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Transient-based frequency domain method for dead-end side branch detection in reservoir-pipeline-valve systems HF Duan<sup>1</sup>, M.ASCE; and PJ Lee<sup>2</sup>, M.ASCE <sup>1</sup> Assistant Professor, Department of Civil and Environmental Engineering, The Hong Kong Polytechnic University, Hung Hom, Kowloon, SAR Hong Kong, China. (Corresponding Author) Email: hf.duan@polyu.edu.hk. <sup>2</sup> Senior Lecturer, Department of Civil and Natural Resources Engineering, The University of Canterbury, Private Bag 4800 Christchurch, New Zealand. Email: pedro.lee@canterbury.ac.nz. **Abstract:** Unexpected, unknown or unused side pipe branches, termed dead-end side branches in this study, commonly exist in water pipe systems due to incorrect device installations or illegal connections. These connections are detrimental to the water quality as well as the operation and management of pipe networks. This paper investigates a transient-based frequency domain method for detecting side branches in pipe systems. The frequency response function for a pipe system with a single dead-end side branch is first derived by the transfer matrix method and the side branch was found to cause shifts in the system resonant frequencies. The nature of the resonant frequency shifts can be used to inversely determine the location and size of the side

21 is validated through numerical simulations and the results demonstrate the feasibility of this

method for detecting side branches. The accuracy for locating the side branch is higher than the

branch. A two-step Genetic algorithm based optimization is proposed in this study to efficiently

solve the derived analytical expression for the resonant frequency shifts. The developed method

accuracy in sizing the branches. The sensitivity of the method to the magnitude and bandwidth of

the transient wave signal is also discussed in the paper.

25 Keywords: Water pipe system; dead-end side branch; transient-based frequency domain method;

transients; transfer matrix; genetic algorithm

#### Introduction

Urban water pipe systems support the water supply, drainage, sewage as well as sea water systems of cities and are central to the social, economic and health sectors. The understanding and monitoring of such pipe systems are crucial to their management and efficient operation. Unexpected dead-end side branches commonly exist in pipe networks because of careless pipe installation, historical system modifications resulting in orphan pipes and illegal connections (Stephens 2008; and Meniconi et al. 2011a). At times, severe blockages and accidentally shut isolation valves can also lead to the formation of such side branches. These pipe branches are normally hidden underground and the fluid in these branches is static or under very small flowrate such that they are hard to detect using steady-state based methods. In this study, the flow is assumed to be static within these branches.

It is important to locate these side branches as the stagnant water within can compromise the water quality across the networks and the impedance of the branches can affect the hydraulics responses of the networks under normal operations (Carter et al. 1997). Pressure surges from normal pump and valve operations can enter the side branch and the reflection at the termination end can magnify the wave front, creating pressure fluctuations in the pipe networks that exceed the original transient design capacity (Karney and McInnis 1990; and Ferrante et al. 2009). Currently the methods used for side branch detection are physical inspection and steady hydraulics based analysis (Walski et al. 2003; Meniconi et al. 2011a). Such methods are limited to short range or long duration tests that require many months to complete the survey of the system.

Due to the magnification of the transient signal at a side branch, a logical approach is to use small amplitude (e.g., 1 to 10 m in pressure head), controlled fluid transients as a means of

detecting these hidden pipe sections (Meniconi et al. 2011b). An advantage of this approach is that fluid transient waves in elastic pipes travel at speeds exceeding 1000 m/s, allowing rapid detection procedures. Transient-based method has been applied to the detection of leaks and blockages in pipelines with good success (Liggett and Chen 1994; Brunone 1999; Vítkovský et al. 2000, 2007; Wang et al. 2002, 2005a; Beck et al. 2005; Covas et al. 2005; Lee et al. 2006, 2008; Mohapatra et al. 2006; Ferrante et al. 2007; AL-Khomairi 2008; Kim 2008; Sattar et al. 2008; Stephens 2008; Tuck et al. 2013; Duan et al. 2011, 2012a, 2012b, 2013, 2014; Meniconi et al. 2011a, 2011b, 2011c, 2013; and Kim et al. 2014). The principle of the transient based diagnostic method is: (1) to induce a small amplitude pressure signal (2) to measure the responses at accessible points in the system (such as hydrant points or maintenance locations) and (3) to compare the data with analytical models, in order to inversely determine the size and location of faults either in the time, or frequency domains (Colombo et al. 2009). It has been shown in the literature that transient fault diagnostics in the frequency domain can allow the fault to be isolated from complex system phenomena such as unsteady friction and pipe wall viscoelasticity as well as providing an increased tolerance to random noises (Lee et al. 2006, and Duan et al. 2011, 2012a). Therefore, in this study, the transient frequency responses of pipeline systems are used to detect side braches.

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Wang et al. (2005b) investigated the effect of lateral dead ends on the pipeline transients by forward analysis, where a linear analytical solution in the form of a Fourier series was derived for the shifting of harmonic frequencies by the lateral dead ends. Their results demonstrated the significant impact of lateral dead ends on the resonant frequencies, which is a function of the location, volume, and wave speed of the dead end. However, this study had a number of limitations. The linear assumptions of small amplitude transients and dead-ends were adopted in

their study and the lumped volume of the potential dead end was used in the derivation such that it is difficult to obtain the individual influence of the size and length of the dead ends on the pipe transients respectively. Moreover, their study did not produce a method for the detection of lateral dead ends in pipe systems.

A recent study by Meniconi et al. (2011a) has demonstrated the detection of inactive pipe side branches using time domain techniques. The location of the illegal side branch is given by the arrival time of the reflection from the branch, while the size of the side branch is determined by a simple frictionless relationship with the size of the reflection. While this technique is shown to work well, the number of reflections in a complex network can quickly become overwhelming and the size of the branch may be underestimated using the frictionless approximation of the system.

This paper presents a new frequency domain method for detecting side branches in pipe systems. The impact of a side branch on the transient response of a pipeline is first presented to illustrate how the response may be used for the detection and location of the branch. The impact of the side branch on the system frequency response is derived analytically using the transfer matrix method (Lee et al. 2006; Duan et al. 2011, 2012a), which is then used in a two-stage inverse calibration method to show the accurate detection and location of the side branch. The methodology is validated using numerical data simulated by the one-dimensional (1D) transient method of characteristics model.

# 1D water hammer models

The classical 1D water hammer model is adopted in this study for analytical and numerical investigations, mathematically expressed as (Chaudhry 1987; Wylie et al. 1993; and Ghidaoui et

97 al. 2005)

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$$\frac{gA}{a^2}\frac{\partial H}{\partial t} + \frac{\partial Q}{\partial x} = 0, \tag{1}$$

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$$\frac{\partial Q}{\partial t} + gA \frac{\partial H}{\partial x} + \frac{\pi D}{\rho} \tau_{w} = 0, \qquad (2)$$

where H = pressure head, Q = pipe discharge, A = pipe cross-sectional area, D = pipe diameter, a = 101 wavespeed, a = time, a = spatial coordinate along pipeline, a = gravitational acceleration, a = 102 fluid density, and a = pipe wall shear stress, and a = pipe a = 103 area averaged velocity; a = pipe friction factor, and a = steady and unsteady 104 components of pipe friction factor. Particularly, for laminar flow, a = 105 Reynolds number and a = kinematic viscosity of water; while for turbulent flow a = 106 approximately by the Colebrook equation (e.g., Chaudhry 1987; Wylie et al. 1993)

### Importance of side branches to the transient modeling and analysis

It has been widely evidenced in the literature that the side branches may have intense effect on the transient responses of pipelines systems (e.g., Chaudhary 1987; Wylie et al. 1993). To clearly demonstrate this effect in present study, the simple pipe system with a single side branch shown in Fig. 1 is used for this example. The main pipeline (pipe sections no. 1 and no. 2) is bounded by two constant head tanks ( $N_1$  and  $N_2$  in the figure) and the discharge in the pipeline is controlled by the valve located at the downstream tank ( $N_2$ ). A side branch is located at the junction  $N_3$  with zero discharge with the valve at terminal  $N_4$  fully shut under the steady state condition. The parameters of the system are as follows: (1) length (l): 400 m, 600 m, and 50 m; (2) diameter (D): 0.6m, 0.5m, and 0.2m; and (3) wave speed (a): 1000 m/s, 1200 m/s, and 1300

m/s, for pipe sections no. 1, no. 2, and no. 3, respectively. The steady state discharge ( $Q_0$ ) in the main pipeline (pipe sections 1 and 2) is assumed to be 0.2 m<sup>3</sup>/s.

### [Figure 1 will be here]

The transient is generated by the sudden closure of the downstream end valve at  $N_2$ . Based on the 1D numerical models in Eqs. (1) through (3), the transient pressure heads in the time domain at  $N_2$  for a pipeline system without the side branch (pipe no. 3) and with the side branch are plotted in Fig. 2(a) for comparison. Note that the grid size of the discretization is fixed at 1 m for all pipe sections and spatial interpolation is applied in the numerical simulation (Chaudhry 1987). The results show that the existence of the side branch could increase the magnitudes of the system transient responses in certain locations in the system (e.g., the location  $N_2$  shown here). These increases in the transient response may exceed the designed capacity of the pipe system.

Furthermore, the side branch has induced more wave oscillations by creating reflections in the transient responses, which can affect the dynamics of air-chambers, surge tanks and air-valves in the system (Duan et al. 2010a). The increased pressure oscillations are confirmed by the frequency domain results in Fig. 2(b), which demonstrates the shifting of resonant peaks and energy transfers between different frequency modes (amplitude changes). Note that current transient-based pipe leakage/blockage detection methods rely on the reflection and damping information of the measured/simulated transient responses (Duan et al. 2010b), therefore the understanding and characterization of the impact of side branches on transient reflections and damping is important for system design (modeling and analysis) as well as system management

(pipe defects detection). This paper aims to characterize the transient behavior of dead-end side branches in the pipe system and develop a transient-based frequency domain method for the side branch detection.

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# [Figure 2 will be here]

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### Analytical results of frequency responses for pipe systems with a side branch

To derive the frequency responses, the 1D mass and momentum equations in Eqns. (1) and (2) are first linearized and transformed into frequency domain equivalents to describe the behaviors of the transient pipe system (Chaudhry 1987; Duan et al. 2011, 2012a). The result of the downstream transient responses in terms of the upstream quantities can be obtained as (Lee et al. 2006),

$$\begin{cases} q \\ h \end{cases}^{n+1} = \begin{bmatrix} \cos(\mu l) & i \frac{1}{Y} \sin(\mu l) \\ iY \sin(\mu l) & \cos(\mu l) \end{bmatrix} \begin{bmatrix} q \\ h \end{bmatrix}^{n}, \tag{3}$$

where  $\mu = C_R \frac{\omega}{a}$  = wave number;  $Y = -C_R \frac{a}{gA}$  = characteristic impedance;  $C_R = \sqrt{1 - i \frac{gAR}{\omega}}$  =

friction coefficient;  $R = \frac{fQ}{gDA^2}$  = friction damping factor; q, h = discharge and pressure head in

the frequency domain; n, n+1 = upstream and downstream ends of pipe section;  $\omega$  = frequency;

and i = imaginary unit. The expression for friction damping factor R, derived in Vítkovský et al.

(2003) is used in this study. Note that this friction damping factor includes the effect of unsteady

friction.

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The transient frequency response function is a fundamental description of the system behavior and this function has been applied for the detection of leaks (Lee et al. 2006; Duan et al.

2011), blockages (Mohapatra et al. 2006; Lee et al. 2008; Sattar et al. 2008; Duan et al. 2012a, 2013, 2014), and the description of pipe wall visco-elasticity (Duan et al. 2012b). Using similar procedures in these studies, the transient frequency responses are derived for side branches in this paper. Using the system configuration of Fig. 1, the transient frequency response function at the valve  $N_2$  is obtained as (see Eq. 10 in the appendix),

$$Y_{2}Y_{3}\cos(\mu_{3}l_{3})\sin(\mu_{2}l_{2})\cos(\mu_{1}l_{1}) - Y_{2}Y_{1}\sin(\mu_{3}l_{3})\sin(\mu_{2}l_{2})\sin(\mu_{1}l_{1})$$

$$h^{N_{2}} = i\frac{+Y_{1}Y_{3}\cos(\mu_{3}l_{3})\cos(\mu_{2}l_{2})\sin(\mu_{1}l_{1})}{Y_{1}Y_{3}\cos(\mu_{3}l_{3})\sin(\mu_{2}l_{2})\sin(\mu_{1}l_{1}) + Y_{1}Y_{2}\sin(\mu_{3}l_{3})\cos(\mu_{2}l_{2})\sin(\mu_{1}l_{1})}.$$

$$-Y_{2}Y_{3}\cos(\mu_{3}l_{3})\cos(\mu_{2}l_{2})\cos(\mu_{1}l_{1})$$

$$(4)$$

Detailed derivations are presented in the appendix. The resonant condition for the single side branch system is given by Eq. (13) in the appendix,

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$$\begin{bmatrix} Y_{2}Y_{1}\sin(\mu_{3}l_{3})\cos(\mu_{2}l_{2})\sin(\mu_{l}l_{1}) \\ -Y_{3}Y_{2}\cos(\mu_{3}l_{3})\cos(\mu_{2}l_{2})\cos(\mu_{l}l_{1}) \\ +Y_{3}Y_{1}\cos(\mu_{3}l_{3})\sin(\mu_{2}l_{2})\sin(\mu_{l}l_{1}) \end{bmatrix} = 0,$$
 (5)

171 or alternatively,

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$$\begin{bmatrix}
(Y_{2}Y_{3} + Y_{1}Y_{2} + Y_{1}Y_{3})\cos(\mu_{1}l_{1} + \mu_{2}l_{2} + \mu_{3}l_{3}) \\
+ (Y_{2}Y_{3} - Y_{1}Y_{2} - Y_{1}Y_{3})\cos(\mu_{1}l_{1} - \mu_{2}l_{2} - \mu_{3}l_{3}) \\
+ (Y_{2}Y_{3} + Y_{1}Y_{3} - Y_{1}Y_{2})\cos(\mu_{1}l_{1} + \mu_{2}l_{2} - \mu_{3}l_{3}) \\
+ (Y_{2}Y_{3} - Y_{1}Y_{3} + Y_{1}Y_{2})\cos(\mu_{1}l_{1} - \mu_{2}l_{2} + \mu_{3}l_{3})
\end{bmatrix} = 0.$$
(6)

Note that in this study only the case of a single dead-end side branch is investigated and tested, and similar derivation process can be applied for the situations of multiple side branches in complex pipe system.

### Simplification and analysis of analytical results

To use Eq. (5) for side branch detection, the resonant frequencies of the system have to be obtained by experimental measurement or numerical simulations, and then inversely fitted to Eq.

(5). To obtain the solution, an inverse method such as GA-based optimization used in Duan et al. (2012a) is needed. However, Eq. (5) is too implicit to show the direct relationship between the resonant frequencies and properties of the side branch. It is useful to further simplify the equation to obtain a clear illustration of this dependency, in a similar analytical procedure used for extended blockages in Duan et al. (2013, 2014). Under the assumption that the side branch is small relative to the main pipeline and using a first-order approximation, the simplified expression for the resonant frequency shifts by the side branch can be obtained as,

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$$\frac{\Delta \omega_{rf}}{\omega_{th0}} = -\frac{2\varepsilon}{\pi} \frac{\sin\left[\lambda_3 \omega_{rf0}\right] \cos^2\left[\lambda_2 \omega_{rf0}\right]}{\cos\left[\lambda_3 \omega_{rf0}\right]},\tag{7}$$

where  $\Delta \omega_{rf} = \Delta \omega_{rf}(k) =$  resonant frequency shift;  $\omega_{rf0} = \omega_{rf0}(k) =$ system resonant frequency with subscript 0 indicating the system without a side branch; k = resonant peak number;  $\omega_{th0} = \frac{\pi}{2\lambda_0} =$  theoretical (first resonance) frequency of the reservoir-pipe-valve system, that is,  $\omega_{th0} = \omega_{rf0}(k=1)$ ;  $\lambda = \mu l/\omega =$  wave propagation operator; and  $\varepsilon = Y_0/Y_3 =$  the ratio of characteristic impedances of main pipeline and side branch, and  $\varepsilon <<1$  is assumed in the derivation of Eq. (7). Note that for simplification the resonant frequency peak number has been neglected in Eq. (7), but this equation is valid for all resonant peaks in the frequency domain.

Eq. (7) clearly shows the dependence of resonant frequency shifts on the side branch properties (i.e., branch size, length and location derived from the parameters  $\varepsilon$ ,  $\lambda_3$ , and  $\lambda_2$  respectively). Meanwhile, compared to the full form of Eq. (5), the simplified Eq. (7) is simpler to solve using the inverse fitting process. The solution of Eq. (7), however, could only provide a reasonable initial guess for solving the full form Eq. (5) due to the assumption of small side branch and linearity.

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#### **Application procedure of side branch detection**

For accuracy, Eq. (5) has to be used to identify and obtain the properties of the side branch in the system. The difficulty and complexity for solving Eq. (5) is very similar to that for the equation of extended blockage detection in Duan et al. (2012a). Therefore, the genetic algorithm (GA)-based optimization procedure adopted in Duan et al. (2012a) is also used herein for solving Eq. (5). The objective function for the optimization is the absolute value of the left hand side of Eq. (5), which will be minimized to close to zero as much as possible. To avoid singularity problems during the optimization process, the variables are chosen as  $D_3$ ,  $l_2$ ,  $l_3$ , and  $a_2$ , instead of  $Y_3$ ,  $\lambda_3$ , and  $\lambda_2$  in Eq. (5). The rewritten form of Eq. (5) is given by Eq. (14) in the appendix.

The efficiency and accuracy of the GA-based optimization is highly dependent on the initial values and the search range. In this study, a two-step method is used, where Eq. (7) is first solved to provide the initial values and searching ranges of all parameters for the GA-based optimization of Eq. (5). It is assumed that the side branch is smaller than the main pipeline and the search range of the size and length of the side branch range between zero to the values of the main pipeline (that is  $[0, \max(D_1, D_2)]$  and  $[0, l_0]$  with  $l_0 = l_1 + l_2$  as in Fig. 1).

### [Figure 3 will be here]

The GA scheme developed in Vítkovský et al. (2000) is used for the optimization. The procedure of the GA-based optimization is shown in Fig. 3. The optimization process ends when either one of the two conditions are met (1) the relative error of all the results for any two continuous iterations is smaller than 1%; (2) the maximum iteration number, M = 5000, is

reached.

# **Numerical applications**

Numerical tests are used to examine the two-step optimization procedure for detecting side branches. The test settings are listed in Table 1, covering a wide range of system scales and flow regimes. Ten numerical tests were conducted in this study to investigate the influence of different side branch configurations and hydraulic conditions on the accuracy of the proposed method. The total length of pipeline ( $l_0$ ) without a side branch is kept at 1000 m, while the length, size and wavespeed of each pipe section divided by the side branch change with different tests. Smooth pipe friction factors are assumed for all pipe sections. Numerical transient simulations are conducted by the 1D water hammer equations coupled with unsteady friction models. The quasi-steady friction effect is represented by the Darcy-Weisbach formula; and the unsteady friction is calculated by Zielke's model for laminar flows (Zielke, 1968) and modeled by Vardy-Brown's model for turbulent flows (Vardy and Brown, 1995).

# [**Table 1** will be here]

Transients in the system are caused by the operation of a side-discharge valve at point  $N_2$  in Fig. 1. The discharge perturbation  $(Q_V)$  consists of a fast close-open-close operation of the valve as shown in Fig. 4, in order to generate a sharp wave signal with a large bandwidth in the system (Duan et al. 2010b, Lee et al. 2014). The influence of different input signals is discussed later in the paper. In Fig. 4, the valve discharge perturbation on the vertical coordinate is normalized by the initial steady state flowrate  $(Q_0)$ , and the axial time coordinate is normalized

by the wave period for the main pipeline  $(4l_0/a_0)$ . The pressure head at point  $N_2$  in Fig. 1 is collected for analysis.

### [**Figure 4** will be here]

The simulated pressure head for test case no. 1 is plotted in Fig. 5 (a), where the vertical coordinate is the transient pressure head ( $\Delta H_t$ ) normalized by the *Joukowsky* head rise (i.e.,  $\Delta H_0$ ) =  $a\Delta V_{v0}/g$  with  $\Delta V_{v0}$  = the flow velocity change of the initial side discharge valve operation) and the axial coordinate is dimensionless time in terms of the wave period of the main pipeline. The frequency response function (Lee et al. 2006) is shown in Fig. 5 (b) with the vertical coordinate normalized by fundamental harmonic response of the main pipeline and the axial coordinate normalized by the fundamental frequency ( $a/4l_0$ ) of main pipeline. For comparison, the results for a system without a side branch are also plotted in Fig. 5. The results demonstrate clearly the additional wave reflections in the time domain in Fig. 5(a) and the frequency shifts of resonant peaks in the frequency domain in Fig. 5(b).

### [**Figure 5** will be here]

### Results analysis and discussion

The results of Fig. 5 for test no. 1 are first used to illustrate the efficiency of two different optimization approaches for side branch detection: (1) inverse fitting the results directly to Eq. (5); and (2) using the two-step procedure shown in Fig. 3 to first inversely fit the results to the simplified Eq. (7) to provide a starting guess and refine the search space for a second inverse

optimization with Eq. (5). The first 15 resonant peaks in Fig. 5(b) are extracted for analysis. The operations of the two optimization schemes are shown in Fig. 6, which clearly indicates a faster convergence for the two-step optimization scheme. The CPU computation time for the two-step method is less than 1/3 of that for the direct GA-based optimization scheme for the same accuracy. Particularly in this test case here, the computation time required by two-step method is about 3 minutes, while that by direct optimization method is over 10 minutes.

# [Figure 6 will be here]

The result of the optimization for test no. 1 is plotted in Fig. 7, which reveals the good agreement between the results from the two-step optimization and the numerical data. Meanwhile, the result of Fig. 7 also shows clearly the improvement of the proposed two-step optimization compared to the single step optimization for Eq. (7) only. The two-step GA-based optimization scheme in Fig. 3 is used for side branch detection in the following analysis.

The side branch characteristics determined for all the ten tests in Table 1 are shown in Table 2. The relative error of the prediction ( $\gamma$  in Table 2) for each parameter is given by,

$$\gamma_P(\%) = \frac{\left| P_r - P_p \right|}{P_r} \times 100, \qquad (8)$$

where P is the parameter being considered and subscripts "r" and "p" refer to the real and predicted parameter values respectively. It is also necessary to point out that the predicted values

during the GA-based optimization process are real numbers with more than 3 decimal places, but only the rounded integer values are shown here in Table 2.

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### [Table 2 will be here]

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The results in Table 2 demonstrate that the proposed approach can be applied for side branch detection, as the prediction errors are within 20% for all parameters in the numerical tests in this study. The maximum relative errors for  $D_3$ ,  $l_3$ ,  $a_3$ , and  $l_2$  are 14.7%, 14.0%, 4.6% and 3.7% respectively (numbers in bold in Table 2), which implies that the proposed method is more accurate for locating  $(l_2)$  than sizing (size  $D_3$  and length  $l_3$ ) the side branch. In addition, Table 2 indicates that the accuracy of the method is decreasing with the scales of the side branch under the same other conditions during the testing (e.g., tests no. 3, 5 & 6; and tests no. 7 & 8), with the worst prediction for test no. 3, which has the largest side branch amongst all the cases considered. This large error is due to the violation of the small side branch approximation used in the analytical derivation (e.g.,  $\varepsilon \le 10\%$ ,  $l_3/l_0 < 10\%$ ). Furthermore, the results in Table 2 also reveal that detection accuracy under laminar flow (test no. 1 in Table 1) is higher than that for the turbulent system (tests no. 2 through 10). For all the turbulent flow cases (tests no. 2 through 10), the accuracy of the detection results is by and large decreasing with initial turbulent conditions (i.e.,  $Re_0$ ). This variation trend of the detection accuracy is mainly because of the linearization approximation of the quasi-steady friction term (relatively small transient flow to the steady base flow) used in the analytical derivation of Eq. (5) (Duan et al. 2013; and Meniconi et al. 2013). It is necessary to note that this conclusion will require further systematical analysis and experimental validations under a wide range of system configurations and conditions in the

future work.

# **Influence of input transient signals**

For the practical application of the proposed method, it is also necessary to consider the impact of different transient signals on the detection accuracy. Real valves will generate signals of different properties and the sensitivity of the method to the shape of the signal is an important consideration. Test no. 1 in Table 1 is used for this analysis and three different input signals shown in Fig. 8 are used in conjunction with the proposed side branch detection method. In Fig. 8, signal #1 represents the "fast and sharp" (large frequency bandwidth) signal; while signal #2 represent a larger magnitude input signal but of the same shape and signal #3 is a signal of the same magnitude as the original signal but of a much smoother form (small frequency bandwidth).

### [Figure 8 will be here]

For a fair comparison, only the first 7 resonant frequency peaks are used for side branch detection as signal #3 do not have detectable resonant peaks beyond this point. The detection results for the three different input signals are shown in Table 3, with the detection accuracies for signals #2 and #3 being much lower than that for signal #1. This result implies that a small amplitude but sharp signal produces the best accuracy from the method. This result is consistent with previous studies in Duan et al. (2010b) and Lee et al. (2014). The larger errors with the other signals are caused by the violation of the linearization approximation used in the analytical derivation of Eq. (5) in the case for signal #2 (Lee and Vítkovský 2010) and the inaccurate representation of high frequency resonant peaks (Lee et al. 2014) in the case for signal #3.

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### [Table 3 will be here]

Furthermore, the results in Table 3 show that the influence of the input signal is more profound for the prediction of the side branch size and length than for the side branch location and wave speed. This result is consistent with the results and findings in previous studies for the extended blockage detection (Duan et al. 2012a, 2013; and Meniconi et al. 2013), and is due to the insensitivity of the resonant frequency to the size and length of the faults (e.g., side branch and extended blockage) (Duan et al. 2014).

### **Summary and conclusions**

This paper investigates the detection of dead-end side branch in pipe systems by using transient-based frequency domain method. The transient responses for systems with and without side branches are compared in this study to illustrate the need to identify such side branches in the pipe system. Based on the transfer matrix method, the analytical expression of the frequency response function for a single dead-end side branch in the pipe system is derived and used for the detection of dead-end side branches in this study. The derived implicit equation is further simplified based on the assumptions of small side branches and linearity, in order to give a clear relationship between the resonant frequency shifts and the side branch information. This simplified expression improves the initial guesses and search space for the genetic algorithm (GA) based optimization. For this purpose, a two-step GA-based optimization scheme is proposed in this study to provide a more efficient solution.

The developed dead-end side branch detection scheme has been validated by numerical

simulations in this study. The results show that the developed method coupled with the optimization scheme can provide an accurate detection of the side branches in pipe system. It is also found that the proposed two-step optimization scheme is more efficient than the common direct GA-based optimization method. At the same time, the results indicate that the developed method in this study can provide a more accurate prediction for the location and wavespeed of the dead-end side branches than that for the size and length of the branches. It is also noted that further experimental tests will be needed to test the robustness and accuracy of the proposed method for the detection of multiple side branches, and more investigations will be required to examine the influences of other practical complexities and factors such as fluid-structure interaction (FSI) and system uncertainties to the proposed method.

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- Appendix: analytical derivation of transient frequency responses for pipe systems with
- 379 dean-end side branches
- Based on the method of transfer matrix in Eq. (3), the frequency responses at downstream end
- " $N_2$ " in the system of Fig. 1 can be expressed as,

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$$\begin{cases} q \\ h \end{cases}^{N_2} = \begin{bmatrix} \cos(\mu_2 l_2) & i \frac{1}{Y_2} \sin(\mu_2 l_2) \\ i Y_2 \sin(\mu_2 l_2) & \cos(\mu_2 l_2) \end{bmatrix} \begin{bmatrix} 1 & Z_{DESB} \\ 0 & 1 \end{bmatrix} \begin{bmatrix} \cos(\mu_1 l_1) & i \frac{1}{Y_1} \sin(\mu_1 l_1) \\ i Y_1 \sin(\mu_1 l_1) & \cos(\mu_1 l_1) \end{bmatrix} \begin{cases} q \\ h \end{cases}^{N_1},$$
 (9)

where  $Z_{DESB}$  = the hydraulic impendences of the dead-end side branch in the system of Fig. 1.

384 After mathematical manipulation and rearrangement, the perturbations of transient pressure head

at downstream B can be obtained as,

386 
$$h^{N_2} = -\frac{iY_2 \sin(\mu_2 l_2) \cos(\mu_1 l_1) - Z_{DESB} Y_2 Y_1 \sin(\mu_2 l_2) \sin(\mu_1 l_1) + iY_1 \cos(\mu_2 l_2) \sin(\mu_1 l_1)}{\cos(\mu_2 l_2) \cos(\mu_1 l_1) - \frac{Y_1}{Y_2} \sin(\mu_2 l_2) \sin(\mu_1 l_1) + iY_1 Z_{DESB} \cos(\mu_2 l_2) \sin(\mu_1 l_1)}.$$
(10)

387 Therefore, the resonant condition for the branched pipe system is,

388 
$$\cos(\mu_2 l_2) \cos(\mu_1 l_1) - \frac{Y_1}{Y_2} \sin(\mu_2 l_2) \sin(\mu_1 l_1) + i Y_1 Z_{DESB} \cos(\mu_2 l_2) \sin(\mu_1 l_1) = 0. \tag{11}$$

Considering the dead-end boundary at the side branch, the impedance can be given by

390 (Chaudhry 1987),

$$Z_{DESB} = \frac{i \sin\left(\mu_3 l_3\right)}{Y_3 \cos\left(\mu_3 l_3\right)}.$$
 (12)

392 Substituting Eq. (12) into Eq. (11) gives,

$$\begin{bmatrix}
Y_{2}Y_{1}\sin(\mu_{3}l_{3})\cos(\mu_{2}l_{2})\sin(\mu_{1}l_{1}) \\
-Y_{3}Y_{2}\cos(\mu_{3}l_{3})\cos(\mu_{2}l_{2})\cos(\mu_{1}l_{1}) \\
+Y_{3}Y_{1}\cos(\mu_{3}l_{3})\sin(\mu_{2}l_{2})\sin(\mu_{1}l_{1})
\end{bmatrix} = 0,$$
(13)

or in terms of pipe parameters (D, l, and a),

$$\begin{bmatrix}
C_{R2}C_{R1}A_{3}a_{2}a_{1}\sin\left(C_{R3}\frac{l_{3}}{a_{3}}\omega_{rfb}\right)\cos\left(C_{R2}\frac{l_{2}}{a_{2}}\omega_{rfb}\right)\sin\left(C_{R1}\frac{l_{1}}{a_{1}}\omega_{rfb}\right) \\
-C_{R3}C_{R2}a_{3}a_{2}A_{1}\cos\left(C_{R3}\frac{l_{3}}{a_{3}}\omega_{rfb}\right)\cos\left(C_{R2}\frac{l_{2}}{a_{2}}\omega_{rfb}\right)\cos\left(C_{R1}\frac{l_{1}}{a_{1}}\omega_{rfb}\right) \\
+C_{R3}C_{R1}a_{3}A_{2}a_{1}\cos\left(C_{R3}\frac{l_{3}}{a_{3}}\omega_{rfb}\right)\sin\left(C_{R2}\frac{l_{2}}{a_{2}}\omega_{rfb}\right)\sin\left(C_{R1}\frac{l_{1}}{a_{1}}\omega_{rfb}\right) \\
+C_{R3}C_{R1}a_{3}A_{2}a_{1}\cos\left(C_{R3}\frac{l_{3}}{a_{3}}\omega_{rfb}\right)\sin\left(C_{R2}\frac{l_{2}}{a_{2}}\omega_{rfb}\right)\sin\left(C_{R1}\frac{l_{1}}{a_{1}}\omega_{rfb}\right)
\end{bmatrix} = 0.$$
(14)

Eq. (13) and Eq. (14) can be further simplified by neglecting the friction effect, i.e.,  $C_R$ 

=1. Therefore,

$$Y = -\frac{a}{gA}; \ \mu = \frac{\omega}{a}. \tag{15}$$

399 Furthermore, considering the single main pipeline case, gives,

400 
$$A_1 = A_2 = A_3$$
;  $a_1 = a_2 = a_3$ ; and  $Y_1 = Y_2 = Y_3$ . (16)

401 Under the conditions of Eq. (15) and Eq. (16), Eq. (13) can be simplified as,

$$\begin{bmatrix}
(2+\varepsilon)\cos\left[(\lambda_{1}+\lambda_{2}+\lambda_{3})\omega_{rfb}\right] \\
-\varepsilon\cos\left[(\lambda_{1}-\lambda_{2}-\lambda_{3})\omega_{rfb}\right] \\
+(2-\varepsilon)\cos\left[(\lambda_{1}+\lambda_{2}-\lambda_{3})\omega_{rfb}\right] \\
+\varepsilon\cos\left[(\lambda_{1}-\lambda_{2}+\lambda_{3})\omega_{rfb}\right]
\end{bmatrix} = 0,$$
(17)

403 where 
$$\varepsilon = \frac{Y_0}{Y_3} = \frac{a_0}{a_3} \frac{A_3}{A_0}$$
; and  $\lambda = \frac{l}{a}$ .

It is easy to obtain the resonance condition for original intact case of single main pipeline

405 system,  $\omega_{rf0}$ , given by

$$\cos\left|\left(\lambda_1 + \lambda_2\right)\omega_{rf0}\right| = 0. \tag{18}$$

407 That is,

$$\cos\left[\left(\frac{l_1 + l_2}{a_0}\right)\omega_{rf0}\right] = \cos\left[\frac{l_0}{a_0}\omega_{rf0}\right] = 0.$$
 (19)

By setting  $\omega_{r,fb} = \omega_{r,f0} + \Delta\omega_{r,f}$  or  $\Delta\omega_{r,f} = \omega_{r,fb} - \omega_{r,f0}$  where  $\Delta\omega_{r,f} = \Delta\omega_{r,f}(k)$  is the shift (difference) of

resonant frequency between the systems with and without side branches for a particular resonant

peak (k), the co-sinusoidal functions in Eq. (17) can be expanded about  $\omega_{rf0}$  with a first-order

approximation as (Duan et al. 2013),

$$\cos\left[\alpha\omega_{x_0}\right] = \cos\left[\alpha\omega_{x_0}\right] - \alpha\sin\left[\alpha\omega_{x_0}\right] \Delta\omega_{x_0} + O\left[\left(\Delta\omega_{x_0}\right)^2\right] + \cdots, \tag{20}$$

where  $\alpha$  represents the coefficients of the resonant frequency terms in Eq. (17). Note that the

resonant peak number has been ignored in Eq. (20) for convenience. Combining Eq. (17) through

416 Eq. (20), gives,

$$\Delta \omega_{rf} = \frac{C_u}{C_d}, \tag{21}$$

418 where,

419
$$C_{u} = \begin{bmatrix} (2+\varepsilon)\sin\left[\lambda_{3}\omega_{rf0}\right] \\ \varepsilon\sin\left[(2\lambda_{2}+\lambda_{3})\omega_{rf0}\right] \\ -(2-\varepsilon)\sin\left[\lambda_{3}\omega_{rf0}\right] \\ -\varepsilon\sin\left[(2\lambda_{2}-\lambda_{3})\omega_{rf0}\right] \end{bmatrix}; \quad C_{d} = \begin{bmatrix} -(2+\varepsilon)(\lambda_{1}+\lambda_{2}+\lambda_{3})\cos\left[\lambda_{3}\omega_{rf0}\right] \\ +\varepsilon(\lambda_{1}-\lambda_{2}-\lambda_{3})\cos\left[(2\lambda_{2}+\lambda_{3})\omega_{rf0}\right] \\ -(2-\varepsilon)(\lambda_{1}+\lambda_{2}-\lambda_{3})\cos\left[\lambda_{3}\omega_{rf0}\right] \\ -\varepsilon(\lambda_{1}-\lambda_{2}+\lambda_{3})\cos\left[(2\lambda_{2}-\lambda_{3})\omega_{rf0}\right] \end{bmatrix}. \tag{22}$$

420 Particularly, for relative small side branch, it has,

421 
$$\varepsilon = \frac{Y_0}{Y_3} = \frac{a_0}{a_3} \frac{A_3}{A_0} << 1$$
, and  $\frac{\lambda_3}{\lambda_1 + \lambda_2} << 1$ . (23)

422 As a result, Eq. (21) becomes,

$$\Delta\omega_{rf} = -\frac{\varepsilon}{\lambda_0} \frac{\sin\left[\lambda_3\omega_{rf0}\right] \cos^2\left[\lambda_2\omega_{rf0}\right]}{\cos\left[\lambda_3\omega_{rf0}\right]},\tag{24}$$

424 or in dimensionless form,

425 
$$\frac{\Delta \omega_{rf}}{\omega_{th0}} = -\frac{2\varepsilon}{\pi} \frac{\sin\left[\lambda_3 \omega_{rf0}\right] \cos^2\left[\lambda_2 \omega_{rf0}\right]}{\cos\left[\lambda_3 \omega_{rf0}\right]},$$
 (25)

where  $\omega_{th0}$  = theoretical (first resonance) frequency of the intact (DESB-free) pipeline system

427 and  $\omega_{th0} = 2\pi \frac{a_0}{4l_0} = \frac{\pi}{2\lambda_0}$  in this study.

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556		Listing of figure captions
557	Fig. 1:	Illustrative pipe system with single dead-end side branch (pipe section no. 3)
558	Fig. 2:	Comparison of transient responses for system with and without dead-end side branch:
559		(a) time domain results; (b) frequency domain results
560	Fig. 3:	Flowchart of GA-based optimization for dead-end side branch detection
561	Fig. 4:	Flow perturbations of side-discharge valve at downstream end
562	Fig. 5:	Simulated pressure head traces at downstream valve in the pipe system with and
563		without side branch: (a) time domain results; (b) frequency domain results
564	Fig. 6:	GA-based optimization process of DESB detection for test no. 1
565	Fig. 7:	Resonant frequency shifts by optimization process and 1D model simulation
566	Fig. 8:	Input signals for different side-dischange valve operations
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579		Listing of table captions
580	Table 1:	System settings for numerical tests
581	Table 2:	Side branch detection results for all numerical tests
582	Table 3:	Detection results for test no. 1 by different input signals
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Table 1. System settings for numerical tests

	S	ection no	. 1	S	ection no	. 2	Section	on no. 3 (1	Initial	
Test no.	l <sub>1</sub> (m)	D <sub>1</sub> (mm)	a <sub>1</sub> (m/s)	l <sub>2</sub> (m)	D <sub>2</sub> (mm)	a <sub>2</sub> (m/s)	l <sub>3</sub> (m)	D <sub>3</sub> (mm)	<i>a</i> <sub>3</sub> (m/s)	Reynolds number ( $Re_0$ ) (×10 <sup>3</sup> )
1	350	500	1000	650	500	1000	50	100	1200	1
2	350	500	1000	650	500	1000	50	100	1200	5
3	350	500	1000	650	500	1000	50	100	1200	10
4	300	500	1000	700	400	1100	50	200	1200	5
5	350	500	1000	650	500	1100	50	200	1200	10
6	350	500	1000	650	400	1100	150	300	1200	10
7	550	500	1000	450	400	1100	50	100	1200	30
8	550	500	1000	450	400	1100	50	50	1200	30
9	700	500	1000	300	400	1000	50	50	1300	100
10	550	500	1000	450	400	1100	50	50	1300	100

Table 2. Side branch detection results for all numerical tests

Test		$D_3$ (mn	n)	$l_3$ (m)				<i>a</i> <sub>3</sub> (m/s)		$l_2$ (m)		
no.	$D_{3, r}$	$D_{3,p}$	γ(%)	$l_{3, r}$	$l_{3,p}$	γ(%)	$a_{3,r}$	$a_{3,p}$	γ(%)	$l_{2, r}$	$l_{2,p}$	γ(%)
1	100	103	3.0	50	51	2.0	1200	1210	0.8	650	650	0.0
2	100	102	2.0	50	50	0.0	1200	1199	0.1	650	661	1.7
3	100	100	0.0	50	51	2.0	1200	1193	0.6	650	653	0.5
4	200	211	5.5	50	47	6.0	1200	1219	1.6	700	707	1.0
5	200	205	2.5	50	54	8.0	1200	1189	0.9	650	650	0.0
6	300	256	14.7	150	135	10.0	1200	1243	3.6	650	674	3.7
7	100	113	13.0	50	57	14.0	1200	1145	4.6	450	445	1.1
8	100	106	6.0	150	139	7.3	1200	1235	2.9	450	456	1.3
9	50	52	4.0	50	56	12.0	1300	1291	0.7	300	308	2.7
10	50	54	8.0	50	54	8.0	1300	1281	1.5	450	443	1.6

Table 3. Detection results for test no. 1 by different input signals

									1 0			
Signal	$D_3$ (mm)			$l_3$ (m)			<i>a</i> <sub>3</sub> (m/s)			$l_2(m)$		
no.	$D_{3,r}$	$D_{3,p}$	γ(%)	$l_{3, r}$	$l_{3,p}$	γ(%)	$a_{3,r}$	$a_{3,p}$	γ(%)	$l_{2, r}$	$l_{2,p}$	γ(%)
#1	100	103	2.9	50	51	2.0	1200	1210	0.8	650	650	0.0
#2	100	92	8.7	50	54	7.4	1200	1211	0.9	650	650	0.0
#3	100	115	13.0	50	47	6.4	1200	1206	0.5	650	653	0.5

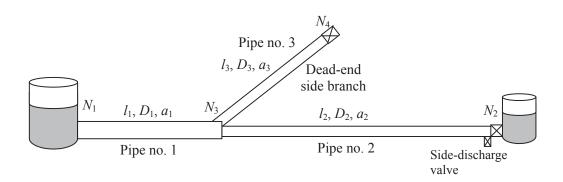
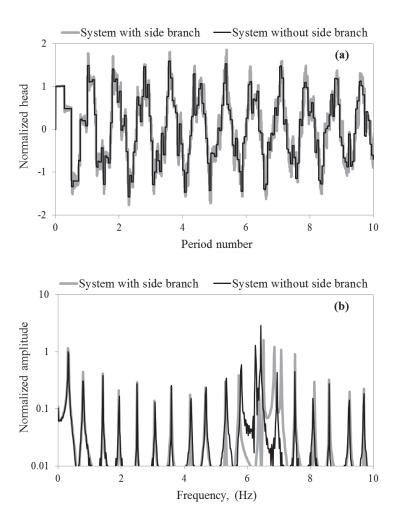


Fig. 1. Illustrative pipe system with single dead-end side branch (pipe section no. 3)



**Fig. 2.** Comparison of transient responses for system with and without dead-end side branch: (a) time domain results; (b) frequency domain results

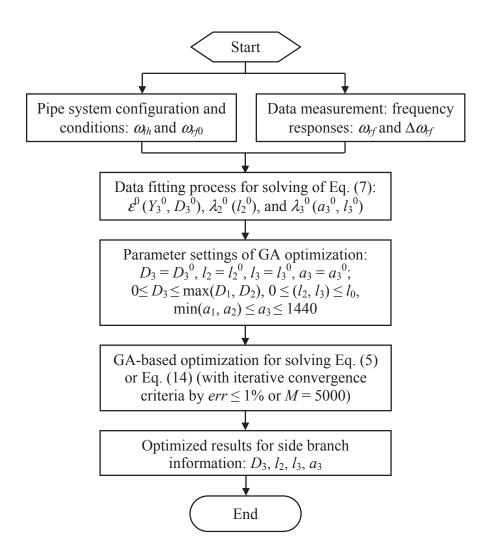


Fig. 3. Flowchart of GA-based optimization for dead-end side branch detection

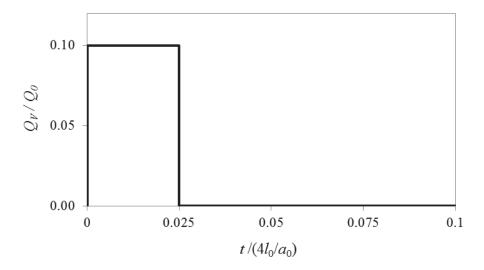
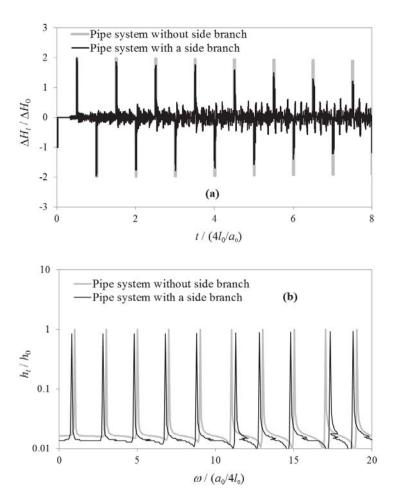


Fig. 4. Flow perturbations of side-discharge valve at downstream end



**Fig. 5.** Simulated pressure head traces at downstream valve in the pipe system with and without side branch: (a) time domain results; (b) frequency domain results

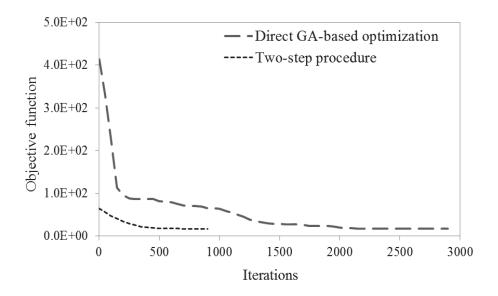


Fig. 6. GA-based optimization process of DESB detection for test no. 1

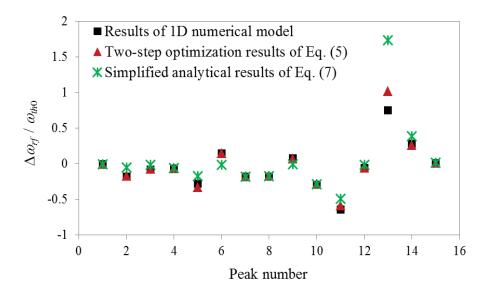


Fig. 7. Resonant frequency shifts by optimization process and 1D model simulation

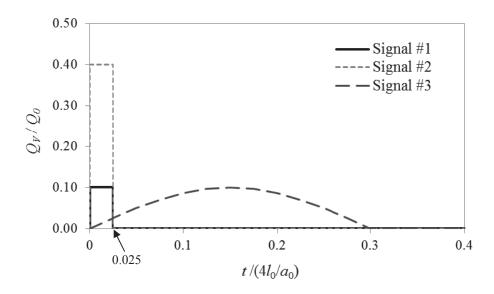


Fig. 8. Input signals for different side-dischange valve operations