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Disordered Layers on WO₃ Nanoparticles Enable Photochemical Generation of Hydrogen from Water

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Luyang Wang,^a Chui-Shan Tsang,^a Wei Liu,^a Xiandi Zhang,^a Kan Zhang,^b Enna Ha,^a Wai-Ming Kwok,^a Jong Hyeok Park,^b Lawrence Yoon Suk Lee,^{*a} and Kwok-Yin Wong^{*a}

Tailored defects on semiconductor surface can provide active catalytic sites and effectively tune the electronic structure for suitable optical properties. Herein, we report that surface modification of WO_3 with disordered layer enables the photochemical hydrogen production from water. A simple room temperature solution process with lithium-ethylenediamine (Li-EDA) alters the surface of WO_3 with localized defects that form a thin disordered layer. Both structural and optical characterizations reveal that such disordered layer induces an upshift in Fermi level and the elevation of the conduction band of WO_3 above the hydrogen reduction potential. Using alkaline sacrificial agent, Li-EDA treated WO_3 shows a co-catalyst-free photochemical hydrogen evolution rate of 94.2 μ mol g⁻¹ h⁻¹ under simulated sun light. To the best of our knowledge, this is the first example of WO_3 as a direct photocatalyst for hydrogen generation from water via simple surface defect engineering.

Introduction

Solar energy-driven hydrogen generation from water can offer a clean and sustainable solution to address the growing energy demand with reduced fossil fuel dependence and associated environmental issues. ¹⁻³ In recent years, much effort has been made to develop effective photocatalysts based on abundant and economic elements. ⁴⁻⁷ However, their overall photocatalytic efficiencies are insufficient to replace those made of noble metals. Tuning the light absorption, charge kinetics, and surface activity of catalysts can elevate the photocatalytic performance, but it is still challenging to achieve the activity of noble metal-based ones.

Semiconductor metal oxides, with advantages of natural abundance, chemical stability, and low cost, have been widely used as photocatalysts and electrode materials for energy conversion and storage devices.⁸ Both electronic and optical properties of metal oxides can be regulated by the type and population of defect introduced. The creation of intrinsic defects, such as oxygen vacancies and doping of low valence metal centers or non-metal atoms, are the most commonly used defect engineering methods for tailoring the photocatalytic activity.⁹⁻¹² In particular, oxygen vacancies are known to serve as shallow donors for metal oxides, such as TiO₂, ¹³ WO₃, ¹⁴ and SnO₂. ¹⁵ It is generally accepted that the increased carrier density improves the electrical conductivity and facilitates the charge transfer at the interface between metal oxide and substrate.

The introduced defects can be extended to form an internal boundary⁷ or even disordered layer,¹⁰ which are relatively less studied than atomic point defects. Surface disordered layer arises from the large amount of lattice disorder on the crystalline surface after being treated under reducing conditions. Such layers usually contain a significant amount of oxygen vacancies and the low valence metal centers (selfdoped), which may create the mid-gap states below the Fermi level of the metal oxide. 16 The mid-gap states form a continuum that extends to and overlaps with the conduction band (CB) or merges with the valence band (VB). This can effectively reduce the band gap and enhance the light absorption in visible and near-IR regions. 17, 18 These extended energy states and the newly produced energy levels may become dominant for optical excitation and relaxation, and thus new application may be endowed. Furthermore, the localized holes generated in the mid-gap state of the disordered layer efficiently inhibit the charge recombination, 16, 19-22 where the low valence metal centers act as a trapping sites for photogenerated carriers.

The pioneering work of black TiO₂ with disordered surface demonstrated a significant enhancement in photocatalytic hydrogen evolution reaction (HER) rate and has sparked the investigations of creating disordered layer in other metal oxides, including Fe₂O₃,²³ ZnO,²⁴ BiVO₄,²⁵ and WO₃.^{26, 27} Among them, tungsten oxide (WO_3) is an important *n*-type semiconductor of various applications ranging from dyesensitized solar cell²⁸ to electrochromic devices,²⁹ gas sensors,³⁰ and photoelectrochemical water splitting,31 owing to low cost and good conductivity. With the desirable bandgap energy (2.5-2.8 eV) for capturing visible light and a number of stable oxygendeficient sub-stoichiometric structures (WO_{3-x}, where $0 < x \le$ 0.28) available, WO₃ is one of the most attractive candidates for photocatalytic water splitting.32 Nonetheless, restricted by its intrinsic electronic structure, WO₃ is only suitable for water oxidation.33 The positon of valence band maximum (VBM) of WO₃ is located at a relatively high energy level (3.0 eV vs.

^a Department of Applied Biology and Chemical Technology and the State Key Laboratory of Chemical Biology and Drug Discovery, The Hong Kong Polytechnic University, Hung Hom, Hong Kong SAR, China.

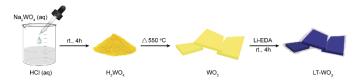
b. Department of Chemical and Biomolecular Engineering, Yonsei University, 50 Yonsei-ro, Seodaemun-gu, Seoul 120-749, Republic of Korea.

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reversible hydrogen electrode, RHE), whereas the conduction band minimum (CBM) lies at a position lower than hydrogen reduction potential (0 V vs. RHE).

Recently, thin disordered layer on WO₃ surface has been reported to significantly increase the charge carrier density and facilitate the interface charge transfer, thus improve the charge separation efficiency and conductivity in photoelectrocatalytic water oxidation.^{26, 27} Localized surface atomic disorder was also engaged to turn the inert WO_3 to an active electrocatalyst for hydrogen generation from acidic water.³⁴ Nevertheless, it is still challenging to trigger solar water reduction using WO₃ and the effect of surface disordered layer on the band structure of WO₃ has not been clearly understood yet. Herein, we demonstrate that a simple room-temperature solution processing method can modify the surface morphology and electronic structure of WO₃ to enable the unprecedented capability of photocatalytic HER without co-catalyst. Using Li-ethylenediamine (Li-EDA) as a disordering agent, surface disordered layer was successfully created on the crystalline WO₃ nanoplates (see Scheme 1 for synthetic schematics). Experimental results confirm the formation of mid-gap state that causes the blue-shift in both CBM and Fermi level (E_f) with respect to pristine WO₃. The reformed band structure of Li-EDA treated WO₃ (LT-WO₃) allows the photocatalytic hydrogen evolution without the assistance of co-catalyst at a steady rate of 94.2 µmol g⁻¹ h⁻¹ under simulated sun light. With the first example of using WO₃ as a co-catalystfree photocatalyst for HER, we demonstrate that Li-EDA treatment can be applied as a simple yet effective strategy to tailor the optical band structure of metal oxide via the formation of thin disordered surface layer.



Scheme 1. Schematic illustration of synthetic procedure for WO₃ with surface disordered layer (LT-WO₃).

EXPERIMENTAL SECTION

Synthesis of WO₃ nanoplates.

Sodium tungstate solution (0.4 mol L $^{-1}$) was added to excess hydrochloric acid solution (1 mol L $^{-1}$) under vigorous stirring to obtain the tungstic acid nanoparticle precipitate. The yellow precipitate was collected after centrifugation, and washed with water for 3 times. Afterwards, the as-prepared solids were calcinated at 550 °C under Ar atmosphere for 3 h. The product was collected after the furnace cooled down to room temperature.

Synthesis of LT-WO3 nanoplates.

The as-synthesized WO $_3$ nanoplate was immersed into the solvated electrons solution, Li-EDA (1 mol L $^{-1}$), at room temperature with continuous stirring for 4 hours. After quenching the reaction with deionized water, the resulting dark

blue precipitate was collected using centrifugation. The solids were washed with HCl solution (0.5 mol L^{-1}) and ethanol alternately in order to remove any physically adsorbed Li and EDA. It was then dried under vacuum at room temperature overnight.

Photocatalytic H₂ production.

In a typical run, 5 mg sample was suspended in a 25 mL mixture solution of Na₂S (0.35 M) and Na₂SO₃ (0.25 M). The solution was degassed with N₂ for 30 minutes. Afterwards, the as-prepared solution was irradiated with a Newport solar simulator (150 W Xe lamp, ozone free, Air Mass Filter, AM 1.5 Global) of 100 mW cm⁻² light intensity under stirring. The amount of generated H₂ was analyzed by an Agilent 7890B gas chromatograph system equipped with a TCD detector using N₂ as carrier gas.

Sample characterization.

The crystal structure of samples was identified by X-ray diffractometer (XRD, Rigaku SmartLab, 9kW) with a diffraction angle (2θ) ranging from 10 to 70°. TEM images were obtained using a STEM (JEOL JEM-2100F) operated with a field emission gun at 200 kV. The samples were first dispersed in ethanol and drop-cast onto holey carbon-coated 400 mesh copper grids. The UV-vis diffuse reflectance spectroscopy was recorded using a UV-Vis spectrophotometer (Cary 4000) with ultrapure BaSO₄ as the reference. Raman spectra were obtained by a Raman spectrometer (Renishaw inVia) with the excitation wavelength of 785 nm at room temperature. XPS measurements were performed using a K-Alpha X-ray photoelectron spectrometer (Thermo Fisher Scientific, UK) with a monochromatic Al Kα X-ray source (excitation energy = 1,468.6 eV) and a base pressure of 3×10^{-9} Torr in the analytical chamber, and with an internal reference of C 1s at 284.0 eV. EPR measurements were performed on an EPR spectrometer (SPINSCAN X, ADANI) operating at 9.3 GHz (X band) using a 100 kHz field modulation. All the spectra were recorded at 298 K using a microwave power of 10 mW and accumulated for three times.

Electrochemical measurements.

All the measurements were performed on a CHI1030A electrochemical workstation. 4 mg catalyst and 200 μL 5 wt% Nafion solution were dispersed in 4 mL water/ethanol (3:1 v/v) by 10 min sonication to form homogeneous suspension. Then, 10 μL catalyst ink was drop-cast onto a glassy carbon electrode of 3 mm diameter (surface area = 7.07 mm²) as the working electrode with a constant mass loading of 0.14 mg cm⁻². A linear sweep voltammogram from 0 to -0.6 V was measured at a scan rate of 1 mV s^{-1} in 0.5 M H_2SO_4 using a saturated calomel electrode and a Pt wire as the reference electrode and counter electrode, respectively. Pt/C (20 wt% platinum on graphitized carbon) is commercial available from Sigma-Aldrich. The Mott-Schottky plot was constructed based on the electrochemical data obtained using a typical three electrode system in phosphate buffer solution of pH 7.4. The Ag/AgCl (in 3 M NaCl solution) electrode and graphite were used as reference and counter electrode, respectively. The data were collected by Journal Name ARTICLE

measuring the space-charge layer capacitance over a range of applied potentials (from -1.0 to 1.0 V) at 1,000 Hz using a CHI760E electrochemical station. The recorded potential was referenced to the reversible hydrogen electrode (RHE) according to the following equation: E (vs. RHE) = E (vs. Ag/AgCl) + 0.197 V + $0.0592 \times pH$.

Results and discussion

Monoclinic WO₃ nanoplate was first synthesized by directly annealing tungstic acid at 550 °C, which was afforded by a simple co-precipitation of sodium tungstate solution with HCl. To introduce disordered layer, the as-prepared WO₃ nanoplates were dispersed and stirred in Li-EDA solution³⁵ at room temperature for 4 h. During the Li-EDA treatment, a distinct color change from pale yellow to dark blue was observed (Figure 1a). The UV-vis diffusive reflectance spectra (UV-vis DRS, Figure 1b) reveal that the light response of Li-EDA treated WO₃ (LT-WO₃) has been extended, as evident from a significant absorption tail in the visible region (450 to 800 nm).

Such change in optical properties can arise from the oxygen vacancies created on the surface or modified crystal phase.^{36, 37} X-ray diffraction (XRD) analysis was conducted to investigate the structural changes induced by Li-EDA treatment. As shown in Figure 1c, the XRD pattern of pristine WO₃ matches well with the monoclinic phase of WO₃ (JCPDS no. 43-1035) where the sharp diffraction peaks indicate its crystalline nature. LT-WO₃ also shows a similar pattern of diffraction peaks, indicating that the crystal structure is retained. However, the peaks are broader with much weaker intensities compared to those observed from pristine WO₃, suggesting reduced crystallinity due to the localized modification of crystal structure.³⁶

Similar change was observed from the Raman spectra (Figure 1d), where the characteristic Raman peaks of WO_3 are broadened after the Li-EDA treatment. Two peaks at 273 and 326 cm⁻¹, which are assigned to the W⁶⁺-O-W⁶⁺ bending vibration mode of bridging oxygen in WO₃ crystal lattice,³⁸ almost disappeared in LT-WO₃ spectrum, suggesting the partial removal of oxygen atoms. The peak at 716 cm⁻¹ that is attributed to the long W⁶⁺–O–W⁶⁺ bond stretching mode was also markedly reduced with a slight shift. In addition, the peak for the short W⁶⁺–O–W⁶⁺ bond stretching at 808 cm⁻¹ shows an appreciable broadening that is correlated to the weakening of W-O bond.³⁸ Furthermore, a small peak at 430 cm⁻¹ appeared after the Li-EDA treatment indicates the presence of W species at lower oxidation states.³⁸⁻⁴⁰ The reduction of W⁶⁺ to a lower oxidation state by Li-EDA is believed to weaken the W-O bond, which in turn leads to the partial loss of oxygen from the lattice. This is supported by the electron paramagnetic resonance (EPR) studies where only LT-WO₃ shows a signal at g-value of 2.002, an indication of oxygen vacancies (Figure 2).14 Both Raman and EPR results are in good agreements with the observations from XRD measurements.

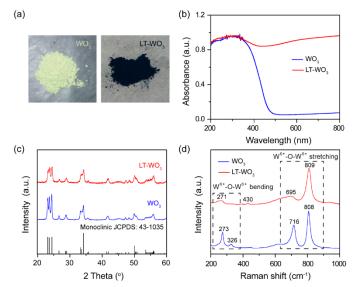


Figure 1. (a) Photographs of pristine WO_3 and LT- WO_3 obtained via Li–EDA treatment. (b) UV-vis diffusive reflectance spectra (UV-vis DRS) of WO_3 and LT- WO_3 . (c) XRD and (d) Raman spectra of WO_3 and LT- WO_3 .

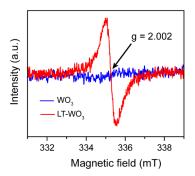


Figure 2. EPR spectra of WO₃ and LT-WO₃.

Figure 3 compares the transmission electron microscopic (TEM) images of as-prepared WO₃ and LT-WO₃ nanoplates. The typical morphology of WO₃ is a rectangular nanoplates (mean side length of ca. 130 nm) with some smaller nanoplates of irregular shapes (Figure 3a). A high resolution TEM (HR-TEM) image shown in Figure 3b clearly reveals well-ordered lattice fringes of 0.366 and 0.378 nm that can be indexed to the (200) and (020) planes of monoclinic phase WO₃, respectively. The corresponding selected area electron diffraction (SAED) pattern (inset in Figure 3b) shows the single-crystal spots that are in good agreements with the (200) and (020) planes. On the other hand, the TEM image of LT-WO₃ (Figure 3c) shows a similar but more round-shaped morphology with a slightly reduced average side length (ca. 77 nm). Moreover, from the HR-TEM (Figure 3d), thin disordered layer (average thickness = ca. 2 nm) is evident along the edge, which is supported by the disappearance of lattice fringes and its corresponding SAED pattern (inset in Figure 3d). The formation of disordered layer is attributed to the deprivation of surface lattice oxygen by solvated electrons during the treatment with Li-EDA, which is a strong reducing agent.⁴¹ By reacting with Li-EDA, W⁶⁺ is reduced to a lower oxidation state, and oxygen atoms are removed from the surface lattice. 18, 35, 42 The removal of oxygen atoms would

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lead to the crystal lattice deformation on the surface, leading to the decrease in both crystallinity and size of the nanoplates as observed from the XRD, TEM, and Raman spectroscopic studies.

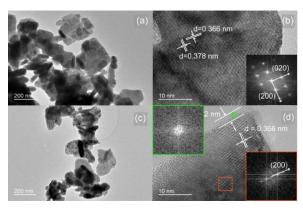


Figure 3. Electron microscopic analysis of WO_3 and LT- WO_3 . (a) Low-resolution and (b) high-resolution TEM images of WO_3 . (c) Low-resolution and (d) high-resolution TEM images of LT- WO_3 . The insets in (b) and (d) show the FFT patterns of selected areas (orange squares: monoclinic phase core; green squares: disordered layer).

It is interesting that the core part of LT-WO₃ still retained the crystallinity as manifested by the clear lattice fringes (Figure 3d). We have conducted a series of time-resolved Li-EDA treatments to investigate the formation of disordered layer and control its thickness (Figure 4). The disordered layer starts to form on the edges after 2 h of reaction and gradually grows thicker with treatment time. The thickness of disordered layer reached up to 5 nm after 7 h, however, further prolonged reaction did not change the thickness, but only reduced the nanoplate size. The reaction between Li-EDA and WO₃ initially occurs at the Li-EDA/WO₃ interface. At early stage, solvated electrons can diffuse through thin disordered layer formed on the surface to reduce the W6+ ions in the inner crystal lattice. As the disordered layer becomes thicker, it acts as a passivation layer that hinders the diffusion of the electrons into the inner part. With further extended reaction time, the size of core shrinks, as the outer layer is lost. These results suggest that the surface disordered layer is closely associated with the formation of oxygen vacancies.

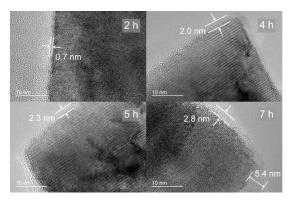


Figure 4. Time-resolved TEM analysis of the thickness of disordered layer in LT-WO₃ nanoplates treated for 2, 4, 5, and 7 h.

To confirm the formation of disordered layer, the surface species was characterized using X-ray photoelectron

spectroscopy (XPS). Similar survey XPS spectra were obtained from pristine WO₃ and LT-WO₃, where only W and O were identified, with a small C 1s peak due to the adsorbed carboncontaining species, eliminating the possibility of Li contamination (binding energy of Li 1s is ca. 55.0 eV), especially in LT-WO₃ (Figure 5). From the normalized W 4f core level spectra (Figure 6a), only W6+ ions are detected in pristine WO3, with the peaks at binding energy of 35.7 and 37.8 eV assigned to W $4f_{5/2}$ and W $4f_{7/2}$. In contrast, the W 4f spectrum of LT-WO₃ shifts to a lower binding energy and becomes broader, indicates the presence of lower oxidation states. In addition to W⁶⁺ state, W⁵⁺ ions appear as a shoulder at binding energies of 34.5 and 36.4 eV. This is in good accordance with the Raman studies and further validates the reduction of W6+ to a lower oxidation state by Li-EDA. Additionally, a broad peak was seen from the O 1s spectrum of LT-WO₃ (Figure 6b), which can be deconvoluted into two peaks. The peak centered at 530.4 eV is ascribed to the lattice O²⁻ species (W–O–W),¹⁴ while the other peak centered at 531.6 eV is assigned to O2- in the vicinity of oxygen vacancies.43 Meanwhile, the O 1s spectrum of pristine WO₃ shows a major peak for lattice oxygen atoms at 530.3 eV with a small peak at 531.6 eV, presumably due to the surface contamination during sample preparation. It is noted that both W 4f and O 1s peaks of LT-WO₃ exhibit binding energies slightly shifted to the lower values, ca. 0.2 and 0.1 eV, respectively. Similar shifts were observed from WO₃ after the chemical treatment as a result of the change in the electronic structure.^{26,}

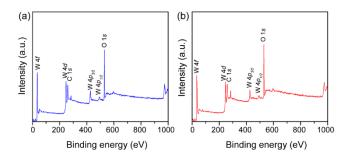


Figure 5. XPS survey spectra of (a) WO₃ and (b) LT-WO₃.

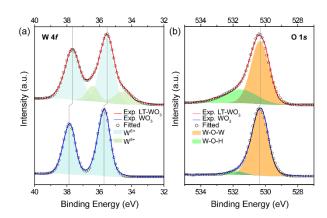


Figure 6. (a) W 4f and (b) O 1s XPS spectra of LT-WO₃ (top) and WO₃ (bottom).

Disordered atomic arrangement is known to affect the electronic structure, and thus physicochemical properties of

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WO₃. 16, 25, 26 The observed change in the optical absorption of LT-WO₃ is related to the surface disordered layer. The band gap energies calculated by fitting the UV-vis DRS to Tauc equation are 2.68 and 2.74 eV for WO₃ and LT-WO₃, respectively (Figure S1). Despite the obvious color change, the optical band gap energies (E_g) did not show much difference. However, the valence band XPS (VB-XPS) spectra reveal the changes in the electronic structure and band offset (Figure 7a). The position of VB maximum (VBM) was first evaluated by linearly extrapolating the onset of XPS valence band spectra to the baseline, which reflects a band edge position with respect to the Fermi level (E_f). The VBM edge levels of WO₃ and LT-WO₃ were determined to be 2.63 and 2.43 eV, respectively, indicating the upshift of VBM by induced oxygen vacancies. Interestingly, the VB-XPS spectrum of LT-WO₃ also exhibits the VB tail state (1.31 eV with respect to E_f), suggesting the formation of additional electronic state. This mid-gap state would essentially broaden the VB width to create a reduced band gap ($E_g' = 1.62 \text{ eV}$), and consequently allows the wider visible light absorption and change in sample color to dark blue.⁴⁵ We also engaged the Mott–Schottky plots to determine E_f levels and thus obtain the band structures of samples. 46, 47 Figure 7b shows the flat band potentials, which shifted from 0.5 to 0.25 V after Li-EDA treatment. The flat band potential reflects the difference between the Ef level and the water reduction potential of the sample. This shift of E_f confirms again that the surface disordered layer reconstructs the electronic structure of WO₃ by introducing oxygen vacancies. The formation of mid-gap is closely related to the disordered surface layer on LT-WO₃, which contains W5+ ions and oxygen vacancies (Figure 6). The selfdoping of W5+ ions also modifies the electronic structure by creating a mid-gap and enhances the surface electron density, contributing to the upshift of the E_f. 16, 48, 49

Combining the results from UV-Vis DRS, VB-XPS, and Mott–Schottky plot, we have constructed an energy level diagram qualitatively representing the changes in band structure of LT-WO₃, and presented in Figure 7c. The conduction band minimum (CBM) was estimated according to the equation of $E_{\text{CBM}} = E_{\text{VBM}} - E_{\text{g}}$. The upshift of the E_{f} induced by the disordered layer on LT-WO₃ caused the VBM shift from 3.13 to 2.68 V. With the E_{g} of 2.74 eV, the CBM in LT-WO₃ is determined to be –0.06 V, which satisfies the thermodynamic requirement for H₂ generation.

To verify this reformed band structure of LT-WO₃, we have carried out the photocatalytic HER using a mixed solution of Na₂S and Na₂SO₃ as a sacrificial agent under one sun illumination. LT-WO₃ indeed shows a remarkable performance in photocatalytic HER with the average rate of 94.2 μ mol g⁻¹ h⁻¹ (Figure 8). The H₂ generation rate was stable during the continuous 25 h reaction, showing no discernable degradation in the catalytic activity. The control experiment conducted under the dark generated no detectable H₂. To the best of our knowledge, this is the first report of photocatalytic H₂ evolution from WO₃ without any co-catalysts. The creation of disordered surface layer on WO₃ is demonstrated to modify its electronic structure, thus endows a new capability of photocatalytic H₂ generation. It also seems to affect the generation and lifetime

of charge carriers. The steady state photoluminescence (PL) spectra (Figure S2a) taken at room temperature show a much enhanced PL emission intensity, attributed to the higher concentration of photogenerated excitons. The time-resolved PL spectra (Figure S2b) indicates that the charge carriers in LT-WO₃ have more than 4 times longer lifetime (1.05 \pm 0.05 ns) than those in WO₃ (0.25 \pm 0.05 ns).

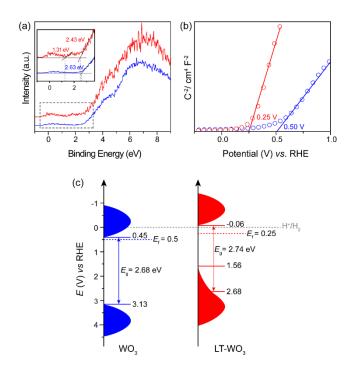


Figure 7. (a) VB-XPS and (b) Mott-Schottky plot of WO_3 and LT- WO_3 . Inset in (a) is the enlarged area near the band edges indicated with a dotted box. (c) Schematic illustration of the energy levels of WO_3 and LT- WO_3 nanoplates.

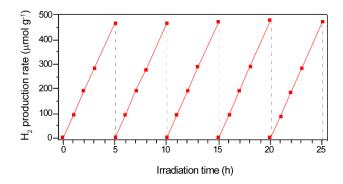


Figure 8. Photocatalytic H_2 production using LT-WO $_3$ as a photocatalyst for 25 h reaction of five cycles.

In order to gain further insight on the disordered surface layer of LT-WO₃, the surface reactivity was studied by examining the electrocatalytic activity in H_2 production. Figure 9a shows the polarization curves where the normalized current densities are plotted against applied potential without iR correction. LT-WO₃ shows a much lower overpotential (η) of -176 mV (vs. RHE, same below) at the current density of 10 mA cm⁻² compared with pristine WO₃ (-667 mV). The corresponding Tafel slope of

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LT-WO₃ is 59 mV dec⁻¹ (Figure 9b), lower than that of WO₃ (88 mV dec⁻¹), implying its higher activities towards H₂ production. These results indicate that the disordered layer on WO₃ surface can significantly improve the surface reactivity towards H₂ production. The increased population of trapping sites for electron carriers (*i.e.*, oxygen vacancies) is believed to promote the electron transfer, which effectively reduces the energy barrier for proton reduction, and activates the surface reactivity towards H₂ production. 26,27

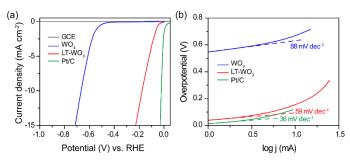


Figure 9. (a) Polarization curves and (b) Tafel plots of WO_3 and LT- WO_3 with Pt/C electrodes included for comparison.

Conclusions

In conclusion, the band structure of WO $_3$ has been successfully tuned by a simple surface defect engineering method, Li-EDA treatment. The formation of oxygen vacancies and self-doping of W $^{5+}$ onto the surface of WO $_3$ through the generation of disordered surface layer create a mid-gap level in the electronic structure and shift the CBM above the H $_2$ reduction potential. This is the first example of tuning the band structure of WO $_3$ to realize co-catalyst-free photocatalytic H $_2$ evolution. Electrochemical investigations also reveal that the oxygen vacancy-rich disordered layer can provide efficient electron trapping sites that promote the surface activity for HER. Our work demonstrates the disordered layer on a semiconductor surface can be utilized to effectively tune the band structure and extend its catalytic applications.

Conflicts of interest

There are no conflicts to declare.

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Notes and references

- T. Hisatomi, J. Kubota and K. Domen, Chem. Soc. Rev., 2014, 43, 7520-7535.
- 2. J. H. Montoya, L. C. Seitz, P. Chakthranont, A. Vojvodic, T. F. Jaramillo and J. K. Nørskov, *Nat. Mater.*, 2016, **16**, 70.
- 3. Y. Wang, H. Suzuki, J. Xie, O. Tomita, D. J. Martin, M. Higashi, D. Kong, R. Abe and J. Tang, *Chem. Rev.*, 2018, **118**, 5201-5241.
- S. Fukuzumi, Y.-M. Lee and W. Nam, Coord. Chem. Rev., 2018, 355. 54-73.
- E. Ha, L. Y. S. Lee, J. Wang, F. Li, K.-Y. Wong and S. C. E. Tsang, *Adv. Mater.*, 2014, 26, 3496-3500.
- J. Zhang, W. Yao, C. Huang, P. Shi and Q. Xu, J. Mater. Chem. A, 2017, 5, 12513-12519.
- W. Liu, E. Ha, L. Wang, L. Hu, L. Y. S. Lee and K.-Y. Wong, Adv. Mater. Interfaces, 2018, DOI: doi:10.1002/admi.201800611, 1800611.
- I. Concina, Z. H. Ibupoto and A. Vomiero, Adv. Energy Mater., 2017, 7, 1700706.
- 9. G. Wang, Y. Yang, D. Han and Y. Li, *Nano Today*, 2017, **13**, 23-39.
- 10. P. Kofstad, Phase Transit., 1996, 58, 75-93.
- J. Nowotny, M. A. Alim, T. Bak, M. A. Idris, M. Ionescu, K. Prince, M. Z. Sahdan, K. Sopian, M. A. Mat Teridi and W. Sigmund, *Chem. Soc. Rev.*, 2015, 44, 8424-8442.
- J. Jia, C. Qian, Y. Dong, Y. F. Li, H. Wang, M. Ghoussoub, K. T. Butler, A. Walsh and G. A. Ozin, *Chem. Soc. Rev.*, 2017, 46, 4631-4644.
- X. Pan, M.-Q. Yang, X. Fu, N. Zhang and Y.-J. Xu, Nanoscale, 2013, 5, 3601-3614.
- J. Meng, Q. Lin, T. Chen, X. Wei, J. Li and Z. Zhang, Nanoscale, 2018, 10, 2908-2915.
- 15. Y. Yang, Y. Wang and S. Yin, Appl. Surf. Sci., 2017, 420, 399-406.
- X. Chen, L. Liu, P. Y. Yu and S. S. Mao, Science, 2011, 331, 746-750.
- 17. Y. Liu, Y. Li, S. Yang, Y. Lin, J. Zuo, H. Liang and F. Peng, ChemSusChem, 2018, DOI: doi:10.1002/cssc.201800940.
- 18. G. Yin, X. Huang, T. Chen, W. Zhao, Q. Bi, J. Xu, Y. Han and F. Huang, *ACS Catal.*, 2018, **8**, 1009-1017.
- P. Yan, G. Liu, C. Ding, H. Han, J. Shi, Y. Gan and C. Li, ACS Appl. Mater. Interfaces, 2015, 7, 3791-3796.
- G. Wang, H. Wang, Y. Ling, Y. Tang, X. Yang, R. C. Fitzmorris, C. Wang, J. Z. Zhang and Y. Li, *Nano Lett.*, 2011, 11, 3026-3033.
- S. W. Kwon, M. Ma, M. J. Jeong, K. Zhang, S. J. Kim and J. H. Park, Chem. Commun., 2016, 52, 13807-13810.
- Y. Cho, S. Kim, B. Park, C. L. Lee, J. K. Kim, K. S. Lee, I. Y. Choi, J. K. Kim, K. Zhang, S. H. Oh and J. H. Park, *Nano Lett*, 2018, 18, 4257-4262.
- S. Sun, T. Zhai, C. Liang, S. V. Savilov and H. Xia, *Nano Energy*, 2018, 45, 390-397.
- 24. T. Xia, P. Wallenmeyer, A. Anderson, J. Murowchick, L. Liu and X. Chen, *RSC Adv.*, 2014, **4**, 41654-41658.
- K. J. Kyu, C. Yoonjun, J. M. Jin, L.-W. Ben, S. Dongguen, Y. Yeonjin, W. D. Hwan, Z. Xiaolin and P. J. Hyeok, *ChemSusChem*, 2018, 11, 933-940.
- M. Ma, K. Zhang, P. Li, M. S. Jung, M. J. Jeong and J. H. Park, Angew. Chem. Int. Ed., 2016, 55, 11819-11823.
- R. Zhang, F. Ning, S. Xu, L. Zhou, M. Shao and M. Wei, *Electrochim. Acta*, 2018, 274, 217-223.
- 28. H. Zheng, Y. Tachibana and K. Kalantar-zadeh, *Langmuir*, 2010, **26**, 19148-19152.
- S.-H. Lee, R. Deshpande, P. A. Parilla, K. M. Jones, B. To, A. H. Mahan and A. C. Dillon, *Adv. Mater.*, 2006, **18**, 763-766.
- 30. R. Xing, Y. Du, X. Zhao and X. Zhang, Sensors, 2017, 17, 710.

Journal Name ARTICLE

- 31. X. Shi, I. Y. Choi, K. Zhang, J. Kwon, D. Y. Kim, J. K. Lee, S. H. Oh, J. K. Kim and J. H. Park, *Nat. Commun.*, 2014, **5**, 4775.
- 32. H. Zhen-Feng, S. Jiajia, P. Lun, Z. Xiangwen, W. Li and Z. Ji-Jun, *Adv. Mater.*, 2015, **27**, 5309-5327.
- 33. C. A. Bignozzi, S. Caramori, V. Cristino, R. Argazzi, L. Meda and A. Tacca, *Chem. Soc. Rev.*, 2013, **42**, 2228-2246.
- 34. Y. H. Li, P. F. Liu, L. F. Pan, H. F. Wang, Z. Z. Yang, L. R. Zheng, P. Hu, H. J. Zhao, L. Gu and H. G. Yang, *Nat. Commun.*, 2015, **6**, 8064.
- 35. K. Zhang, L. Wang, J. K. Kim, M. Ma, G. Veerappan, C.-L. Lee, K.-j. Kong, H. Lee and J. H. Park, *Energy Environ. Sci*, 2016, **9**, 499-503.
- 36. G. Wang, Y. Ling, H. Wang, X. Yang, C. Wang, J. Z. Zhang and Y. Li, Energy Environ. Sci, 2012, 5.
- J. Song, Z.-F. Huang, L. Pan, J.-J. Zou, X. Zhang and L. Wang, ACS Catal., 2015, 5, 6594-6599.
- 38. S. Bai, K. Zhang, L. Wang, J. Sun, R. Luo, D. Li and A. Chen, *J. Mater. Chem. A*, 2014, **2**, 7927-7934.
- S.-H. Lee, H. M. Cheong, J.-G. Zhang, A. Mascarenhas, D. K. Benson and S. K. Deb, *Appl. Phys. Lett.*, 1999, **74**, 242-244.
- 40. W. Chun-Kai, D. R. Sahu, W. Sheng-Chang, L. Chung-Kwei and H. Jow-Lay, *J. Phys. D*, 2012, **45**, 225303.

- 41. R. S. Monson, in *Advanced Organic Synthesis*, ed. R. S. Monson, Academic Press, 1971, DOI: https://doi.org/10.1016/B978-0-12-504950-4.50006-X, pp. 25-30.
- 42. M. Ming, Z. Kan, L. Ping, J. M. Sun, J. M. Jin and P. J. Hyeok, *Angew. Chem. Int. Ed.*, 2016, **55**, 11819-11823.
- 43. Y. Sun, W. Wang, J. Qin, D. Zhao, B. Mao, Y. Xiao and M. Cao, *Electrochim. Acta*, 2016, **187**, 329-339.
- 44. H. Cheng, M. Klapproth, A. Sagaltchik, S. Li and A. Thomas, *J. Mater. Chem. A*, 2018, **6**, 2249-2256.
- 45. J. Liu, L. Han, N. An, L. Xing, H. Ma, L. Cheng, J. Yang and Q. Zhang, *Appl. Catal.*, *B*, 2017, **202**, 642-652.
- J. Liu, J. Ke, D. Li, H. Sun, P. Liang, X. Duan, W. Tian, M. O. Tadé,
 S. Liu and S. Wang, ACS Appl. Mater. Interfaces, 2017, 9, 11678-11688.
- L. Liao, Q. Zhang, Z. Su, Z. Zhao, Y. Wang, Y. Li, X. Lu, D. Wei, G. Feng, Q. Yu, X. Cai, J. Zhao, Z. Ren, H. Fang, F. Robles-Hernandez, S. Baldelli and J. Bao, *Nat. Nanotechnol*, 2013, 9, 69.
- 48. F. Zuo, L. Wang, T. Wu, Z. Zhang, D. Borchardt and P. Feng, *J. Am. Chem. Soc.*, 2010, **132**, 11856-11857.
- 49. W. Huang, J. Wang, L. Bian, C. Zhao, D. Liu, C. Guo, B. Yang and W. Cao, *Phys. Chem. Chem. Phys.*, 2018, **20**, 17268-17278.