magnetized magnetostrictive Terfenol-D plates.

II. EXPERIMENTS

A. Structure and fabrication

Figure 1 shows schematic diagram of the proposed sand-

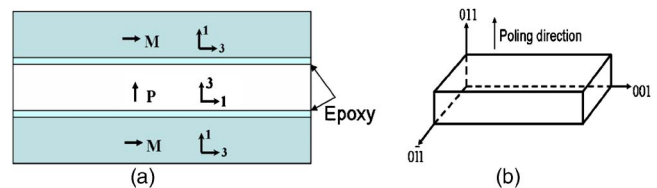


FIG. 1. (Color online) (a) Schematic diagram of the proposed magnetostrictive/piezoelectric laminated composites. The arrows designate the magnetization and polarization directions, respectively. (b) The newly designed special-oriented PMN-PT single crystal. The $\langle 001 \rangle$, $\langle 01\bar{1} \rangle$, and $\langle 011 \rangle$ crystallographic axes are oriented in the length, width, and thickness directions, respectively.

^{a)}Author to whom correspondence should be addressed. Electronic mail: apswor@polyu.edu.hk.

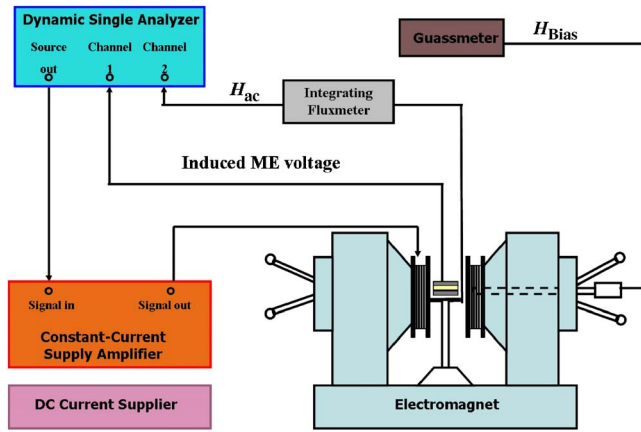


FIG. 2. (Color online) Schematic diagram of the ME measurement setup.

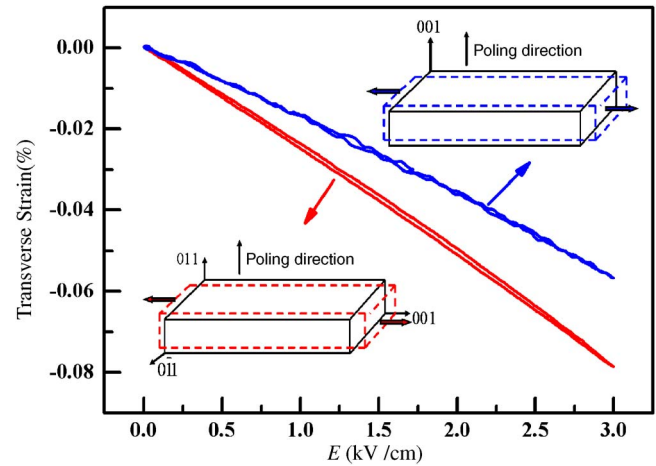
wiched composite structure with the ME effect and the optimal cut-type of PMN-PT plate. The high quality PMN-PT single crystal was grown directly from the melt by modified Bridgman technique.¹⁰ The as-grown single crystals were oriented along $\langle 001 \rangle$, $\langle 011 \rangle$, and $\langle 0\bar{1}1 \rangle$ directions using an x-ray diffraction meter, and then diced to prepare samples with the crystal cut and dimensions of $12\langle 001 \rangle^L \times 6\langle 0\bar{1}1 \rangle^W \times 1\langle 011 \rangle^T \text{ mm}^3$ (L : length, W : width, T : thickness). It has been shown in our previous work that PMN-PT plates, with this specially cut and poled along the $\langle 011 \rangle$ thickness direction, possess ultrahigh transverse piezoelectric performance, i.e., ultrahigh thickness direction voltage response to length direction strain deformation.¹¹ The Terfenol-D plates were commercially supplied (Baotou Research Institute of Rare Earth, China) with the same dimensions as the PMN-PT plate and with the length direction along the $[112]$ direction, which is the highly magnetostrictive crystallographic axis and the magnetization (M) was relatively easy to achieve.¹²

B. Measurement setup and procedure

The ME properties of the laminate composite were characterized at room temperature and zero stress bias using an in-house automated measurement system¹³ shown in Fig. 2. The ME voltages (V) induced in the composites were measured as a function of ac magnetic field (H_{ac}), dc magnetic bias (H_{bias}), and ac magnetic field frequency (f) in the ranges of 10^{-7} – 10^{-3} T, 0–1200 Oe, and 1–100 kHz, respectively. H_{ac} was provided by Helmholtz coils driven by a dynamic signal analyzer (Ono Sokki CF5220) via a constant-current supply amplifier (AE Techron 7572). H_{bias} was supplied by a water-cooled, U-shaped electromagnet (Mytem PEM-8005K) controlled by a dc current supply (Sorensen DHP200-15). H_{ac} and H_{bias} were monitored *in situ* by a pick-up coil connected to an integrating fluxmeter (Walker MF-10D) and a Guassmeter (F. W. Bell 7030), respectively. All quantities were sampled and recorded by the dynamic signal analyzer and stored in a computer.

III. RESULTS AND DISCUSSION

It is well known that PMN-PT single crystal has superior piezoelectric effect and electromechanical coupling perfor-

FIG. 3. (Color online) Strain vs E -field (unipolar) curves in the $\langle 001 \rangle$ -oriented and newly designed cut-type PMN-PT crystals.

mance. However, the piezoelectric effect of PMN-PT is anisotropic, depending significantly on the crystal cut type and the poling direction. Consequently, optimization of the crystal cut and poling processing of PMN-PT single crystals should be considered to produce optimal piezoelectric performance for different piezoelectric resonator modes in versatile applications. Figure 3 shows the electric-field induced strain patterns for the $\langle 001 \rangle$ -oriented and newly designed cut-type 0.70PMN-0.30PT crystal, respectively, using a unipolar field with amplitude of $E < 3$ kV/cm and frequency of 0.2 Hz. From the slope of the plot, the piezoelectric coefficients of d_{31} were determined. The k_{31} was determined using the resonance-antiresonance technique by a HP 4194A impedance analyzer following the IEEE standards, and d_{33} is directly measured by a quasistatic Berlincourt d_{33} meter (50 Hz) (see Table I). From the comparison in Table I, it is obviously observed that the newly designed cut-type PMN-PT crystal is much superior for transverse mode application than the $\langle 001 \rangle$ -oriented one.

The ME voltage coefficient α_E , defined as $|dE/dH_{ac}|$, of the as-prepared crystal cut optimized laminate composite was then measured for various H_{bias} at H_{ac} of 1 Oe peak and f of 1 kHz, as shown in Fig. 4. α_E initially increases rapidly with H_{bias} and reaches a maximum value of ~ 3.02 V/cm Oe at an optimal H_{bias} of 400 Oe, then decreases with increasing H_{bias} . In addition, the laminated composite has almost a linear response to H_{bias} in the range of $0 < H_{bias} < 200$ Oe. Correspondingly, the relationship of the laminated composite can be used to detect small dc magnetic field.⁸

Figure 5 illustrates the induced ME voltage V across the PMN-PT plate as a function of applied H_{ac} over the range of $10^{-7} < H_{ac} < 10^{-3}$ T at $f = 1$ kHz. It is clear that V has an

TABLE I. Piezoelectric parameters for $\langle 001 \rangle$ -oriented and newly designed cut-type PMN-PT crystals.

	d_{31} ($\times 10^{-12}$ C/N)	d_{33} ($\times 10^{-12}$ C/N)	k_{31}	$\varepsilon_{33}^T/\varepsilon_0$
$\langle 001 \rangle$ -oriented	-1126	2420	0.62	4276
$[001]^L \times [\bar{1}10]^W \times [110]^T$ -oriented	-2645	2060	0.95	6200

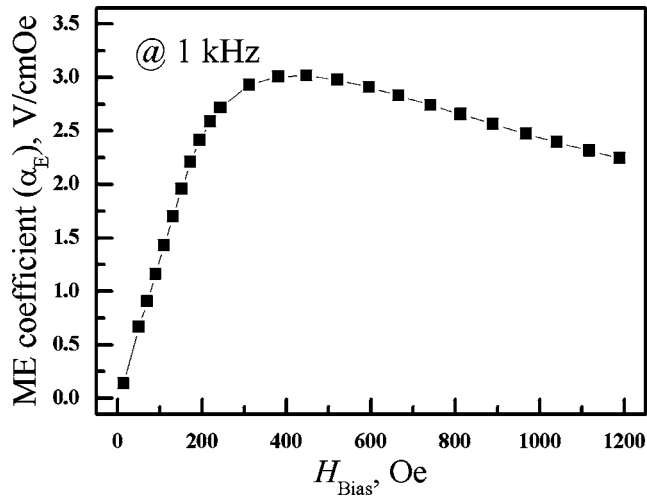


FIG. 4. ME voltage coefficient α_E of the new laminated composite as a function of H_{bias} . These data were taken using $H_{\text{ac}}=1$ Oe peak at $f=1$ kHz.

excellent linear response to H_{ac} in the whole measured range for various H_{bias} . A higher detection sensitivity of 10^{-10} – 10^{-11} T could be obtained if shielding magnetic noise could be adopted and composite fabrication could be improved.⁷ From the slope of the plot, the ME voltage coefficient α_E for various H_{bias} can also be determined, which coincides reasonably well with α_E versus H_{bias} , as shown in Fig. 4.

The dependence of α_E on f at various H_{bias} is shown in Fig. 6(a). It is noted that no remarkable dispersion of α_E is observed for all cases apart from the variations associated with the ME resonances. The largest α_E is observed at $H_{\text{bias}}=400$ Oe for the whole frequency range. In particular, the maximal resonance α_E located at the ME resonances frequency (f^{MER}) of 95.25 kHz under this optimal H_{bias} is as large as 33.2 V/cm Oe. This resonance α_E at f^{MER} is over 13 times larger than its nonresonance α_E of ~ 3.02 V/cm Oe. Figure 6(b) shows H_{bias} dependence of f^{MER} and α_E at f^{MER} . The results clearly demonstrate that the variation of α_E at f^{MER} is similar to the α_E at nonresonant frequency, as shown in Fig. 4. However, f^{MER} decreases with increasing H_{bias} ,

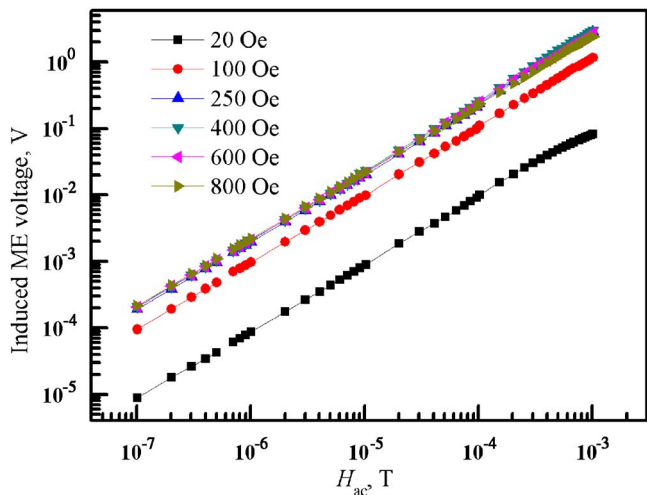


FIG. 5. (Color online) Induced ME voltage V as a function of applied H_{ac} over the range of $10^{-7} < H_{\text{ac}} < 10^{-3}$ T at $f=1$ kHz for various H_{bias} .

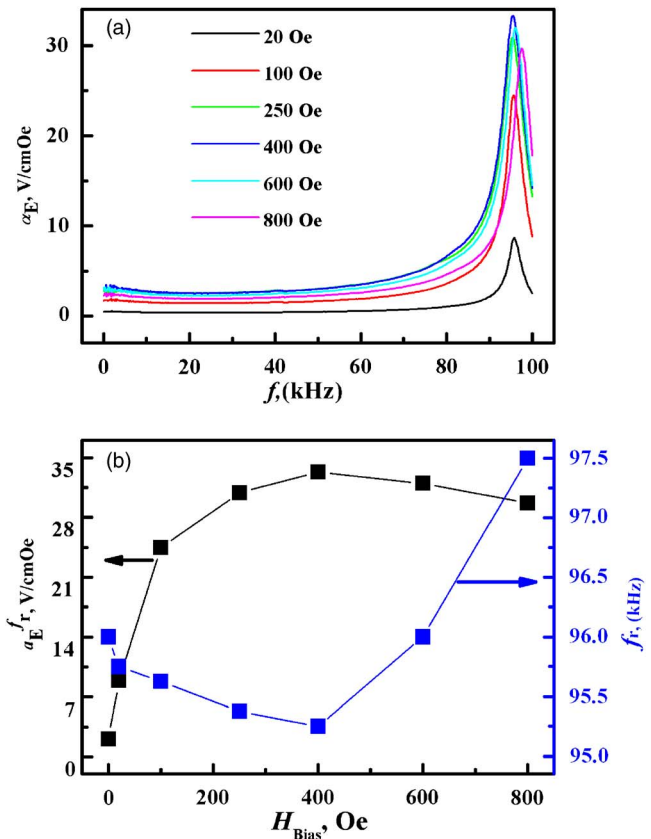


FIG. 6. (Color online) (a) Frequency response of ME voltage coefficient α_E for various H_{bias} , and (b) H_{bias} dependence of ME voltage coefficient at ME resonances α_E^{MER} and ME resonances frequency f_r .

reaching the smallest value of 95.25 kHz at $H_{\text{bias}}=400$ Oe, and then increasing with increasing H_{bias} . Physically, the initial change in both α_E (at f^{MER}) and f^{MER} with increasing H_{bias} can be explained by the H_{bias} -induced motion of the available non- 180° domain walls in the Terfenol-D plates.¹³ That is, as H_{bias} increases to the 400 Oe critical value, the compliance associated with increased deformation contribution from this non- 180° domain-wall motion is maximized, resulting in a maximum in strain (and hence α_E at nonresonant frequency and α_E at f^{MER}) and a minimum in stiffness (and hence f^{MER}). Beyond this optimal and also critical value of H_{bias} , constraining of non- 180° domain-wall motion due to interaction with H_{bias} gives rise to a decrease in strain and an increase in stiffness. It is noted that the deformation contribution from the motion of 180° domain walls is insignificant as it produces changes in magnetization without accompanying strain.⁹

IV. SUMMARY

In summary, our specially designed L - T mode ME laminate composite by sandwiching PMN-PT plate with the optimal cut type of $\langle 001 \rangle^L \times \langle 0\bar{1}1 \rangle^w \times \langle 011 \rangle^T$ between two Terfenol-D plates enhances the elastoelectric coupling in the magnetic-to-electric field conversion, which results in a significantly increased ME voltage coefficient, approximately two times larger than the previous L - T mode laminate com-

posite of $\langle 001 \rangle$ -oriented PMN-PT single crystal. This reveals a new method to enhance the ME effect of the ME composite.

ACKNOWLEDGMENTS

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- ¹L. D. Landau and E. Lifshitz, *Electrodynamics of Continuous Media* (Pergamon, Oxford, 1960), p. 119.
- ²V. J. Folen, G. T. Rado, and E. W. Stalder, *Phys. Rev. Lett.* **6**, 607 (1961).
- ³C. W. Nan, *Phys. Rev. B* **50**, 6082 (1994).
- ⁴G. Srinivasan, E. T. Rasmussen, J. Gallegos, R. Srinivasan, Y. I. Bokhan, and V. M. Laletin, *Phys. Rev. B* **64**, 214408 (2001).
- ⁵H. Zheng, J. Wang, S. E. Lofland, A. Ma, L. Mohaddes-Ardabili, T. Zhao, L. Salamanca-Riba, S. R. Shinde, S. B. Ogale, F. Bai, D. Viehland, Y. Jia, D. G. Schlom, M. Wuttig, A. Roytburd, and R. Ramesh, *Science* **303**, 661 (2004).
- ⁶V. M. Petrov, G. Srinivasan, M. I. Bichurin, and A. Gupta, *Phys. Rev. B* **75**, 224407 (2007).
- ⁷S. X. Dong, J. F. Li, and D. Viehland, *Appl. Phys. Lett.* **83**, 2265 (2003).
- ⁸S. X. Dong, J. Y. Zhai, J. F. Li, and D. Viehland, *Appl. Phys. Lett.* **88**, 082907 (2006).
- ⁹T. L. Li, S. W. Or, and H. L. W. Chan, *J. Magn. Magn. Mater.* **304**, 442 (2004).
- ¹⁰H. S. Luo, G. S. Xu, P. C. Wang, H. Q. Xu, and Z. W. Yin, *J. Appl. Phys.* **39**, 5581 (2000).
- ¹¹J. Peng, H. S. Luo, D. Lin, H. Q. Xu, T. H. He, and W. Q. Jin, *Appl. Phys. Lett.* **85**, 6221 (2004).
- ¹²G. Engdahl, *Magnetostrictive Materials Handbook* (Academic, New York, 2000).
- ¹³S. W. Or, T. Li, and H. L. W. Chan, *J. Appl. Phys.* **97**, 10M308 (2005).