## High magnetoelectric effect at low magnetic basing in heterostructure rod of magnetostrictive fibers, piezoelectric tube, and epoxy binder

Cite as: J. Appl. Phys. **117**, 17D721 (2015); https://doi.org/10.1063/1.4918574 Submitted: 23 September 2014 . Accepted: 21 December 2014 . Published Online: 16 April 2015

Chung Ming Leung, and Siu Wing Or



### **ARTICLES YOU MAY BE INTERESTED IN**

Voltage-mode direct-current magnetoelectric sensor based on piezoelectricmagnetostrictive heterostructure Journal of Applied Physics **117**, 17A748 (2015); https://doi.org/10.1063/1.4919047

Multiferroic magnetoelectric composites: Historical perspective, status, and future directions Journal of Applied Physics **103**, 031101 (2008); https://doi.org/10.1063/1.2836410

Direct current force sensing device based on compressive spring, permanent magnet, and coil-wound magnetostrictive/piezoelectric laminate Review of Scientific Instruments **84**, 125003 (2013); https://doi.org/10.1063/1.4838615





J. Appl. Phys. **117**, 17D721 (2015); https://doi.org/10.1063/1.4918574 © 2015 AIP Publishing LLC.



# High magnetoelectric effect at low magnetic basing in heterostructure rod of magnetostrictive fibers, piezoelectric tube, and epoxy binder

Chung Ming Leung and Siu Wing Or<sup>a)</sup>

Department of Electrical Engineering, The Hong Kong Polytechnic University, Hung Hom, Kowloon, Hong Kong

(Presented 5 November 2014; received 23 September 2014; accepted 21 December 2014; published online 16 April 2015)

A three-phase heterostructure rod of 35 mm length and 8 mm diameter is developed by embedding 18 pieces of length-magnetized Tb<sub>0.3</sub>Dy<sub>0.7</sub>Fe<sub>1.92</sub> magnetostrictive continuous fibers in a wall-thickness-polarized Pb(Zr, Ti)O<sub>3</sub> piezoelectric long-thin tube using an epoxy binder in order to study its magnetoelectric (ME) effect for passive sensing of axial magnetic fields. The heterostructure rod exhibits a high non-resonance ME voltage coefficient ( $\alpha_V$ ) of 38 mV/Oe for frequencies up to 10 kHz and an ultrahigh resonance  $\alpha_V$  of 680 mV/Oe at 32 kHz, both under a reasonably low magnetic bias field ( $H_{\text{Bias}}$ ) of 120 Oe. The simultaneously high axial ME responses at the reasonably low  $H_{\text{Bias}}$ , as confirmed theoretically and experimentally, can be ascribed to the realization of the relatively true magnetostrictive 33-mode and piezoelectric 31-mode of operation in the magnetostrictive continuous fibers and piezoelectric long-thin tube, respectively, as well as the reduction in the effect of demagnetization fields in the high aspect-ratio magnetostrictive continuous fibers. © 2015 AIP Publishing LLC. [http://dx.doi.org/10.1063/1.4918574]

#### I. INTRODUCTION

The magnetoelectric (ME) effect extrinsic in magnetostrictive-piezoelectric laminates has been widely investigated for passive magnetic field sensing over the past decade.<sup>1–5</sup> One of the largest obstacles to practically realize ME sensors is the need of a high magnetic bias field ( $H_{\text{Bias}}$ ), reaching 400-600 Oe in the well-known L-T laminated plates and even  $\sim$ 5000 Oe in the T-T laminated plates, for maximizing the magnetostrictive response for the piezoelectric phase.<sup>1,2</sup> The need of a high  $H_{\text{Bias}}$  implies the need of big and strong permanent magnets in the resulting sensors, which in turn increases the size, mass, magnetic influence inside and outside the sensors, and packaging difficulty.<sup>6</sup> Today, reducing  $H_{\text{Bias}}$  without reducing sensitivity remains a challenging topic to the ME community. Among several published works, the one reporting heterostructures of metglas and piezoelectric fibers has demonstrated a very high ME voltage coefficient ( $\alpha_V$ ) of 1.5 V/Oe at a non-resonance frequency of 1 kHz under an extremely low  $H_{\text{Bias}}$  of ~8 Oe.<sup>7</sup> However, this extremely low  $H_{\text{Bias}}$  not only constrains the heterostructures to small-signal applications but also is easily affected by environmental noises.

In this paper, we report theoretically and experimentally the existence of a high axial ME effect at a reasonably low  $H_{\text{Bias}}$  of 120 Oe in a three-phase heterostructure rod, comprising 18 pieces of length-magnetized Tb<sub>0.3</sub>Dy<sub>0.7</sub>Fe<sub>1.92</sub> magnetostrictive continuous fibers bonded with an epoxy binder in a wall-thickness-polarized Pb(Zr, Ti)O<sub>3</sub> piezoelectric long-thin tube. The use of the magnetostrictive continuous fibers and piezoelectric long-thin tube are aimed to enable the relatively true magnetostrictive 33-response and piezoelectric 31response in the heterostructure rod, respectively. Besides, the high length-to-width aspect ratio (35:1 in our case) of the magnetostrictive continuous fibers can effectively reduce the effect of demagnetization fields, thereby providing a lower field of operation. The use of the epoxy binder is to provide a good mechanical mediation for the magnetostrictive continuous fibers and the piezoelectric long-thin tube. From another view, our heterostructure rod can be considered as a "rod-in-tube" ME rod in which an epoxy-bonded magnetostrictive continuous fiber composite rod is situated in the middle of a piezoelectric ceramic long-thin tube. The creation of the epoxy-bonded magnetostrictive continuous fiber composite rod can also mitigate the eddy-current losses and mechanical brittleness intrinsic in most of the monolithic magnetostrictive materials. As a result, our ME rod exhibits simultaneously high non-resonance and resonance axial ME responses at a reasonably low  $H_{\text{Bias}}$  of 120 Oe.

#### **II. STRUCTURE AND WORKING PRINCIPLE**

Figure 1(a) illustrates the schematic diagram of the proposed three-phase heterostructure rod with an axial magnetic field sensing mode configuration. The heterostructure rod, having 35 mm length and 8 mm diameter, consists of a total of 18 pieces of length-magnetized  $Tb_{0.3}Dy_{0.7}Fe_{1.92}$  magnetostrictive continuous fibers inserted inside a wall-thickness-polarized Pb(Zr, Ti)O<sub>3</sub> piezoelectric long-thin tube and bonded using an epoxy binder. It can also be considered as a "rod-in-tube" ME rod in which the central portion of magnetostrictive continuous fibers and epoxy binder essentially form an epoxy-bonded magnetostrictive continuous fiber composite rod. In fabrication, the piezoelectric long-thin tube was used as the mold, and the epoxy-bonded magnetostrictive continuous fiber composite rod was prepared inside

<sup>&</sup>lt;sup>a)</sup>Author to whom correspondence should be addressed. Electronic mail: eeswor@polyu.edu.hk.



FIG. 1. (a) Schematic diagram of the proposed three-phase heterostructure rod (or ME rod) with an axial magnetic field sensing mode configuration. (b) Dynamic magneto-mechano-electric equivalent circuit of the proposed ME rod.

the mold under vacuum to ensure a good structural integrity. In more detail, a Ferroperm Pz26 hard Pb(Zr, Ti)O<sub>3</sub> piezoelectric long-thin tube, with dimensions of 35 mm length  $(l) \times 8 \,\mathrm{mm}$  outer diameter  $(2 \,r_{\rm p}) \times 6 \,\mathrm{mm}$  inner diameter  $(2r_{\rm m})$ , full-fired silver electrodes on the inner and outer circumferential surfaces normal to the wall-thickness direction, and an electric polarization (P) along the wall-thickness direction, was commercially acquired. Magnetostrictive continuous fibers with a square cross-section of 35 mm long and 1 mm width and the highly magnetostrictive [112] crystallographic axis oriented along their length direction were cut along the length of a Baotou Rare Earth [112]-textured monolithic Tb<sub>0.3</sub>Dy<sub>0.7</sub>Fe<sub>1.92</sub> plate using a wire electrical discharge machining technique. With the piezoelectric longthin tube mold being sealed at one end, a total of 18 pieces of magnetostrictive continuous fibers were placed inside the mold under a dc magnetic alignment field of  $\sim$ 3000 Oe set up by a pair of NdFeB magnets. The predegassed Araldite LY564/HY2954 epoxy was transferred into the mold and degassed again inside the mold, both under vacuum. The other end of the mold was sealed, and the epoxy was allowed to cure at 80 °C for 9 h inside an oven. The volume fraction of the magnetostrictive continuous fibers was theoretically determined to be  $\sim 0.63$  in the epoxy-bonded magnetostrictive continuous fiber composite rod, based on the rule-ofmixture formulation for density. In fact, epoxy-bonded Tb<sub>0.3</sub>Dy<sub>0.7</sub>Fe<sub>1.92</sub>-based composites with Tb<sub>0.3</sub>Dy<sub>0.7</sub>Fe<sub>1.92</sub> volume fraction in excess of 0.5 have been found to be optimal for magnetostrictive and ME applications by taking the balance between the maximization of magnetostrictive response and the minimization of eddy-current losses and material cost into consideration.8,9

The axial magnetic field sensing in the ME rod in Fig. 1(a) can be described as follows. An ac magnetic field

 $(H_3)$  applied along the axial (or length) direction of the ME rod causes the inner epoxy-bonded magnetostrictive continuous fiber composite rod to produce axial magnetostrictive strains in its length direction due to the magnetostrictive 33-mode effect. Since the epoxy-bonded magnetostrictive continuous fiber composite rod is mechanically bonded inside the piezoelectric long-thin tube, these axial magnetostrictive strains will subsequently stress on the outer piezoelectric long-thin tube to generate an ac voltage  $(V_3)$  across its wall-thickness in the wall-thickness direction because of the piezoelectric 31-mode effect. Using the equation of motion to couple the magnetostrictive and piezoelectric constitutive relations,<sup>10,11</sup> applying Newton's second law of motion to the ME rod,<sup>10</sup> taking free-free mechanical boundary conditions at the two ends of the ME rod,<sup>12</sup> considering mechanical losses and internal stresses,<sup>13,14</sup> and finding analogous electrical equivalent parameters,<sup>12</sup> a dynamic magneto-mechano-electric equivalent circuit of the proposed ME rod, shown in Fig. 1(a), was obtained, and is shown in Fig. 1(b). Therefore, the ME voltage coefficient  $(\alpha_V)$  of the ME rod can be determined as

$$\begin{aligned} \alpha_V &= \left| \frac{dV_3}{dH_3} \right| = \left| \frac{-2\varphi_m \varphi_p \beta}{2\varphi_p^2 + j\omega C_0(Z_1 + 2Z_2 + R)} \right|, \quad (1) \\ \varphi_m &= \frac{A_m d_{33,m}}{s_{33}^H}, \quad \varphi_p = \frac{2\pi (r_p - r_m) d_{31,p}}{s_{11}^E}, \\ C_0 &= \frac{\pi l \left( \varepsilon_{33}^T - d_{31,p}^2 \right)}{s_{11}^E}, \\ Z_1 &= j \sqrt{\frac{n\rho}{s_{33}^m} + \frac{(1 - n)\rho}{s_{11}^E}} A \tan(kl/2), \\ Z_2 &= \frac{\sqrt{\frac{n\rho}{s_{33}^m} + \frac{(1 - n)\rho}{s_{11}^E}}}{j\sin(kl)}, \end{aligned}$$

where  $\varphi_{\rm m}$  and  $\varphi_{\rm p}$  are the magnetomechanical and mechanoelectric coupling coefficients, respectively;  $\omega$  is the angular frequency;  $C_0$  is the clamped capacitance of piezoelectric long-thin tube;  $\beta$  is the internal stresses acting on the magnetostrictive continuous fibers; R is the mechanical damping resistance of the ME rod;  $Z_1$  and  $Z_2$  are the mechanical impedances of the ME rod; A and l are the cross-sectional area and length of the ME rod, respectively;  $A_{\rm m}$  and  $r_{\rm m}$  are the cross-sectional area and radius of the epoxy-bonded magnetostrictive continuous fiber composite rod, respectively;  $r_{\rm p}$ is the outer radius of the piezoelectric long-thin tube;  $d_{33,m}$ and  $d_{31,p}$  are the piezomagnetic and piezoelectric coefficients, respectively;  $s_{33}^H$  and  $s_{11}^E$  are the elastic compliance coefficients at constant magnetic field strength and at constant electric field strength, respectively;  $\varepsilon_{33}^T$  is the dielectric permittivity at constant stress;  $\rho$  is the average mass density of the ME rod;  $n = r_{\rm m}^2/r_{\rm p}^2$  is the cross-section area ratio of the ME rod; and k is the wave number. From Eq. (1), it is noted that  $\alpha_V$  of the ME rod depends on the material properties of the epoxy-bonded magnetostrictive continuous fiber composite rod and the piezoelectric long-thin tube (i.e.,  $d_{33,m}$ ,  $d_{31,p}$ ,  $s_{33}^H$ ,  $s_{11}^E$ ,  $\varepsilon_{33}^T$ , etc.) as well as the geometric parameters of the ME rod.

#### **III. MEASUREMENTS**

The ME effect in our ME rod was characterized at room temperature using an in-house automated ME measurement system.<sup>15</sup> The ME rod was placed between the pole gap of a water-cooled, U-shaped electromagnet (Mylten PEM-8005 K), and the electromagnet was energized by a dc current supply (Sorensen DHP200-15) to give a dc magnetic bias field ( $H_{\text{Bias}}$ ). An arbitrary waveform generator (Agilent 33210 A) connected to a constant-current supply amplifier (AE Techron 7796HF) was employed to drive a pair of Helmholtz coils for providing an ac magnetic drive field  $(H_3)$  over the prescribed frequency (f) range.  $H_{\text{Bias}}$  was monitored using a Hall probe connected to a Gaussmeter (F. W. Bell 7030), while  $H_3$  was measured using a pick-up coil connected to an integrating fluxmeter (Walker MF-10D). All qualities were gathered, together with the induced ME voltage  $(V_3)$ , using a data acquisition unit (Nation Instruments BNC-2110 and NI-PCI6132) under the control of a computer with a Labview program.

#### **IV. RESULTS AND DISCUSSION**

Figures 2(a) and 2(b) show the measured and calculated  $\alpha_V$  of the ME rod driven by an  $H_3$  of 1 Oe peak in the *f* range of 1–60 kHz under various  $H_{\text{Bias}}$  of 50–500 Oe, respectively. The  $\alpha_V$ -*f* curves in Fig. 2(b) were calculated based on Eq. (1),



FIG. 2. (a) Measured and (b) calculated  $\alpha_V$  of the ME rod driven by an  $H_3$  of 1 Oe peak in the *f* range of 1–60 kHz under various  $H_{\text{Bias}}$  of 50–500 Oe.

with the material properties given in Refs. 9 and 16, and the geometric parameters were described in Fig. 1. It is clear that the measured and calculated  $\alpha_V$  agree reasonably well with each other for all  $H_{\text{Bias}}$  levels with no remarkable frequency dispersion, except for the variations associated with the resonance frequency  $(f_r)$  range of 25–40 kHz. The observation indicates that the eddy-current losses are insignificant in the ME rod for f up to, at least, 60 kHz. Moreover, the largest  $\alpha_V$ response is found at a low  $H_{\text{Bias}}$  level of 120 Oe throughout the whole f range for both measurement and calculation. This is mainly due to the reduced effect of demagnetization fields in the high aspect-ratio magnetostrictive continuous fibers of 35:1. In particular, the largest resonance  $\alpha_V$  is detected to be 680 mV/Oe at  $f_r = 32$  kHz by measurement and 710 mV/Oe at  $f_{\rm r} = 33.8 \,\rm kHz$  by calculation. These resonances  $\alpha_V$ , which are caused by the fundamental axial mode resonance with half-wave vibrations along the length of the ME rod, are  $\sim 20$ times larger than their nonresonance  $\alpha_V$  of  $\sim 38$  mV/Oe for f up to 10 kHz. This resonance amplification of  $\alpha_V$  is even more significant when compared to the widely known ME laminated plates of Tb<sub>0.3</sub>Dy<sub>0.7</sub>Fe<sub>1.92</sub> alloy and 0.71Pb(Mg<sub>1/3</sub>Nb<sub>2/3</sub>)O<sub>3</sub>-0.29PbTiO<sub>3</sub> (PMN-PT) crystal of  $\sim$ 10 times, as a result of good magneto-mechano-electric coupling between the epoxybonded magnetostrictive continuous fiber composite rod and the piezoelectric long-thin tube in our ME rod enabled by the relatively true magnetostrictive 33-mode and piezoelectric 31mode of operation in the magnetostrictive continuous fibers and piezoelectric long-thin tube, respectively. When the  $H_{\text{Bias}}$ level is reduced to 50 Oe, the measured resonance  $\alpha_V$  undergoes a reduction of  $\sim 0.91$  times to 617 mV/Oe, while its corresponding  $f_r$  exhibits an up-shift of ~1 kHz-33 kHz. At an elevated  $H_{\text{Bias}}$  level of 500 Oe, the measured resonance  $\alpha_V$  has  $\sim 0.87$  times reduction to 590 mV/Oe, but its corresponding  $f_{\rm r}$ exhibits a smaller up-shift of ~0.2 kHz-33.2 kHz. The observations indicate that besides the presence of a high resonance ME effect in our ME rod, there also exists an interestingly high low-H<sub>Bias</sub>-controllable nonlinear resonance ME tuning effect in the ME rod originated from increasing, maximizing, and decreasing the motion of non-180° domain walls in the epoxy-bonded magnetostrictive continuous fiber composite rod for  $H_{\text{Bias}}$  below, at, and above 120 Oe, respectively.<sup>5,9</sup> This represents the operation of the ME rod in the magnetostrictive *negative*- $\Delta E$  region, at the *negative*- $\Delta E$ -*positive*- $\Delta E$  transition, and in the *positive*- $\Delta E$  region, respectively.<sup>5,9</sup>

Figure 3 plots the measured  $V_3$  over broad ranges of  $H_3$ of 0.1–10 Oe and  $H_{\text{Bias}}$  of 20–500 Oe at 1 kHz and  $f_r$ . It is seen that  $V_3$  varies essentially linearly with  $H_3$  for different  $H_{\text{Bias}}$  at both 1 kHz and  $f_r$ . The good linearity between  $V_3$  and  $H_3$  confirms the validity of the proposed working principle. The largest  $V_3$ – $H_3$  response is observed at  $H_{\text{Bias}} = 120$  Oe for all f, which is in good agreement with the observations in Fig. 2. From the slopes of plot,  $\alpha_V$  can be determined at different  $H_{\text{Bias}}$  and f. The inset of Fig. 3 shows the measured and calculated  $\alpha_V$  as a function of  $H_{\text{Bias}}$  at 1 kHz and  $f_r$ . The measured and calculated  $\alpha_V$  agree reasonably well with each other at both f. The relatively larger discrepancy at  $f_r$  than at 1 kHz is mainly due to the error in estimating the  $H_{\text{Bias}}$ -dependent mechanical damping resistance of the ME rod (R) in Eq. (1). Nonetheless, the measured  $\alpha_V$  is found to increase in



FIG. 3. Measured  $V_3$  over broad ranges of  $H_3$  of 0.1–10 Oe and  $H_{\text{Bias}}$  of 20–500 Oe at 1 kHz and  $f_r$ . The inset shows the measured (symbols) and calculated (lines)  $\alpha_V$  as a function of  $H_{\text{Bias}}$  at 1 kHz and  $f_r$ .

the initial region of  $H_{\text{Bias}}$  of 0–120 Oe, reaching the maximum value of 38 and 680 mV/Oe at 1 kHz and  $f_r$  under a low  $H_{\text{Bias}}$  of 120 Oe, respectively, and then decreasing with increasing  $H_{\text{Bias}}$ . The observation suggests the presence of an interestingly high low-H<sub>Bias</sub>-controllable nonlinear resonance ME tuning effect in the ME rod as described in Fig. 2. For comparison, the well-known L-T laminated plates generally exhibit a non-resonance  $\alpha_V$  of  $\sim 60 \text{ mV/Oe}$  at a higher  $H_{\text{Bias}}$  of 400–600 Oe,<sup>1</sup> while the T–T laminated plates demonstrate a quantitatively similar non-resonance  $\alpha_V$  of  $\sim 60 \text{ mV/Oe}$  at an even higher  $H_{\text{Bias}}$  of  $\sim 5000 \text{ Oe.}^2$ Therefore, the simultaneously high ME responses in our ME rod at the reasonably low  $H_{\text{Bias}}$  of 120 Oe can be attributed to the induction of the relatively true magnetostrictive 33response and piezoelectric 31-response by the magnetostrictive continuous fibers and piezoelectric long-thin tube, respectively, in addition to the reduction in the effect of demagnetization fields through the high aspect ratio magnetostrictive continuous fibers.<sup>10,11</sup>

#### **V. CONCLUSION**

We have found theoretically and experimentally the existence of a high non-resonance  $\alpha_V$  of 38 mV/Oe for *f* up to 10 kHz and an ultrahigh resonance  $\alpha_V$  of 650 mV/Oe at 33 kHz, both under a reasonably low  $H_{\text{Bias}}$  of 120 Oe in a 35 mm long, 8 mm diameter "rod-in-tube" ME heterostructure rod formed by an inner epoxy-bonded magnetostrictive continuous fiber composite rod and an outer piezoelectric long-thin tube. We have also attributed the observed axial ME responses at relatively low  $H_{\text{Bias}}$  to the deployment of the relatively true magnetostrictive 33-mode and piezoelectric 31-mode of operation in the ME rod, respectively, besides the reduction in the effect of demagnetization fields by the high aspect-ratio magnetostrictive continuous fibers. The single-dimensional design of the ME rod makes it high potential as an axial probe for passive sensing of axial magnetic fields.

#### ACKNOWLEDGMENTS

This work was supported by the Research Grants Council of the HKSAR Government (PolyU 5228/13 E) and The Hong Kong Polytechnic University (H-ZG76 and 1-ZV7P).

- <sup>1</sup>S. X. Dong, J. F. Li, and D. Viehland, IEEE Trans. Ultra. Ferroelectr. Freq. Control **50**, 1236 (2003).
- <sup>2</sup>C. W. Nan, L. Gang, and Y. H. Lin, Appl. Phys. Lett. 83, 4366 (2003).
- <sup>3</sup>Y. J. Wang, S. W. Or, H. L. W. Chan, X. Y. Zhao, and H. S. Luo, Appl. Phys. Lett. **92**, 123510 (2008).
- <sup>4</sup>C. M. Leung, S. W. Or, S. Y. Zhang, and S. L. Ho, J. Appl. Phys. **107**, 09D918 (2010).
- <sup>5</sup>Y. F. Duan, C. M. Leung, S. Y. Zhang, L. Zhang, and S. W. Or, J. Appl. Phys. **111**, 07C717 (2012).
- <sup>6</sup>C. M. Leung, S. W. Or, S. L. Ho, and K. Y. Lee, IEEE Sens. J. 14, 4305 (2014).
- <sup>7</sup>Y. Wang, D. Gray, D. Berry, M. Li, J. Gao, J. Li, and D. Viehland, J. Alloys Compd. **513**, 242 (2012).
- <sup>8</sup>S. W. Or, T. L. Li, and H. L. W. Chan, J. Appl. Phys. 97, 10M308 (2005).
- <sup>9</sup>C. Y. Lo, "Giant magnetostrictive composites for smart transducer and actuator applications," M. Phil. thesis, The Hong Kong Polytechnic University, 2007.
- <sup>10</sup>G. Engdahl, *Magnetostrictive Materials Handbook* (Academic Press, New York, 2000).
- <sup>11</sup>T. Ikeda, *Fundamentals of Piezoelectricity* (Oxford University Press, Oxford, 1990).
- <sup>12</sup>S. X. Dong and J. Y. Zhai, Chin. Sci. Bull. **53**, 2113 (2008).
- <sup>13</sup>F. Yang, Y. M. Wen, P. Li, M. Zheng, and L. X. Bian, J. Sens. Actuators, A 141, 129 (2008).
- <sup>14</sup>Z. Chen and Y. Su, J. Appl. Phys. **115**, 193906 (2014).
- <sup>15</sup>C. M. Leung, S. W. Or, F. F. Wang, S. L. Ho, and H. S. Luo, Rev. Sci. Instrum. 82, 013903 (2011).
- <sup>16</sup>Ferroperm Matdata, Ferroperm Piezoceramics A/S (2006).